J-PARCにおける 核子スピン構造の解明

第２回「J-PARCにおける高エネルギーハドロン物理」

２００７年１０月１９日（金）

後藤雄二（理研/RBRC)

セミナーの内容

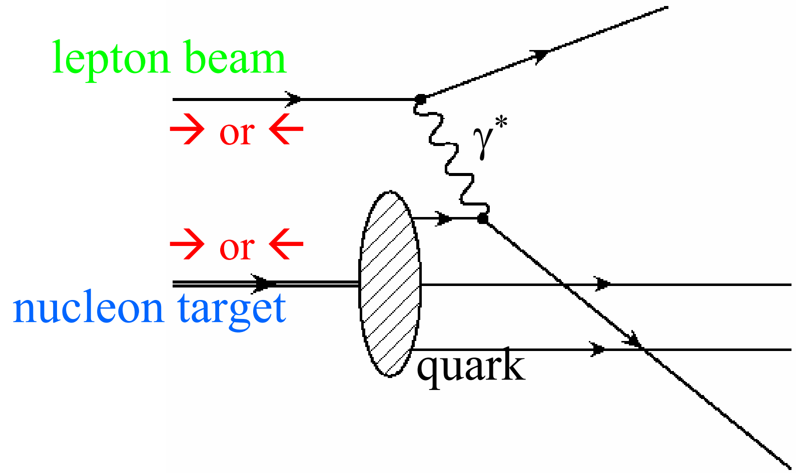
- 長いイントロ
  - 核子スピン構造
  - RHICその他の偏極実験の状況、結果
    - 縦偏極実験
    - 横偏極実験
- J-PARC Drell-Yan実験の物理
  - 縦偏極実験：核子スピンのフレーバー構造
  - 横偏極実験：軌道角運動量の寄与
  - その他

核子スピン構造

- Fundamentalな対象であるにもかかわらず、理解されていない
- QCDを基盤とする研究方法が発展している
- 核子の構造を調べるともに、QCDに対するテスト、理解となる
  - Q^2 evolution + factorization + universality
  - global QCD analysis of e^++e^-, e+p, and p+p(or pbar) data
    - unpolarized/polarized parton distribution functions
    - fragmentation functions
- 縦偏極実験と横偏極実験
  - クォークスピンとグルーオンスピンの寄与
  - クォークとグルーオンの軌道角運動量の寄与

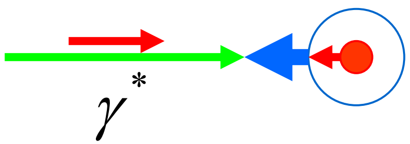
核子スピン1/2の起源？

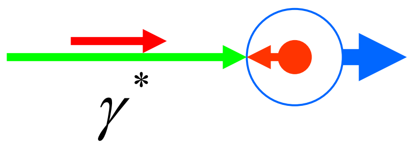
- 偏極レプトン深非弾性散乱実験
  - パートン模型



$$\Delta q(x) = q^{+}(x) - q^{-}(x)$$

$$S_\gamma + S_N = 1/2 \qquad\qquad S_\gamma + S_N = 3/2$$





$$\sigma^{T}_{1/2} \sim \sum_i e_i^2 q^{+}(x) \qquad\qquad \sigma^{T}_{3/2} \sim \sum_i e_i^2 q^{-}(x)$$

$$A_1 = \frac{\sigma^{\mathrm{T}}_{1/2} - \sigma^{\mathrm{T}}_{3/2}}{\sigma^{\mathrm{T}}_{1/2} + \sigma^{\mathrm{T}}_{3/2}} \sim \frac{\sum_i e_i^2 (q_i^{+}(x) - q_i^{-}(x))}{\sum_i e_i^2 (q_i^{+}(x) + q_i^{-}(x))} = \frac{\sum_i e_i^2 \Delta q_i(x)}{\sum_i e_i^2 q_i(x)} = \frac{g_1(x)}{F_1(x)}$$

偏極構造関数 ($g_1(x)$)

非偏極構造関数 ($F_1(x)$)

核子スピン1/2の起源？

- 偏極レプトン深非弾性散乱実験

$$A_1 = \frac{\sigma^{\mathrm{T}}_{1/2} - \sigma^{\mathrm{T}}_{3/2}}{\sigma^{\mathrm{T}}_{1/2} + \sigma^{\mathrm{T}}_{3/2}} \sim \frac{\sum_i e_i^2 (q_i^+(x) - q_i^-(x))}{\sum_i e_i^2 (q_i^+(x) + q_i^-(x))} = \frac{\sum_i e_i^2 \Delta q_i(x)}{\sum_i e_i^2 q_i(x)} = \frac{g_1(x)}{F_1(x)}$$

more correctly

$$A_1 = \frac{g_1(x) - \gamma^2 g_2(x)}{F_1(x)} \qquad A_2 = \frac{2\sigma^{TL}}{\sigma^{\mathrm{T}}_{1/2} + \sigma^{\mathrm{T}}_{3/2}} = \frac{\gamma^2 (g_1(x) + g_2(x))}{F_1(x)}$$

  - 実験データ

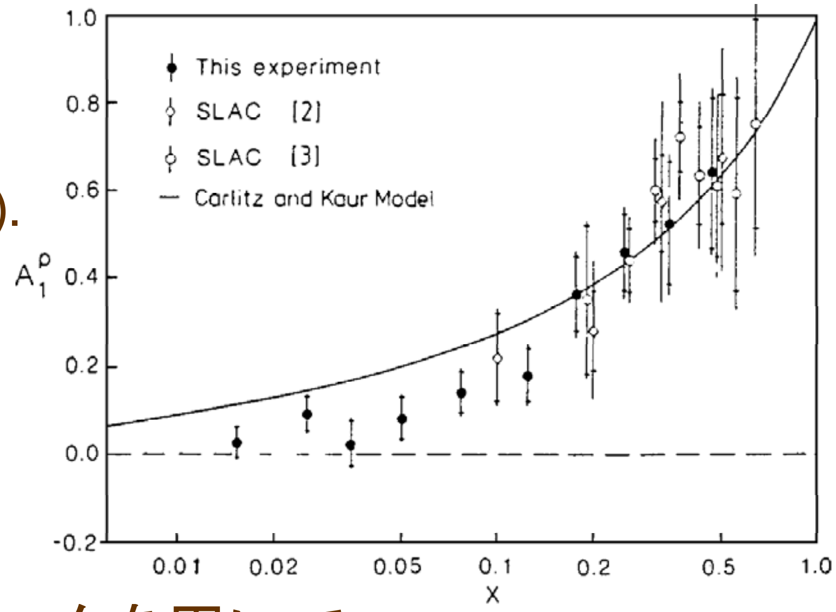
$$A = \frac{1}{P_T P_B} \frac{N_{++} - N_{+-}}{N_{++} + N_{+-}} = D(A_1 + \eta A_2) \sim DA_1$$

P_T target polarization

P_B beam polarization

D depolarization factor

核子スピン1/2の起源？



- EMC実験@CERN

  J. Ashman et al., NPB 328, 1 (1989).

$$\int_0^1 dx g_1^p(x) = \frac{1}{2}\left[\frac{4}{9}\Delta u + \frac{1}{9}\Delta d + \frac{1}{9}\Delta s\right]$$

$$= 0.123 \pm 0.013(\mathrm{stat}) \pm 0.019(\mathrm{syst})$$

  – 中性子およびハイペロン崩壊データを用いて

$$\Delta\Sigma = \Delta u + \Delta d + \Delta s = 12 \pm 9(\mathrm{stat}) \pm 14(\mathrm{syst})\%$$ 「陽子スピンの危機」

    - クォークスピンは核子スピンの小さな割り合いにしか寄与しない

  – x = 0 ~ 1 の積分による不確定性

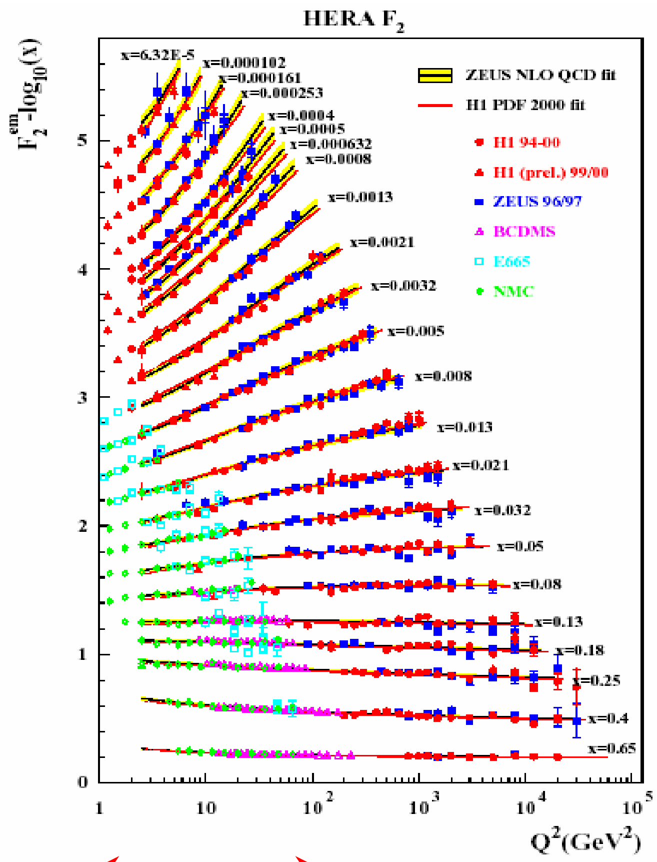
    - より広いx領域を覆う、よりよい精度のデータが必要

➔ SLAC/CERN/DESY/JLAB 実験

偏極レプトン深非弾性散乱実験

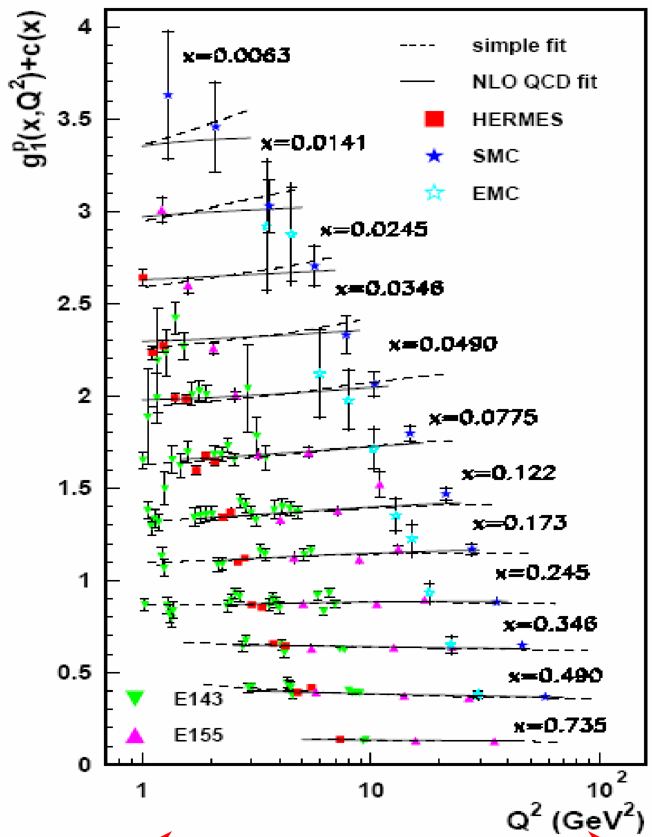
- 固定ターゲット実験
  - *Q^2の範囲が限られている*

unpolarized DIS



$1 < Q^2 < 100\ (\mathrm{GeV}/c)^2$

polarized DIS



$1 < Q^2 < 100\ (\mathrm{GeV}/c)^2$

偏極レプトン深非弾性散乱実験

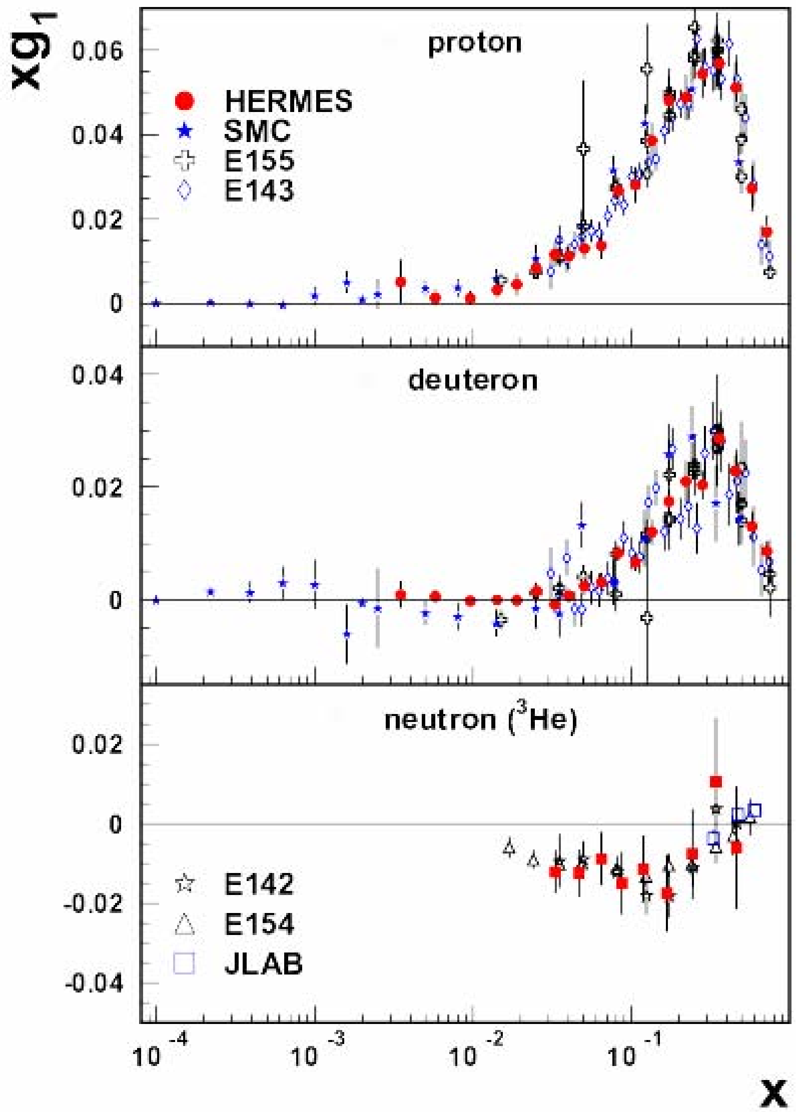
- クォークスピンの寄与

$$\Delta\Sigma \sim 0.2$$

- 核子スピン1/2の起源は何か？

$$\frac{1}{2} = \frac{1}{2}\Delta\Sigma + \Delta g + L$$

  – グルーオンスピンの寄与？
  – 軌道角運動量？

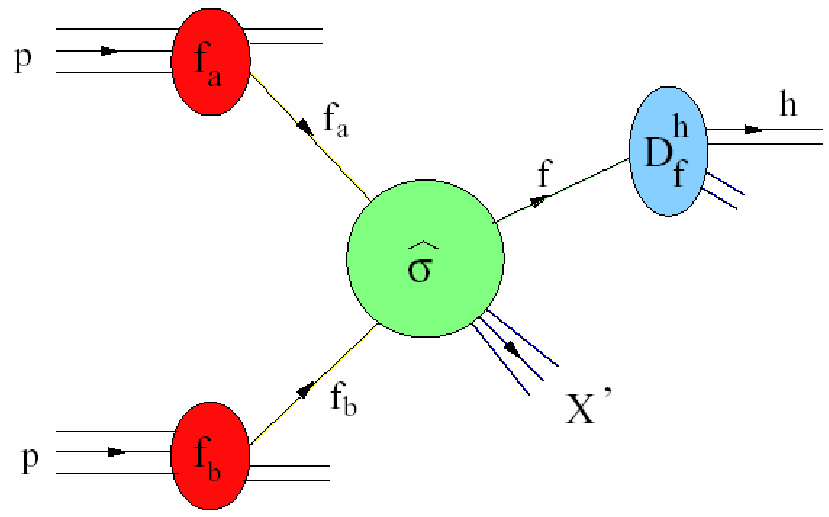


縦偏極の実験

- クォークスピンとグルーオンスピンの寄与の測定
  - QCD collinear factorization
  - Q^2 evolution
  - global analysis

QCD factorization

– ex. hadron production in proton collisions
  - pp→hX



$$d\sigma = \sum_{a,b,c} \int dx_a \int dx_b \int dz_c \; f_a(x_a,\mu) \; f_b(x_b,\mu) \; D_c^h(z_c,\mu) \; d\hat{\sigma}_{ab}^c(x_a P_A, x_b P_B, P_h / z_c, \mu)$$

$f_a(x_a,\mu)$, $f_b(x_b,\mu)$ parton distribution function (PDF)

$D_c^h(z_c,\mu)$ fragmentation function (FF)

} long distance term

$d\hat{\sigma}_{ab}^c(x_a P_A, x_b P_B, P_h / z_c, \mu)$ partonic cross section short distance term

μ factorization scale – boundary between short and long distance

Global QCD analysis

- framework to combine various experimental data into a systematically controlled extraction of the unpol. & pol. PDFs, FFs
  - experimental data $a^{\mathrm{data}}(x, Q^2)$ with experimental errors $\delta a^{\mathrm{data}}(x, Q^2)$
  - function form (parametrizations) of PDFs and FFs satisfying physical requirements at the initial Q^2_0
  - Q^2 evolution of PDFs/FFs and theoretical calculation corresponding to the experimental data $a^{\mathrm{calc}}(x, Q^2)$
  - χ^2 analysis (minimization)

$$\chi^2 = \sum_a \left(\frac{a^{\mathrm{data}}(x,Q^2) - a^{\mathrm{calc}}(x,Q^2)}{\delta a^{data}(x,Q^2)} \right)^2$$

  - parameters (and errors on the parameters) determined

グルーオンスピンの寄与

- 偏極深非弾性散乱実験のスケール則の破れ
  - 摂動QCDの発展方程式の重要な成功
  - Q^2の範囲が限られている

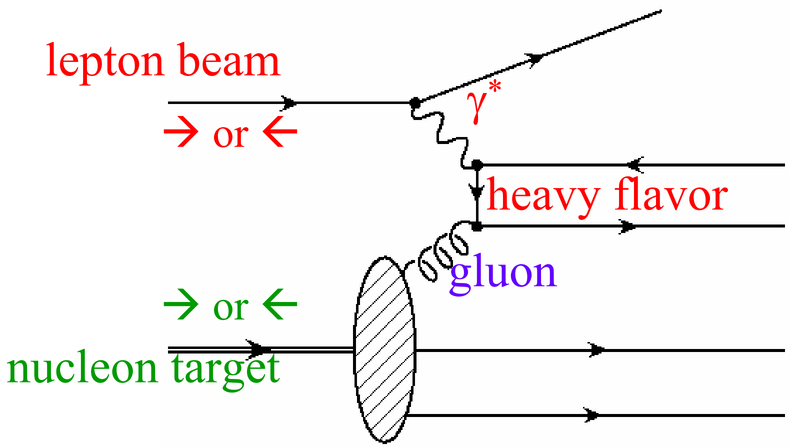
$$\text{SMC}:\ \Delta g(Q^2 = 1\,\text{GeV}^2) = 0.99^{+1.17}_{-0.31}(\text{stat})^{+0.42}_{-0.22}(\text{syst})^{+1.43}_{-0.45}(\text{th})$$

B. Adeba et al., PRD 58, 112002 (1998).

$$\text{E155}:\ \Delta g(Q^2 = 5\,\text{GeV}^2) = 1.6 \pm 0.8(\text{stat}) \pm 1.1(\text{syst})$$

P.L. Anthony et al., PLB 493, 19 (2000).

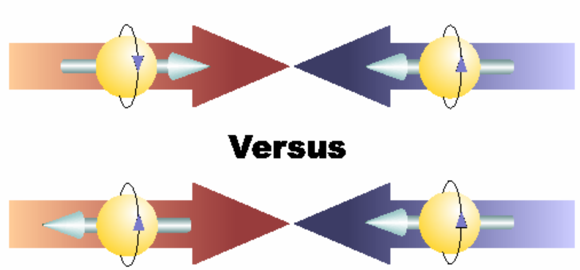
- semi-inclusive 深非弾性散乱実験
  - 高いp_Tのハドロン対生成
  - オープンチャーム生成



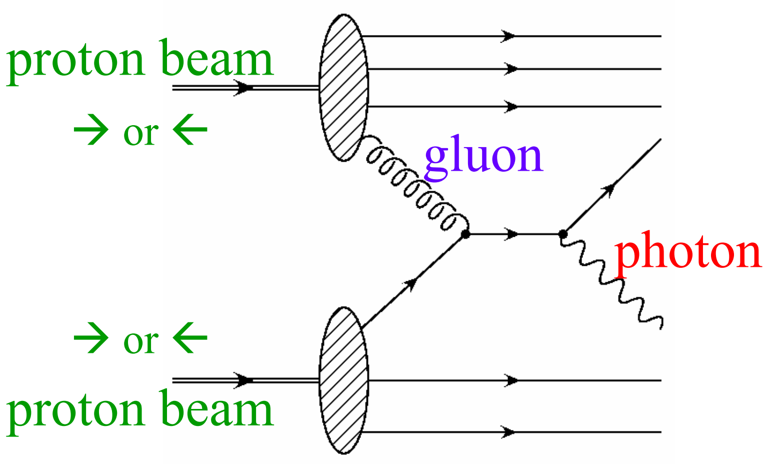
グルーオンスピンの寄与

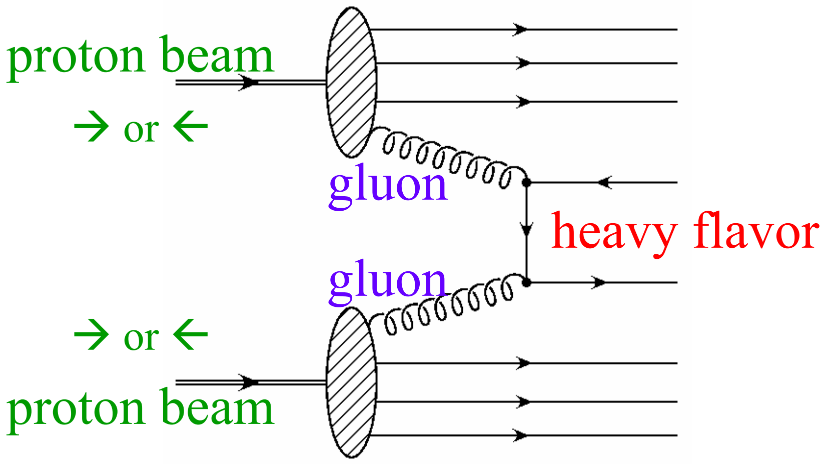
- 偏極ハドロン衝突実験
  - double helicity asymmetry

$$A_{LL} = \frac{d\sigma_{++} - d\sigma_{+-}}{d\sigma_{++} + d\sigma_{+-}}$$



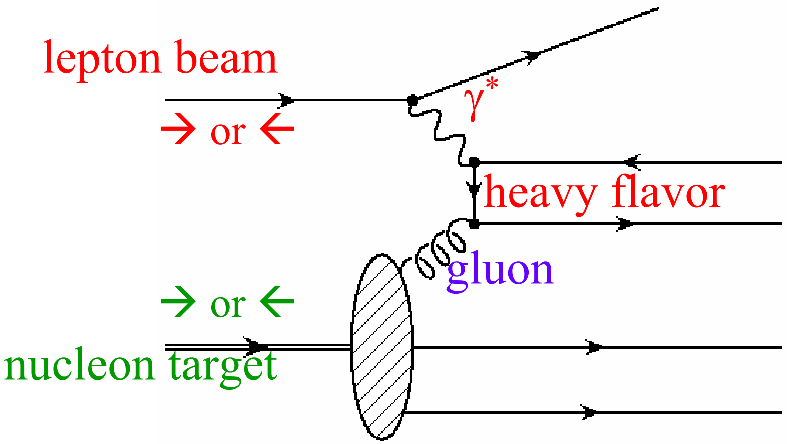
  - leading-order グルーオン測定
    - 光子の直接生成
    - 重いフレーバー生成

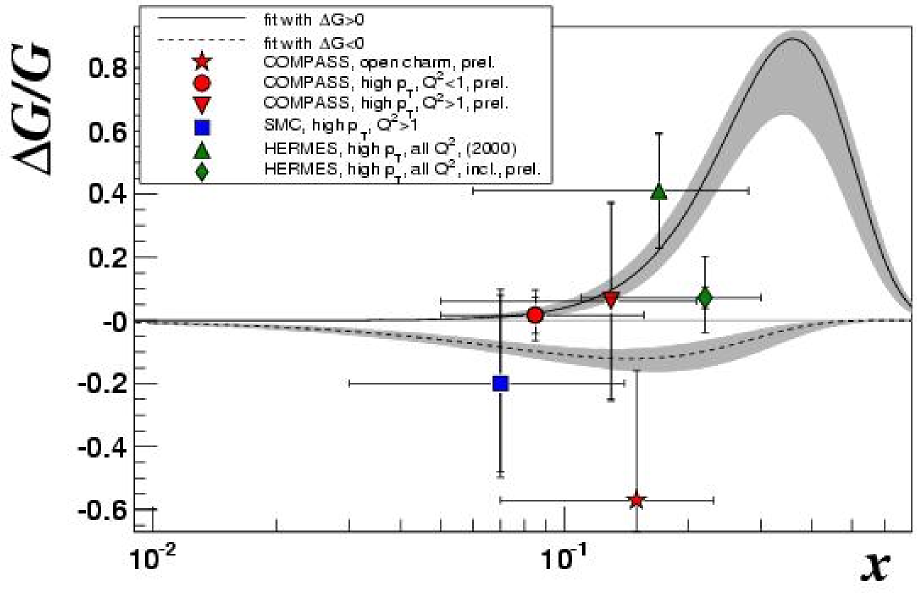


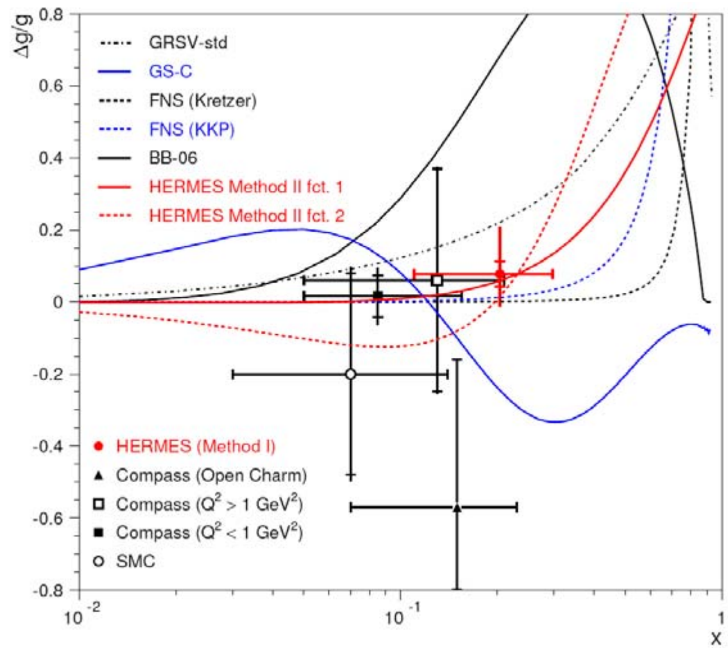


グルーオンスピンの寄与

- semi-inclusive
  深非弾性散乱実験
  - HERMES@DESY
    - 高いp_Tのハドロン対生成
  - SMC@CERN
    - 高いp_Tのハドロン対生成
  - COMPASS@CERN
    - 高いp_Tのハドロン対生成
    - オープンチャーム生成

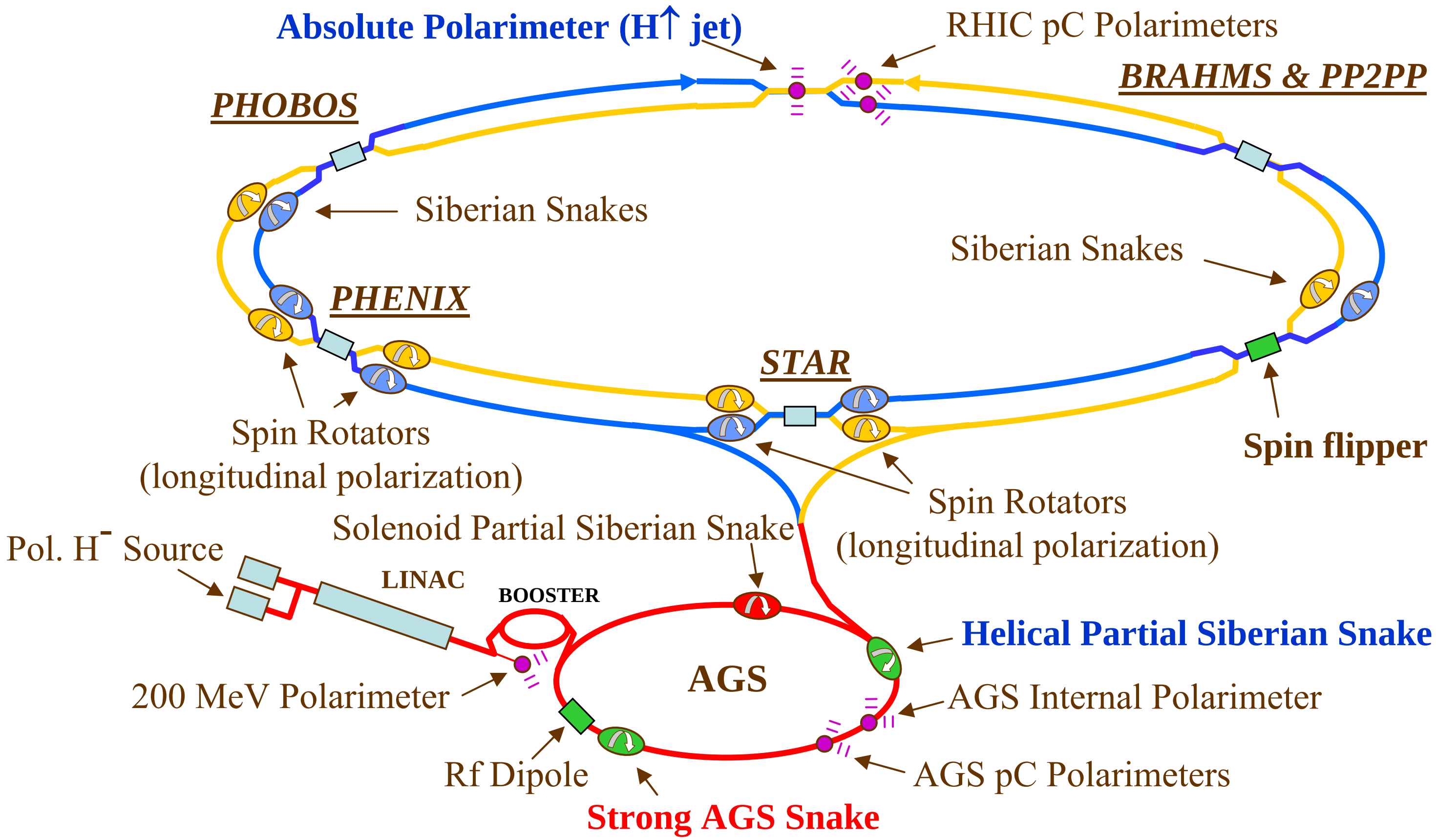






RHICスピン計画@BNL

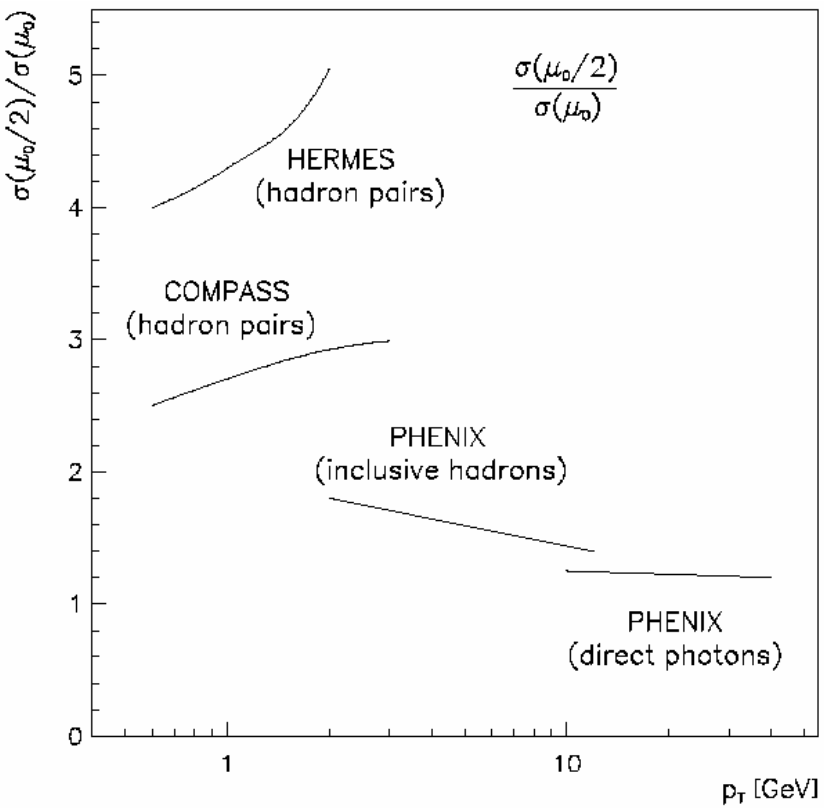
- 偏極陽子衝突型加速器
  - エネルギー 200 GeV (将来は 500 GeV)
  - 偏極度60%以上 (目標は70%)



生成断面積測定

- perturbative QCD applicable ?
  - dependence of the calculated cross section on μ represents an uncertainty in the theoretical predictions

$$\frac{\sigma(\mu_0/2)}{\sigma(\mu_0)}$$



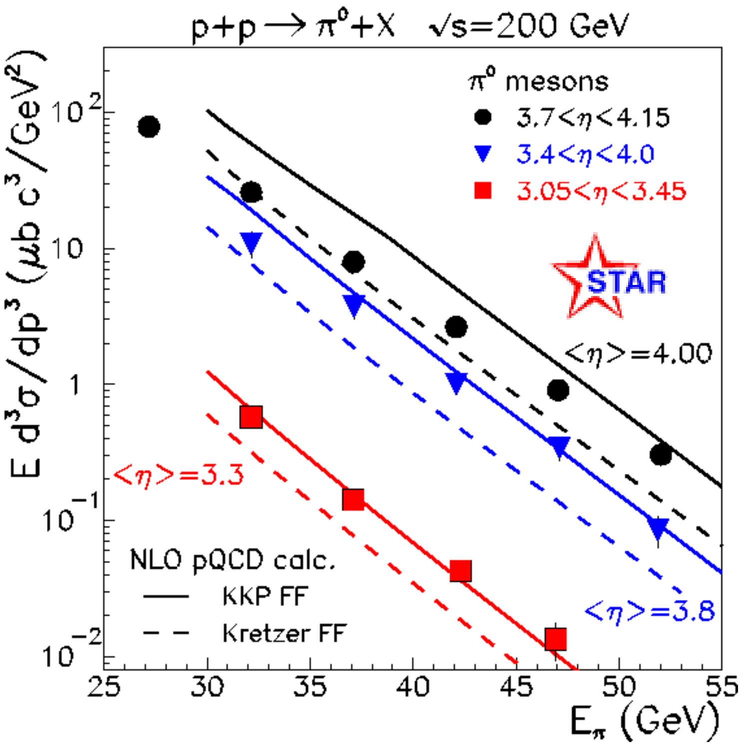
M. Stratmann and W. Vogelsang

p_T (GeV/c)

生成断面積測定

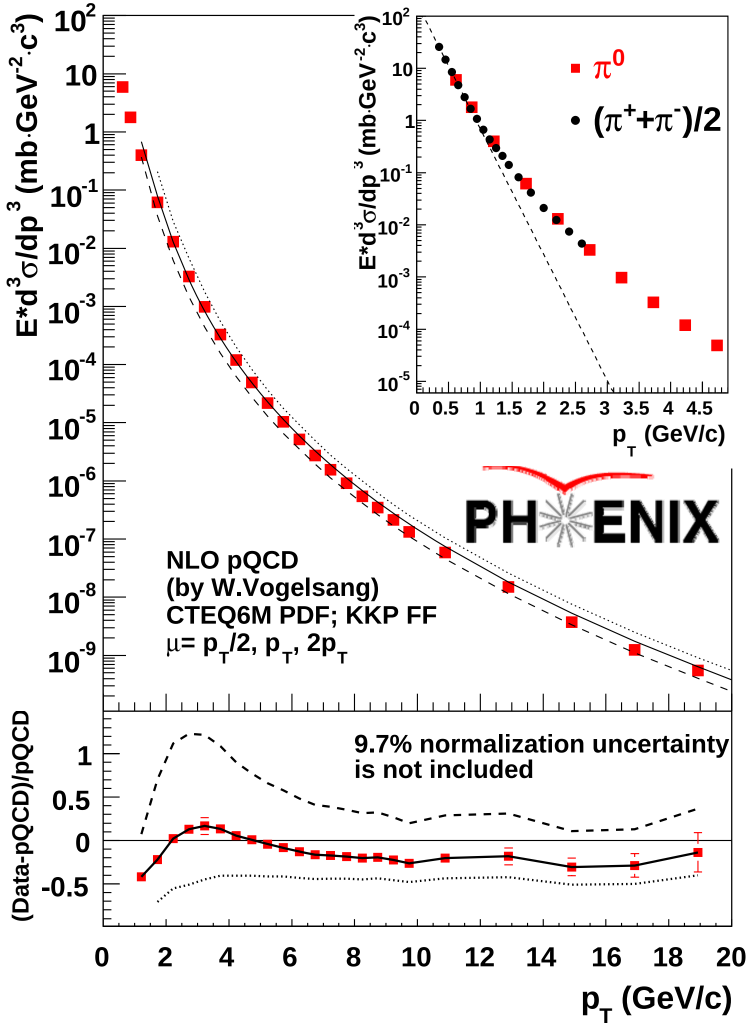
- comparison of π^0 cross section between data and NLO perturbative-QCD calculations
- agreement is excellent down even to $p_T \sim 1$ GeV/c

forward rapidity



$\sqrt{s} = 200$ GeV

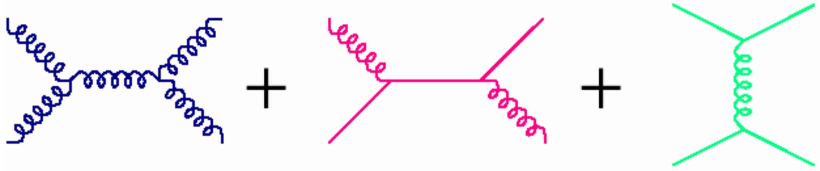
mid-rapidity

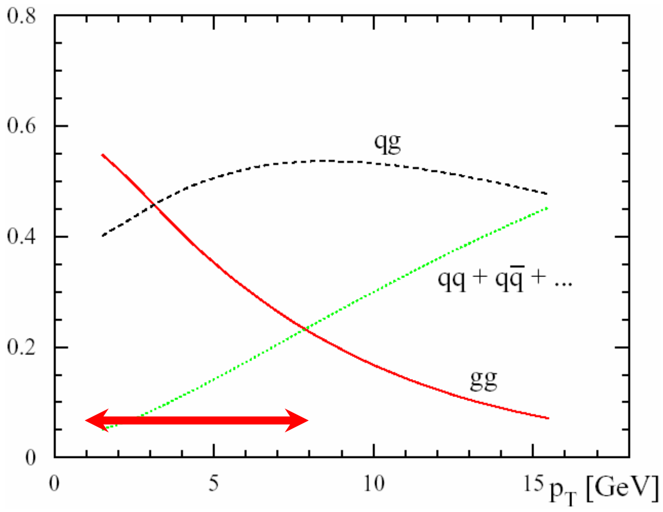


グルーオンスピンの寄与

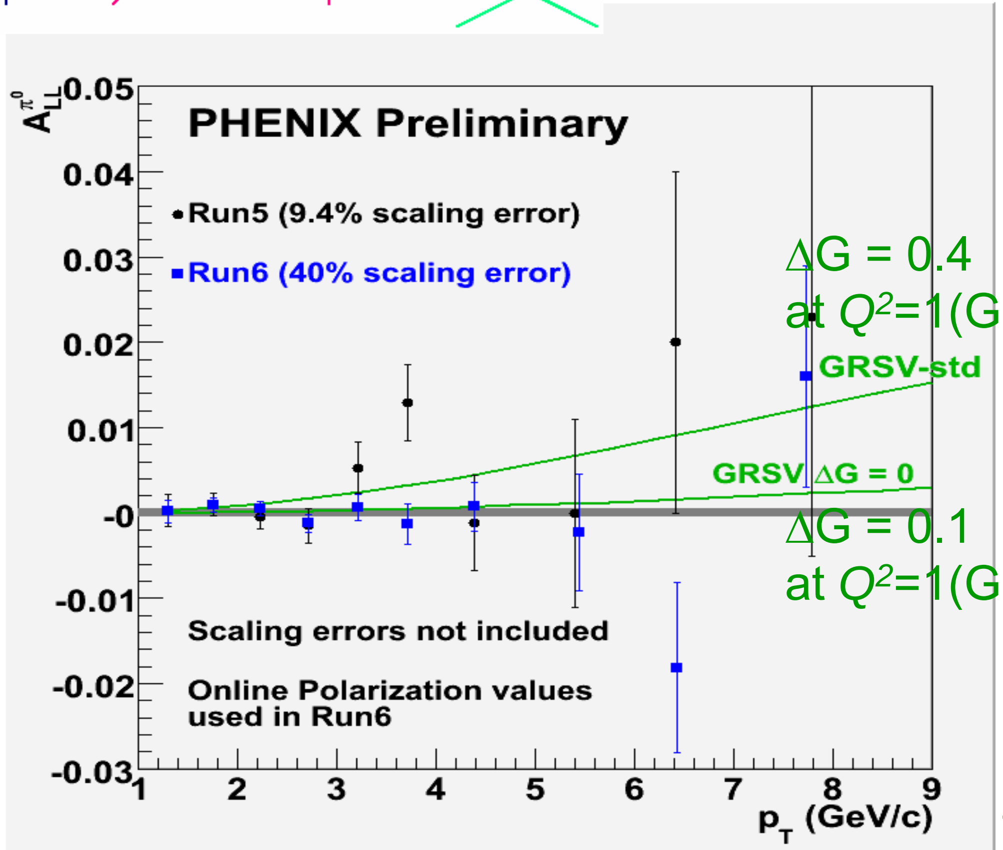
- PHENIX A_{LL} in neutral pion production
  - mid-rapidity |η| < 0.35, √s = 200 GeV

$$A_{LL} = [\omega_{gg}]\Delta g \Delta g + [\omega_{gq}\Delta q]\Delta g + [\omega_{qq}\Delta q \Delta q]$$





gg + qg dominant
sensitive to the gluon reaction



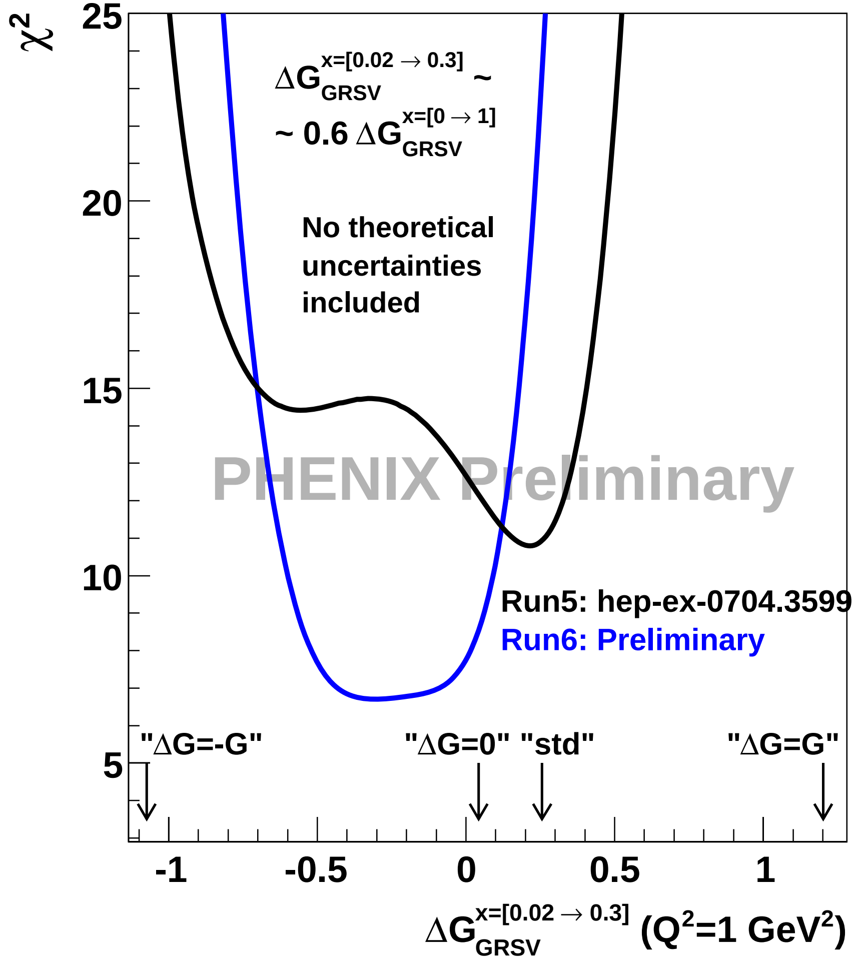
ΔG = 0.4 at Q^2=1(GeV/c)^2

ΔG = 0.1 at Q^2=1(GeV/c)^2

グルーオンスピンの寄与

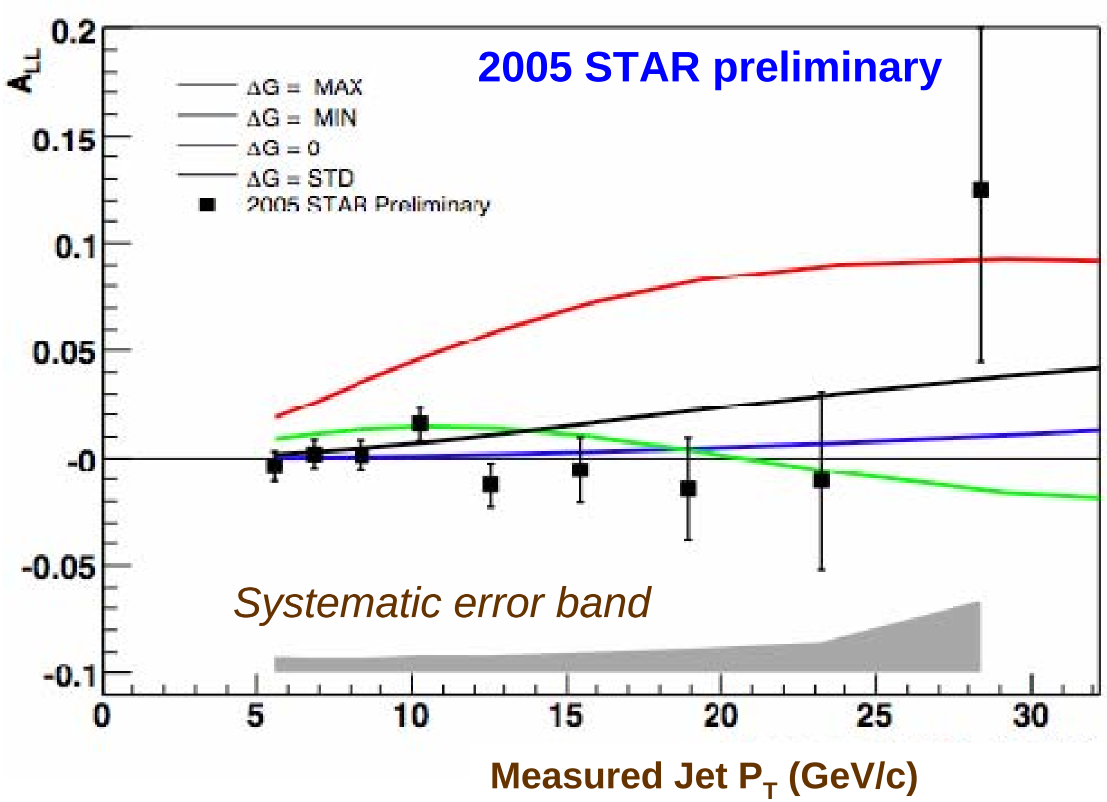
- PHENIX A_{LL} of π^0
  - GRSV-std scenario, ΔG = 0.4 at $Q^2 = 1(\mathrm{GeV}/c)^2$, excluded by data on more than 3-sigma level, $\chi^2(\mathrm{std}) - \chi^2_{min} > 9$
    - only experimental statistical uncertainties included (the effect of systematic uncertainties expected to be small in the final results)
    - theoretical uncertainties not included

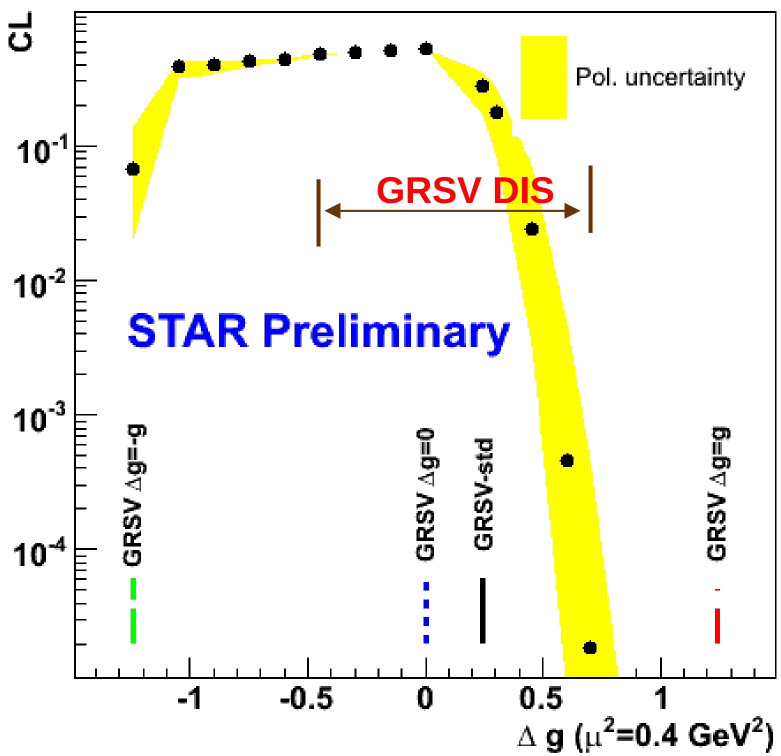
Calc. by W.Vogelsang and M.Stratmann



グルーオンスピンの寄与

- STAR A_{LL} of jet

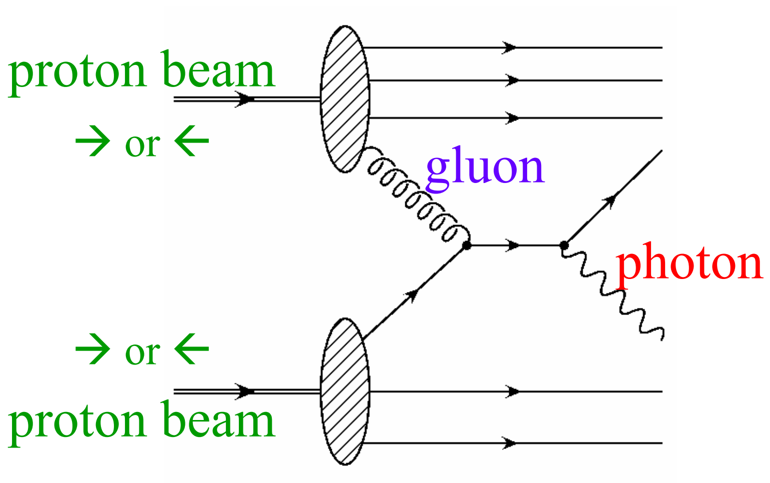


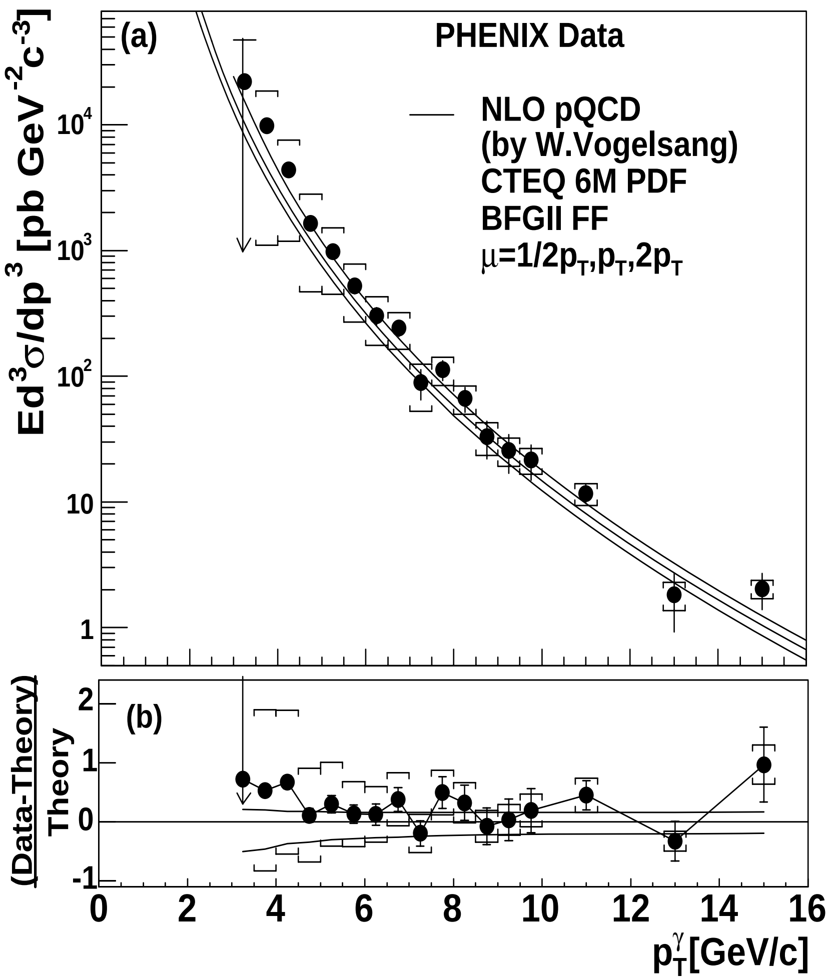


グルーオンスピンの寄与

- 直接生成光子
  - mid-rapidity |η| < 0.35, √s = 200 GeV
  - gluon compton process dominant ~75%

$$A_{LL}(p_T) = \frac{\Delta g(x_g)}{g(x_g)} \cdot A_1^p(x_q) \cdot \hat{a}_{LL}$$





クォーク偏極のフレーバー依存

- various quark and antiquark polarization individually

$$\Delta u, \Delta\bar{u}, \Delta d, \Delta\bar{d}, \Delta s, \Delta\bar{s}$$

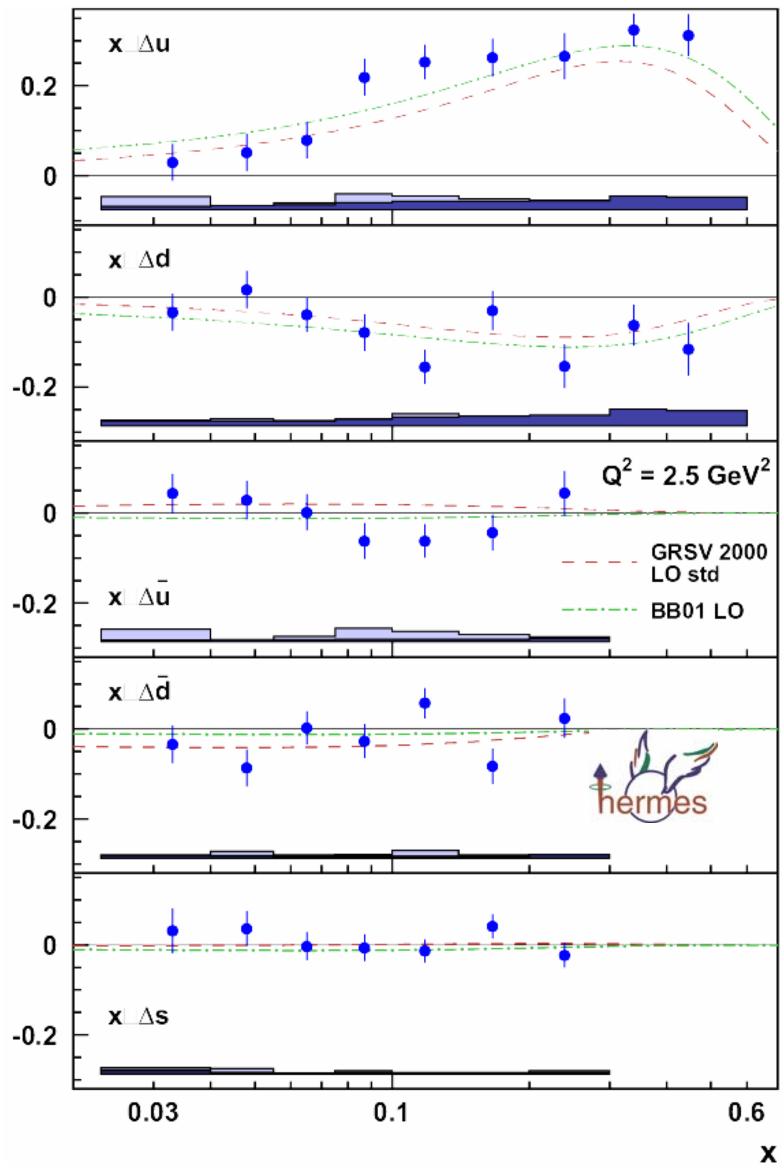
  - reduction of uncertainties to determine the quak spin contribution ΔΣ and gluon spin contribution ΔG to the proton spin

- semi-inclusive DIS
  - HERMES

$$h = \pi^{\pm}, K^{\pm} \quad A_1^h \sim \frac{\sum_i e_i^2 \Delta q_i(x) \int dz D_i^h(z)}{\sum_{i'} e_{i'}^2 q_{i'}(x) \int dz D_{i'}^h(z)}$$

$$= \sum_i P_q^h(x,z) \Delta q_i(x)$$

  - $P_q^h(x,z)$: purity
    - unpolarized quantity

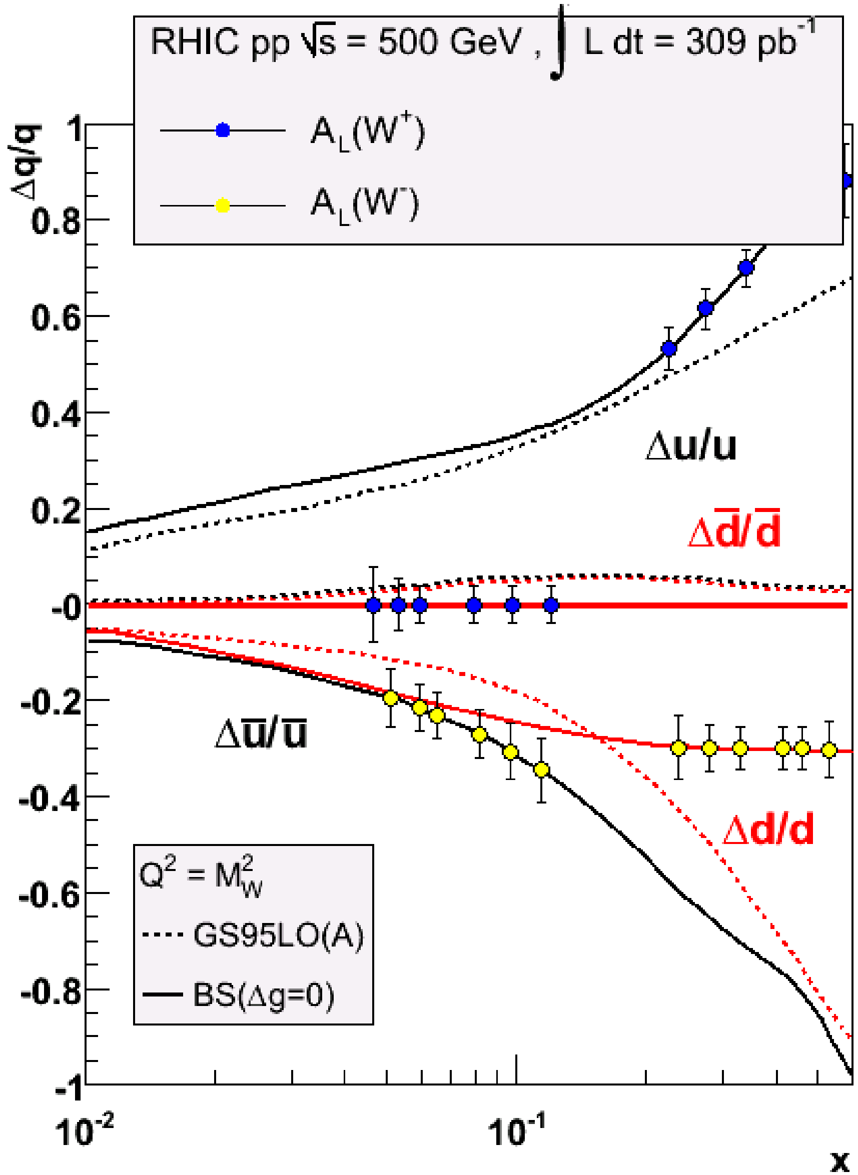


クォーク偏極のフレーバー依存

- weak boson production
  - RHIC spin
    - √s = 500 GeV
    - 2009 –
  - parity-violating asymmetry A_L

$$A_L^{W^+} = \frac{\Delta u(x_a)\bar{d}(x_b) - \Delta\bar{d}(x_a)u(x_b)}{u(x_a)\bar{d}(x_b) + \bar{d}(x_a)u(x_b)}$$

  - no fragmentation ambiguity
  - *x*-range limited



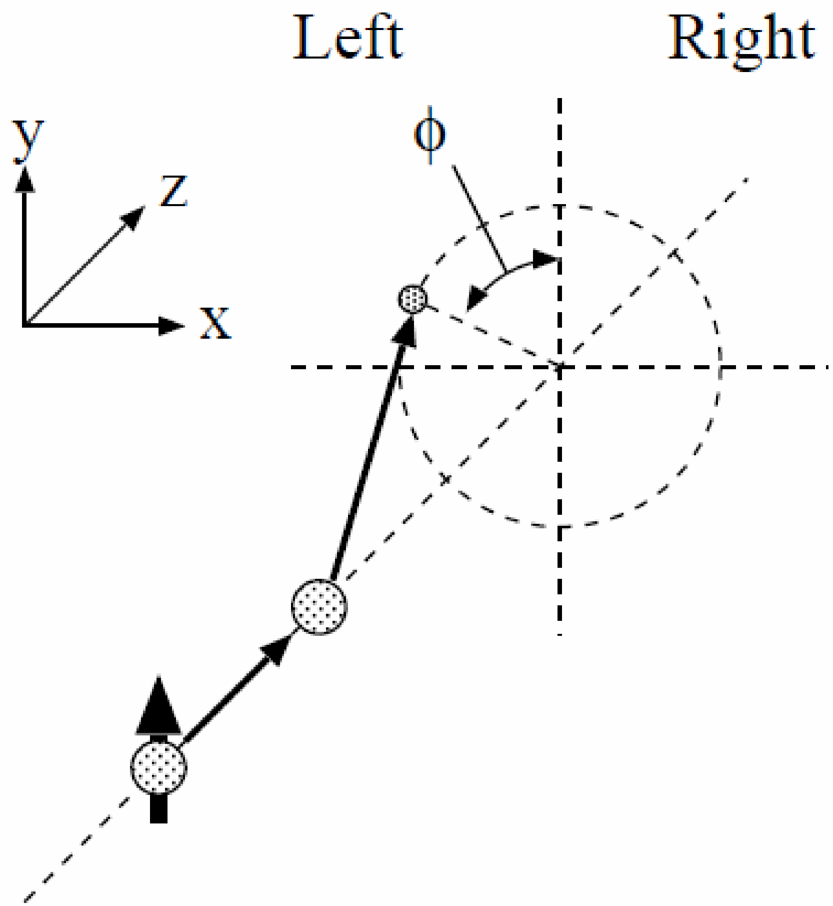
横偏極の実験

- クォークとグルーオンの軌道角運動量の寄与
  - QCDに基づく理論の開発
  - TMD (Transverse-Momentum Dependent) factorization など

Single transverse-spin asymmetry

- 左右非対称度

$$A_N = \frac{d\sigma_{Left} - d\sigma_{Right}}{d\sigma_{Left} + d\sigma_{Right}}$$

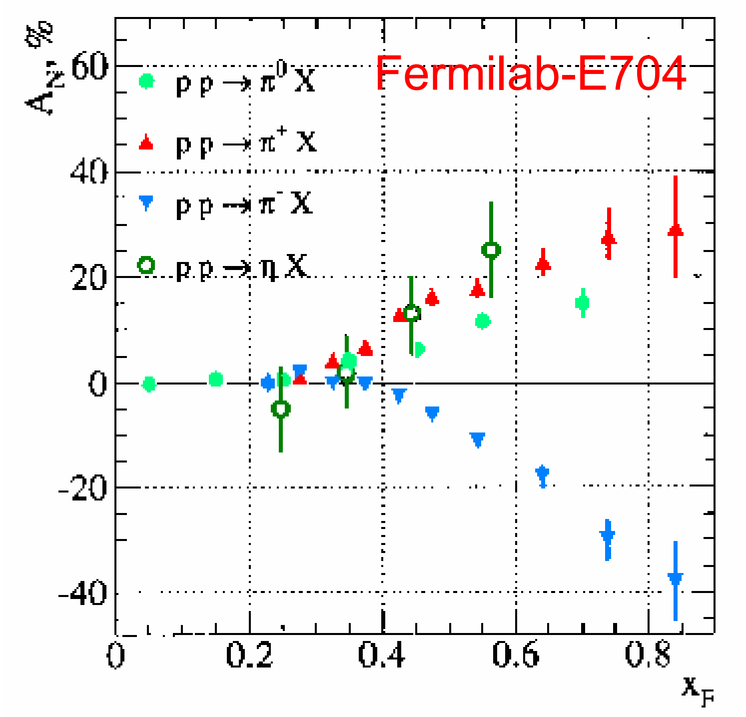


- 前方rapidity
  - Fermilab-E704
    - 固定標的実験
    - $\sqrt{s}$ = 19.4 GeV
  - RHIC-STAR
    - $\sqrt{s}$ = 200 GeV
    - x_F > 0.3 で大きな非対称度

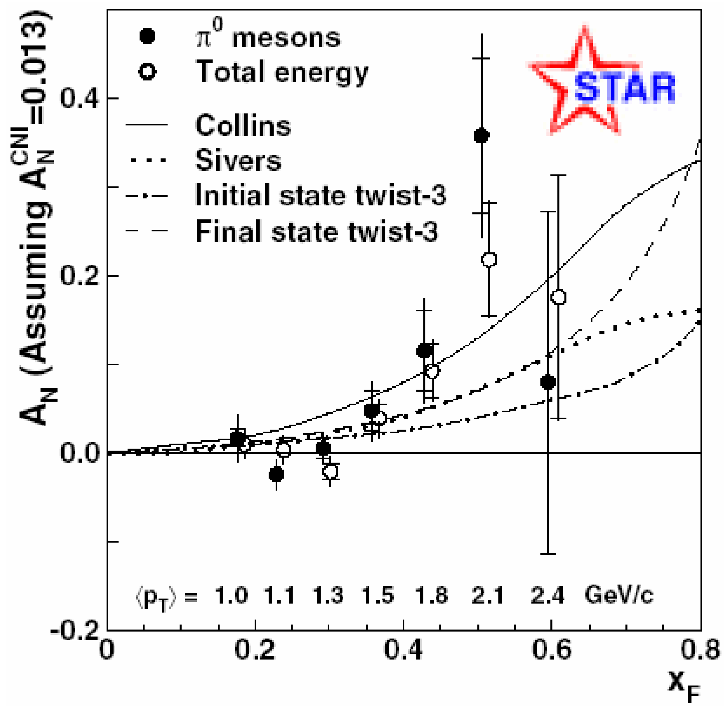
Single transverse-spin asymmetry

- 前方ラピディティー
  - 非対称度 ~20%
  - 多くのQCDに基づく理論の開発

√s = 19.4 GeV



√s = 200 GeV



Phys.Rev.Lett. 92 (2004) 171801

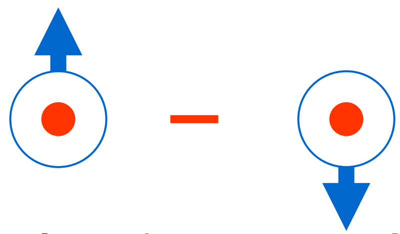
分布関数と破砕関数

- Transversity分布関数

$$\delta q(x) = h_{1T}(x)$$

  - 横方向に偏極した陽子内部におけるパートンの横方向スピンの分布
- Sivers分布関数

$$f_{1T}^{\perp}(x, p_T^2)$$

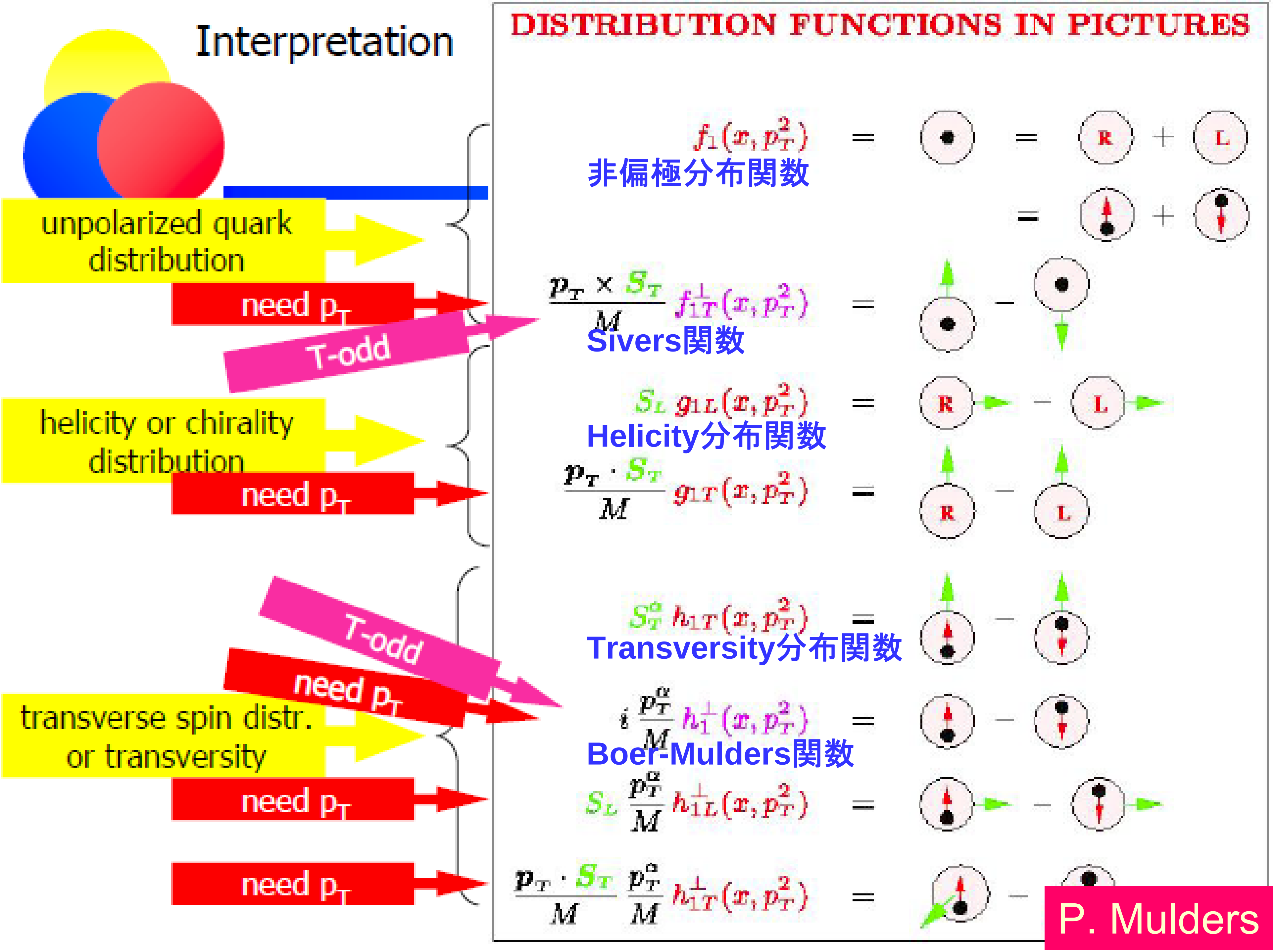


  - 陽子の横方向スピンと、陽子内部の非偏極パートンの横方向運動量(p_T^2)との相関
- Collins破砕関数

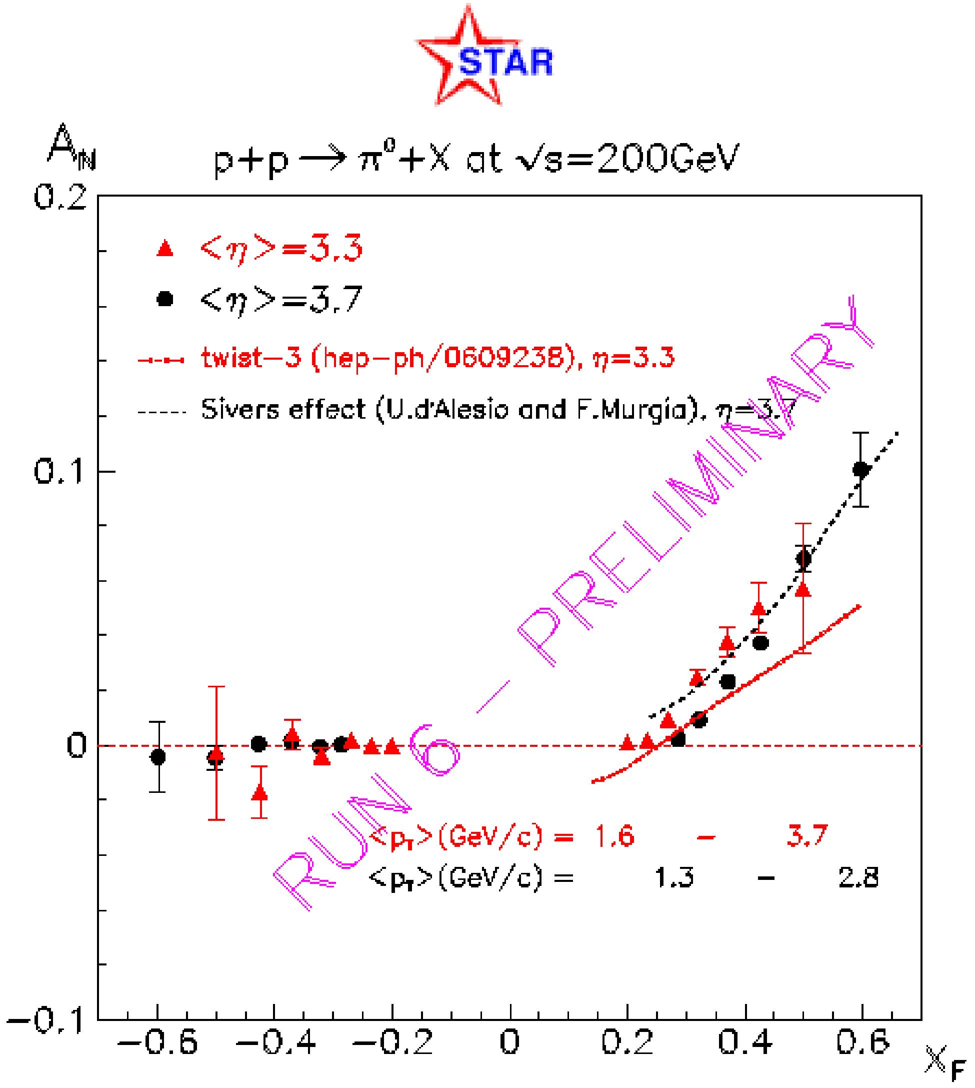
$$H_1^{\perp}(z, k_T^2)$$

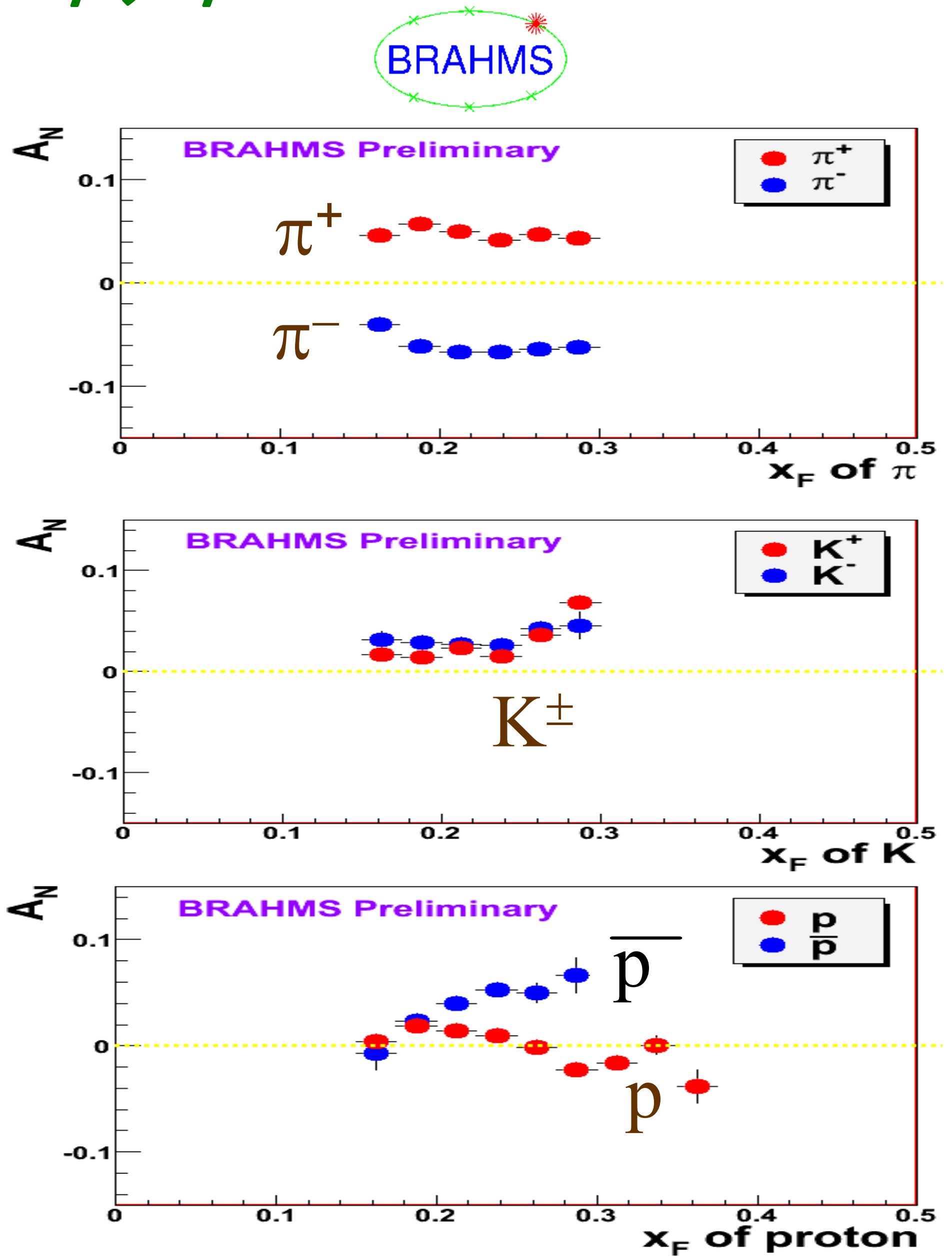
  - 破砕するパートンの横方向スピンと、生成されたハドロンのパートンに対する横方向運動量(k_T^2)との相関

Interpretation



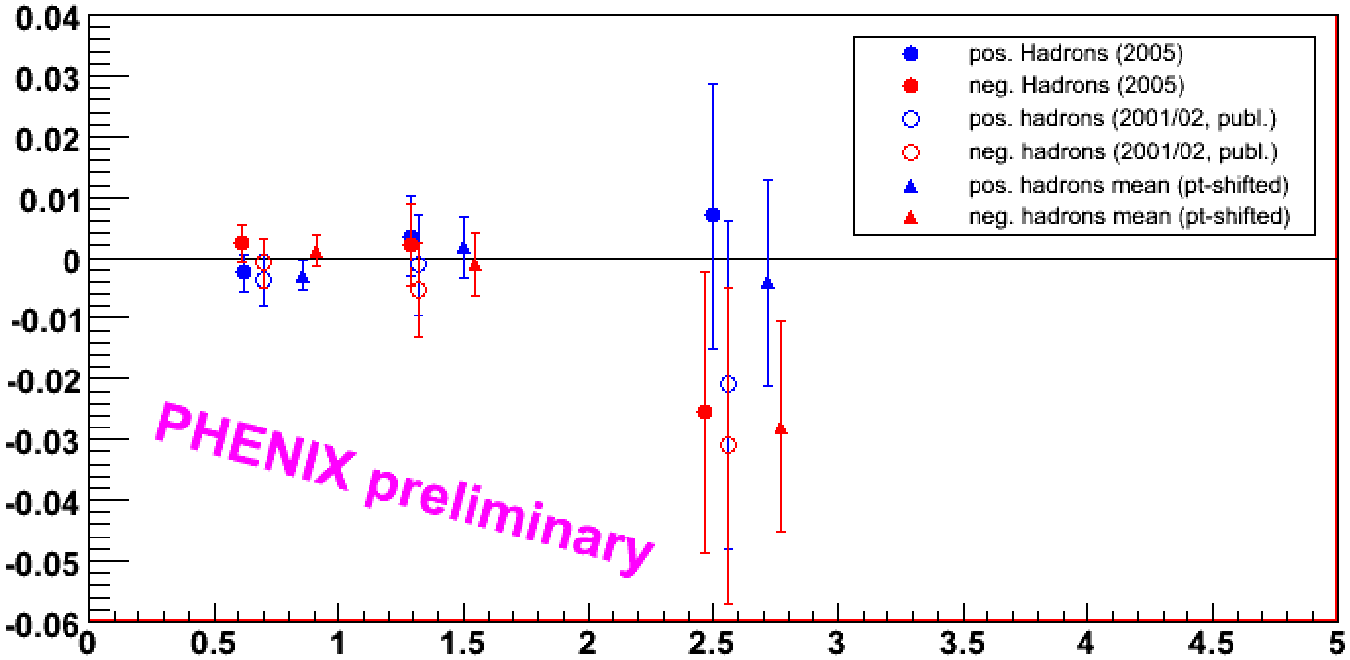
前方ラピディティー


STAR
A_N
p+p → π⁰+X at √s=200GeV
<η>=3.3
<η>=3.7
twist−3 (hep−ph/0609238), η=3.3
Sivers effect (U.d'Alesio and F.Murgia), η=3.7
RUN 6 – PRELIMINARY
<p_T>(GeV/c) = 1.6 – 3.7
<p_T>(GeV/c) = 1.3 – 2.8
0.2
0.1
0
−0.1
−0.6
−0.4
−0.2
0
0.2
0.4
0.6
x_F


BRAHMS
BRAHMS Preliminary
A_N
π⁺
π⁻
0.1
0
-0.1
0
0.1
0.2
0.3
0.4
0.5
x_F of π
BRAHMS Preliminary
A_N
K⁺
K⁻
K±
x_F of K
BRAHMS Preliminary
A_N
p
p̄
x_F of proton

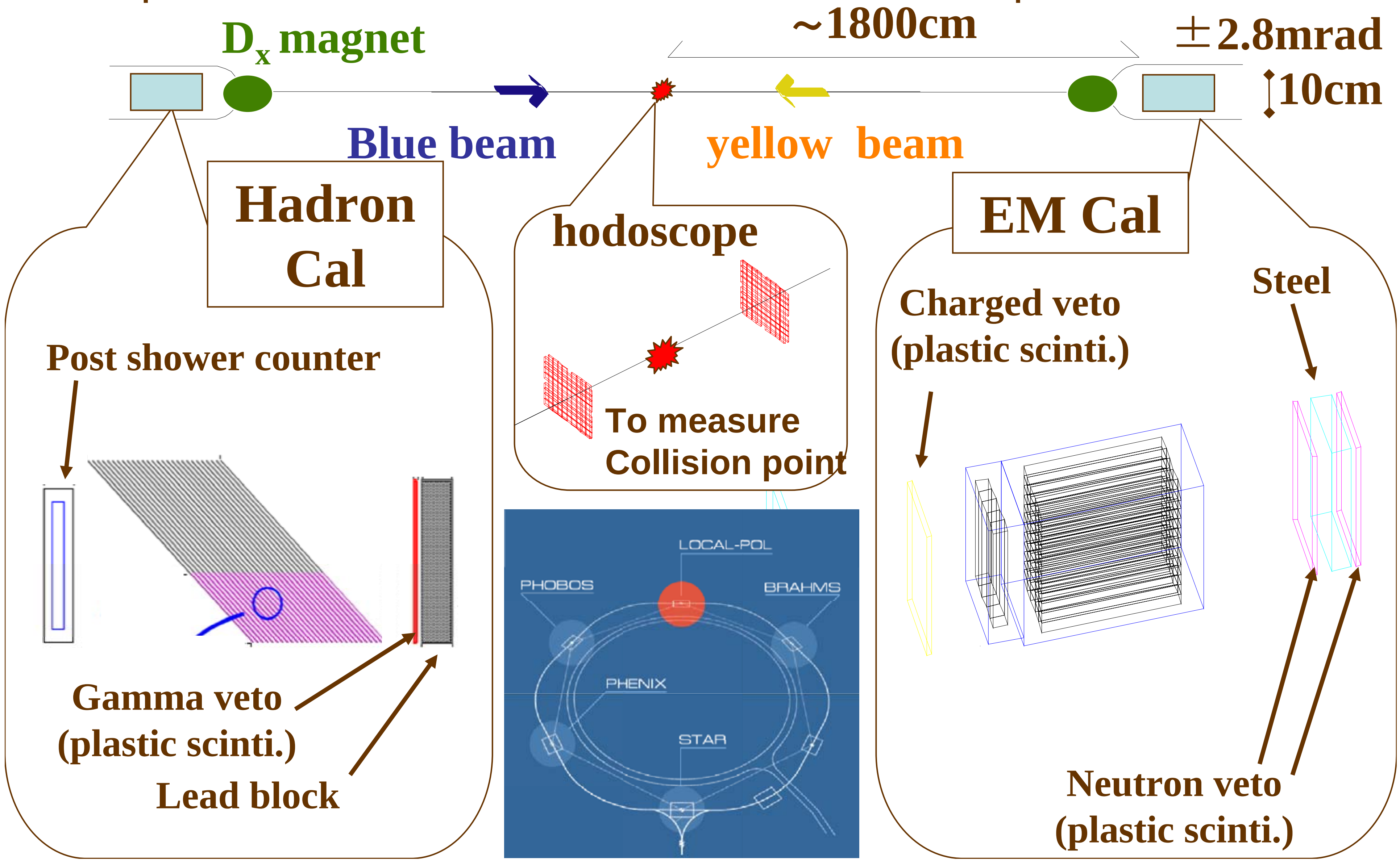
中央ラピディティー

- PHENIX
  - 異なる運動学的領域
    - contribution from both gluon-gluon and quark-gluon reactions
    - x = 0.03 – 0.1
    - small quark polarization/transversity
    - no gluon transversity in leading twist
    - negligible transversity & Collins effect contribution



中性子生成非対称度@IP12

- performed in 2001-2002 with √s = 200 GeV polarized proton collisions at the 12 o'clock collision point



中性子生成非対称度@IP12

- Phys. Lett. B 650 (2007) 325.

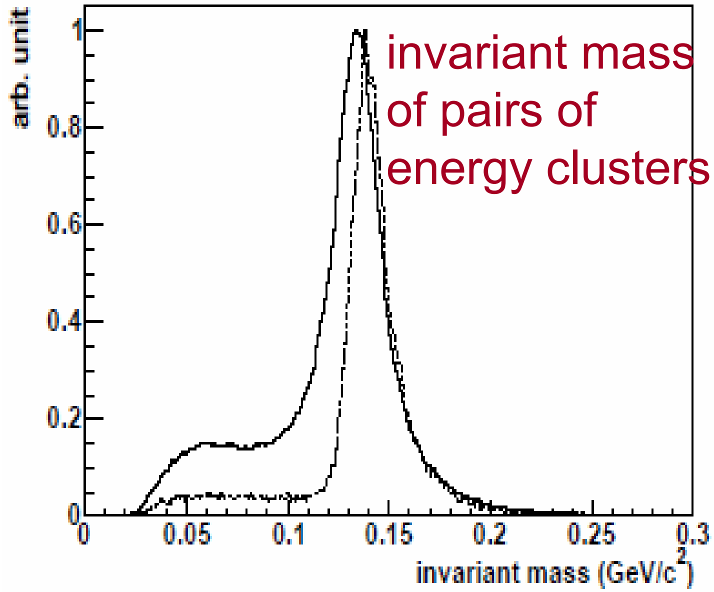
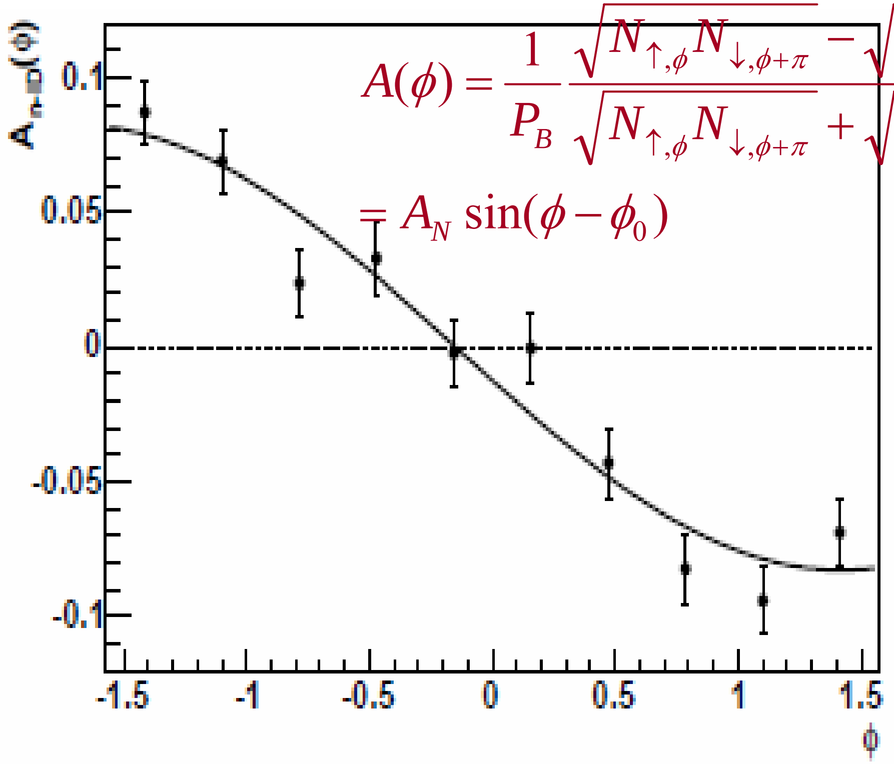


FIG. 3: Invariant mass of pairs of energy clusters in the EMCal, for 444K events with no additional selection requirements (solid) and for 35K events with photon identification (dashed).



$$A(\phi) = \frac{1}{P_B} \frac{\sqrt{N_{\uparrow,\phi} N_{\downarrow,\phi+\pi}} - \sqrt{N_{\uparrow,\phi+\pi} N_{\downarrow,\phi}}}{\sqrt{N_{\uparrow,\phi} N_{\downarrow,\phi+\pi}} + \sqrt{N_{\uparrow,\phi+\pi} N_{\downarrow,\phi}}}$$

$$= A_N \sin(\phi - \phi_0)$$

FIG. 4: Azimuthal dependence of asymmetry for the n-ID sample produced forward with respect to the polarized proton direction, based on the east detector. The error bars are statistical.

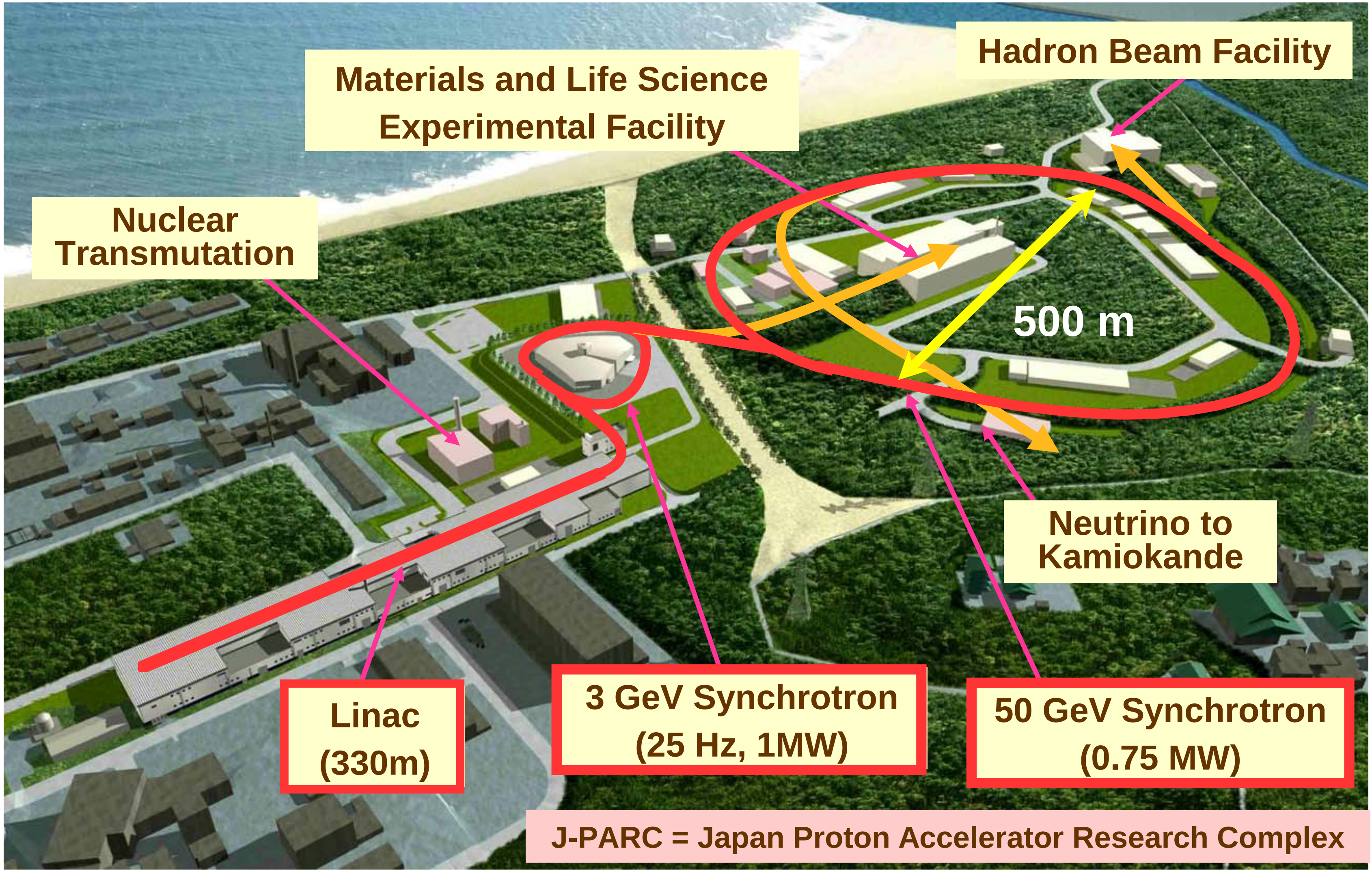
	forward	backward
neutron	$-0.090 \pm 0.006 \pm 0.009$	$0.003 \pm 0.004 \pm 0.003$
photon	$-0.009 \pm 0.015 \pm 0.007$	$-0.019 \pm 0.010 \pm 0.003$
π^0	$-0.022 \pm 0.030 \pm 0.002$	$0.007 \pm 0.021 \pm 0.001$

TABLE I: Asymmetries measured by the EMCal. The errors are statistical and systematic, respectively. There is an additional scale uncertainty, due to the beam polarization uncertainty, of $(1.0^{+0.47}_{-0.24})$.

イントロのまとめ

- RHICその他の偏極実験で、核子スピンに対するグルーオンスピンの寄与についての決着は着く
- 次は核子スピンに対する軌道角運動量の寄与の測定、決定
- 横偏極に対する非対称度について、わかっていないこともまだまだある

J-PARC facility



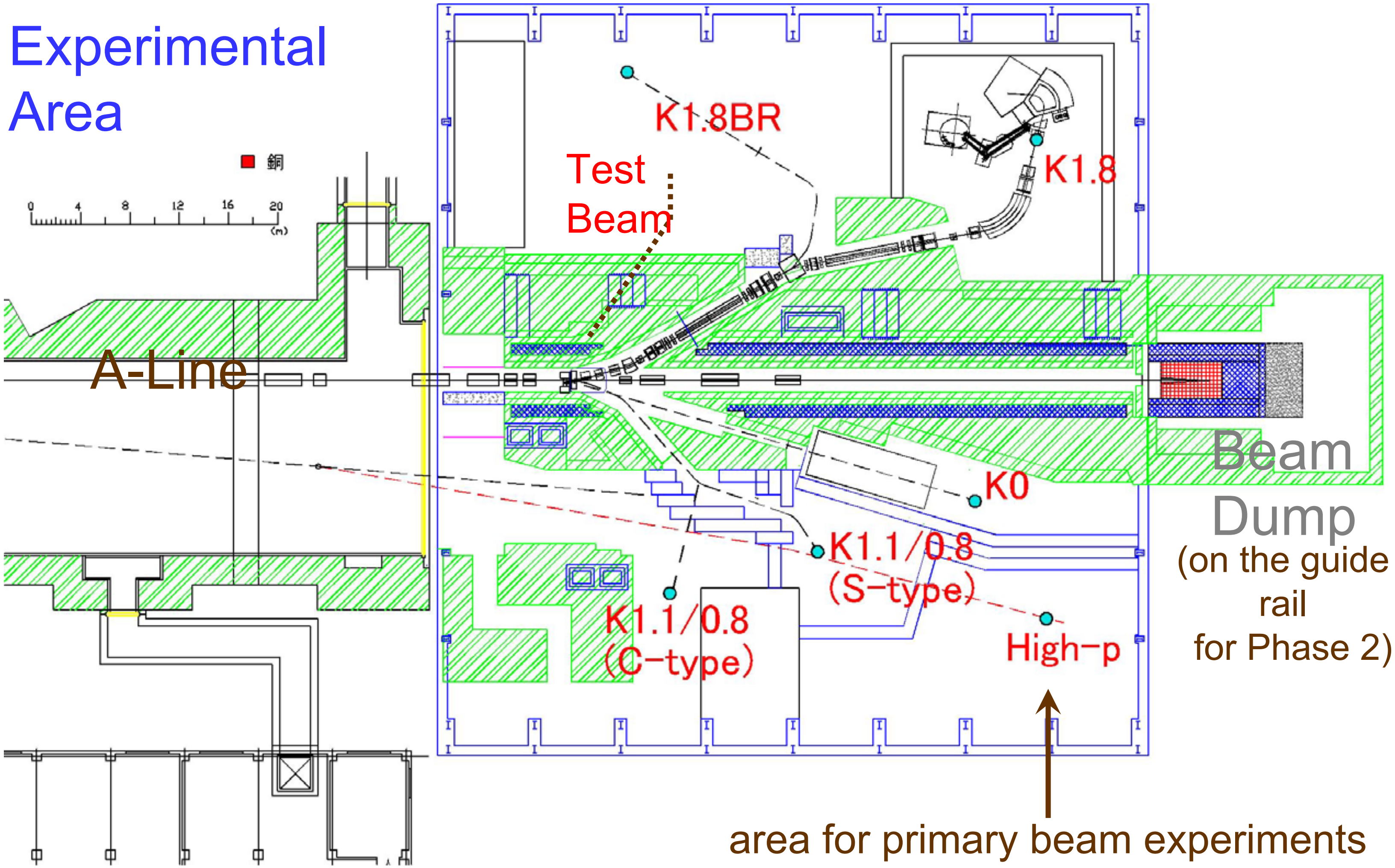
Joint Project between KEK and JAEA

J-PARC parameters

- 50 GeV beam
  - repetition 3.4 ~ 5 (or 6) sec
  - flat top width 0.7 ~ 2 (or 3) sec
  - linac energy 400 MeV
  - 3.3×10^{14} ppp, 15 μA
  - beam power 750 kW
- 30 GeV beam (phase-1)
  - linac energy 180 MeV
  - 2×10^{14} ppp, 9 μA
  - beam power 270 kW

Hadron experimental hall (phase 1)

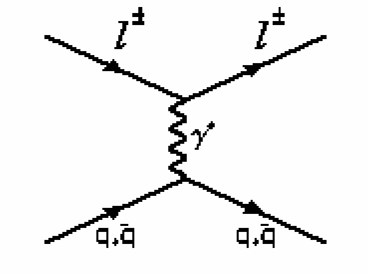
beamlines for secondary beam experiments at the beginning of the phase1



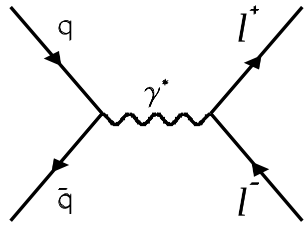
Drell-Yan実験

- The simplest process in hadron-hadron reactions

DIS



Drell-Yan



  - no final-state effect from QCD
  - no polarized Drell-Yan experiment yet
- Flavor asymmetry of the sea quark distribution
  - unpolarized and (longitudinally-)polarized
- Orbital angular momentum in the proton?
  - Sivers effect (no Collins effect)
- Transversity, etc.

Drell-Yan実験

- Why at J-PARC ?
  - polarized beam
    - feasible in discussions with J-PARC and BNL accelerator physicists
  - high intensity/luminosity
    - small Drell-Yan cross section
    - e.g. 10^{12} protons/sec, 10% target $\rightarrow$ $L = 2\times10^{36}/\mathrm{cm}^2/\mathrm{sec}$

Sea-quark分布のflavor非対称性

- Gottfried sum rule
  - in the parton model with isospin symmetry

$$I_G = \int_0^1 \frac{dx}{x}[F_2^p(x,Q^2) - F_2^n(x,Q^2)] = \frac{1}{3} + \frac{2}{3}\int_0^1 dx[\bar{u}(x,Q^2) - \bar{d}(x,Q^2)]$$

  - CERN NMC experiment

$$\int_{0.004}^{0.8} dx[F_2^p(x,Q^2) - F_2^n(x,Q^2)]$$

$$= 0.221 \pm 0.008(\text{stat.}) \pm 0.019(\text{syst.})$$

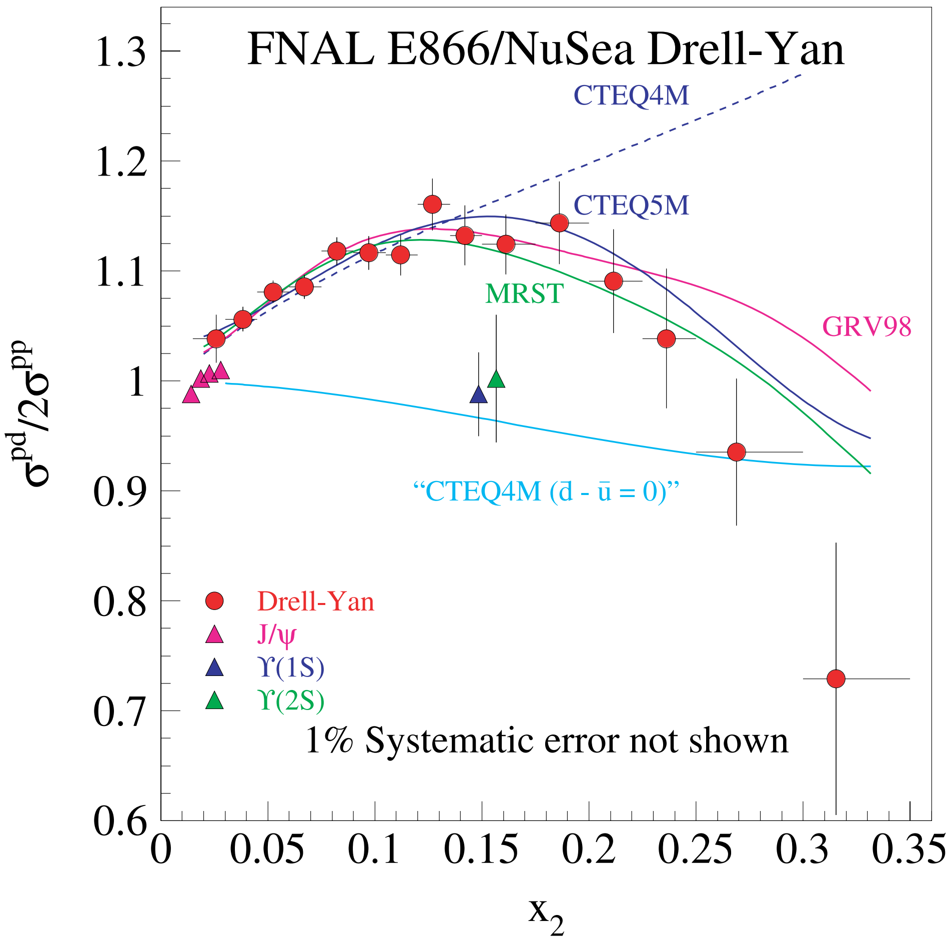
$$\text{at } Q^2 = 4\,\text{GeV}^2$$

$$I_G = 0.235 \pm 0.026$$

$$\int_0^1 dx[\bar{d}(x) - \bar{u}(x)] = 0.148 \pm 0.039$$

Sea-quark分布のflavor非対称性

- Fermilab E866 Drell-Yan experiment
  - flavor asymmetry of the seq quark distribution



$$\frac{\sigma^{pd}}{2\sigma^{pp}} \sim \frac{1}{2}\left[1+\frac{\bar{d}(x_2)}{\bar{u}(x_2)}\right]$$

with CTEQ5M

$$\int_{0.015}^{0.35} dx[\bar{d}(x)-\bar{u}(x)]$$

$$= 0.0803 \pm 0.011$$

$$\int_{0}^{1} dx[\bar{d}(x)-\bar{u}(x)]$$

$$= 0.118 \pm 0.012$$

Sea-quark分布のflavor非対称性

- Physics origin
  - meson cloud model
    - virtual meson-baryon state

$$p \to p\pi^0, n\pi^+, \Delta\pi$$

  - chiral quark model

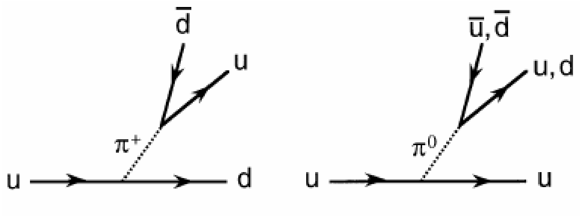


Fig. 17. Valence u quark splitting.

$\bar{d}$ distribution is softer than meson cloud model because it splits from (valence) quark

  - instanton
  - chiral-quark soliton model

Sea-quark分布のflavor非対称性

- π^+ は陽子中の余分なdbar の起源だろうか？

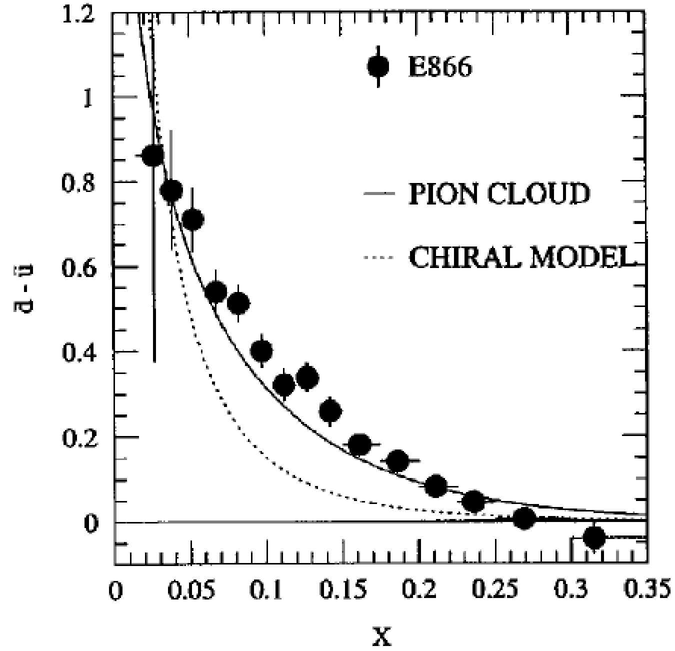


Figure 14: Comparison of the E866 [70] $\bar{d}-\bar{u}$ results at $Q^2 = 54\ \mathrm{GeV}^2/c^2$ with the predictions of pion-cloud and chiral models as described in the text.

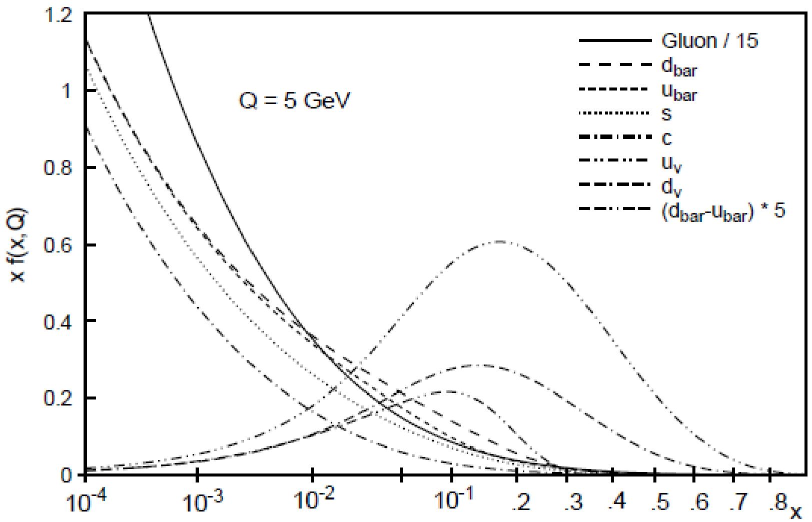


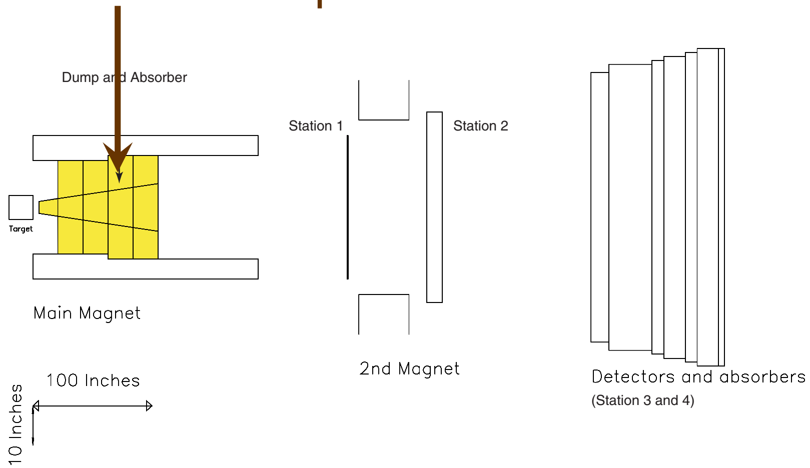
Fig. 2. Overview of CTEQ5M parton distributions at $Q = 5$ GeV. The gluon distribution is scaled down by a factor of 15, and the $(\bar{d}-\bar{u})$ distribution is scaled up by a factor of 5

J-PARCでのdimuon実験

– based on the Fermilab spectrometer for 800 GeV, the length can be reduced but the aperture has to be increased
– two vertically bending magnets with p_T kick of 2.47 GeV/c and 0.5 GeV/c
– tracking is provided by three stations of MWPC and drift chambers
– muon id and tracking are provided
– 2×10^{12} 50 GeV protons/spill

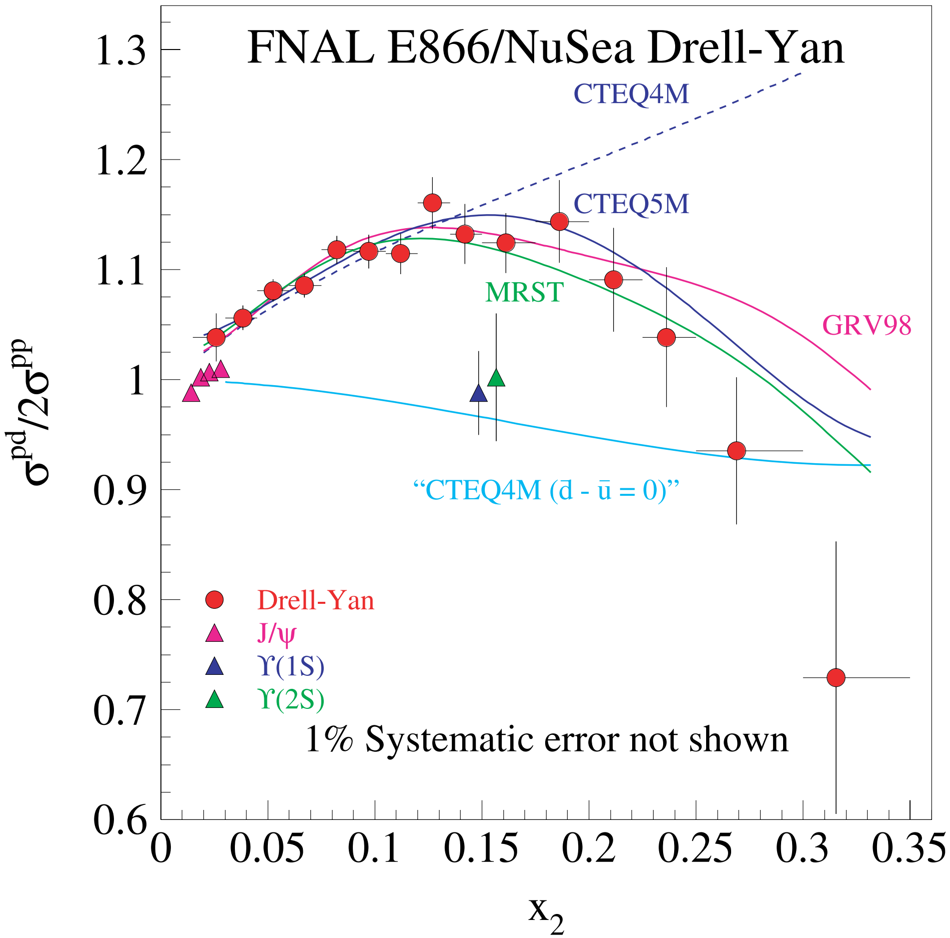
Schematic view in horizontal plane

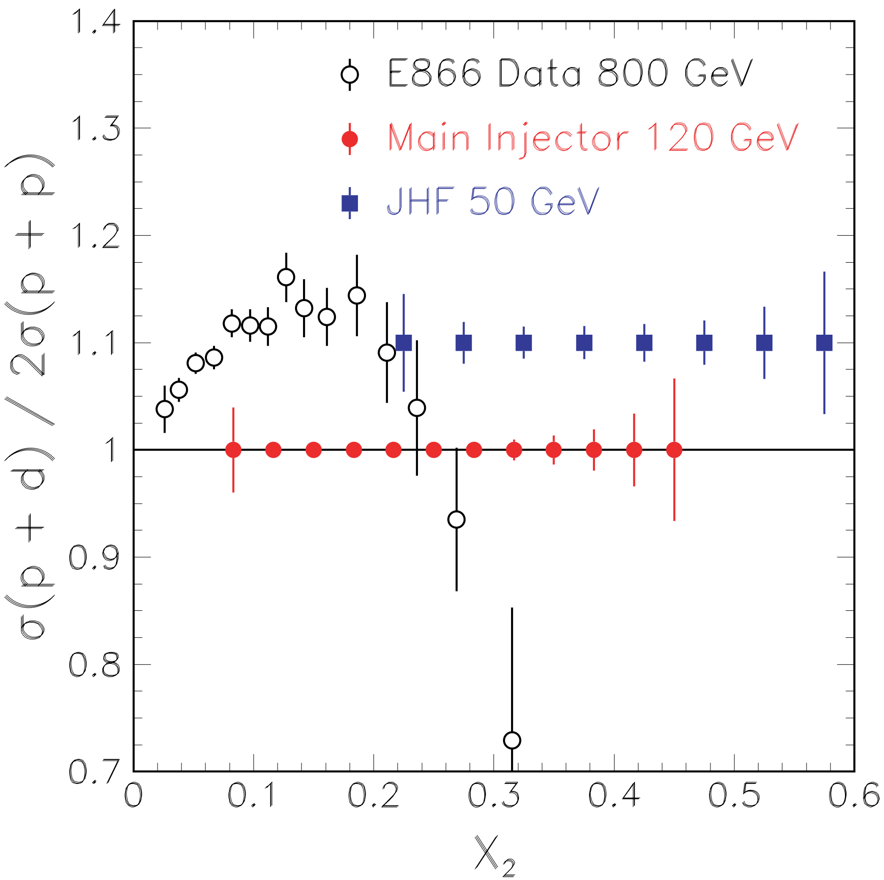
tapered copper beam dump and
Cu/C absorbers placed within the first magnet



J-PARCでのdimuon実験

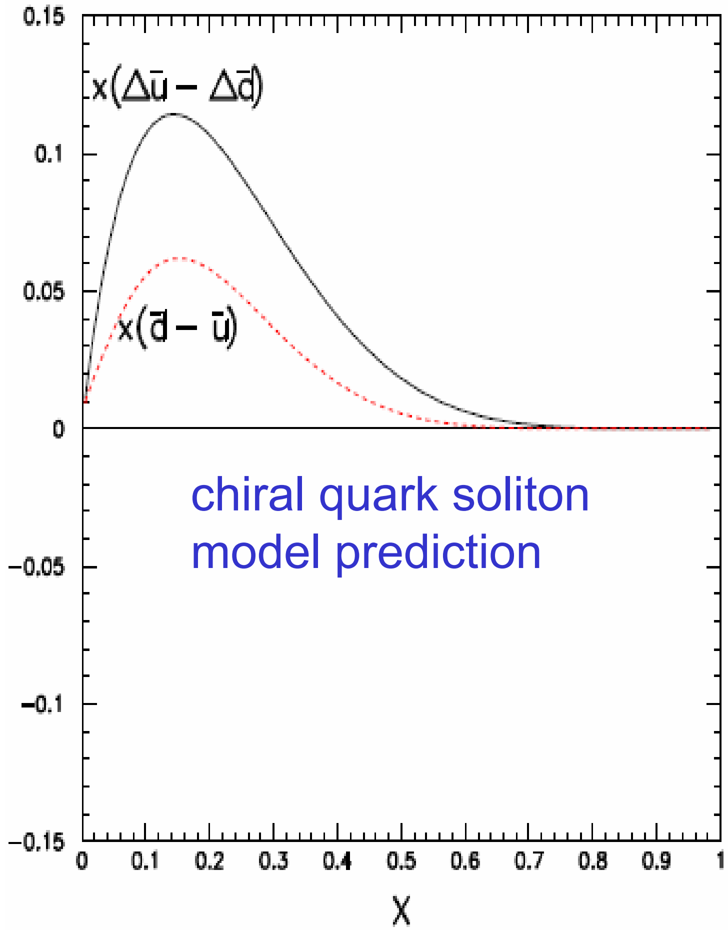
- Unpolarized program
  - proton beam on proton and deuterium target

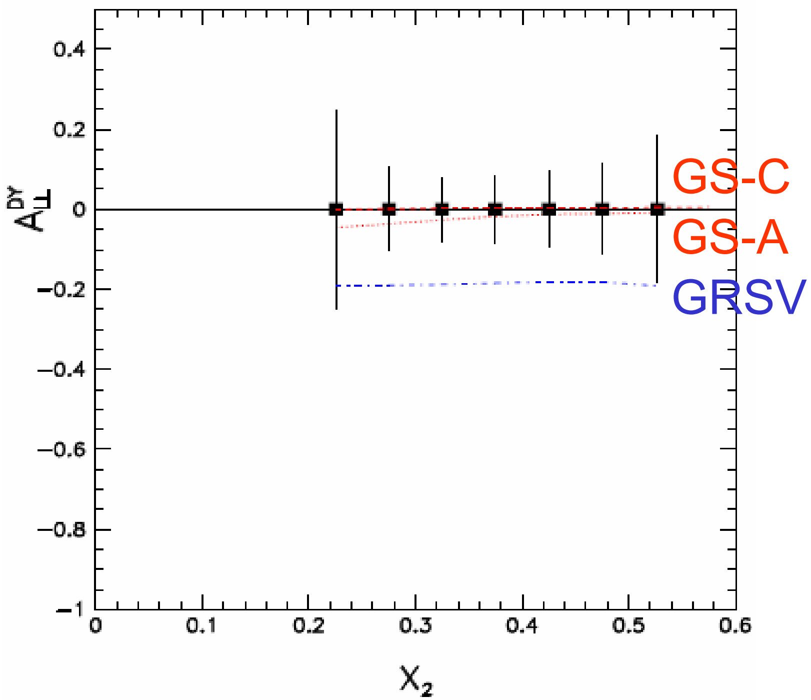




Spin physics at J-PARC

- Polarized Drell-Yan
  - A_{LL} measurement
    - flavor asymmetry of sea-quark polarization

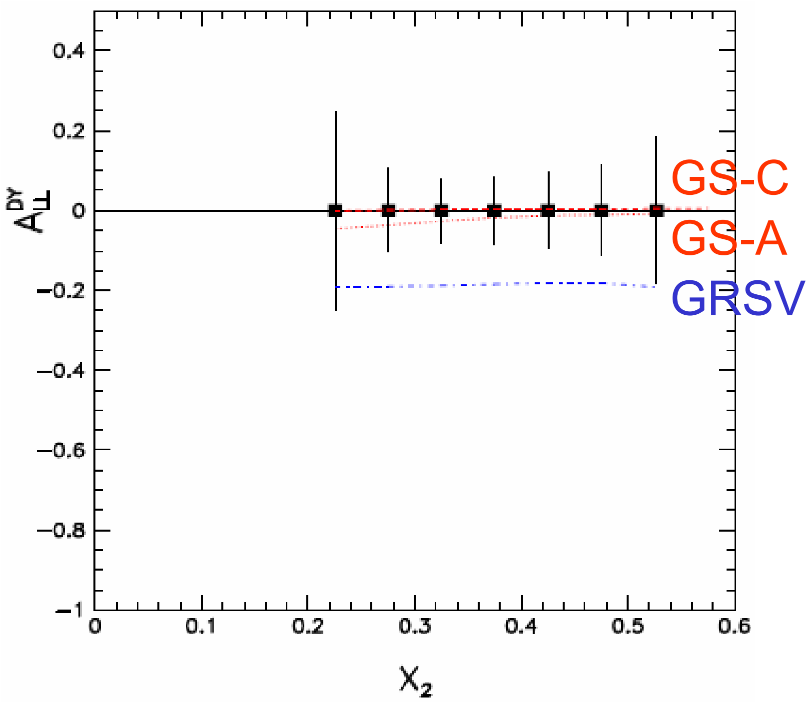


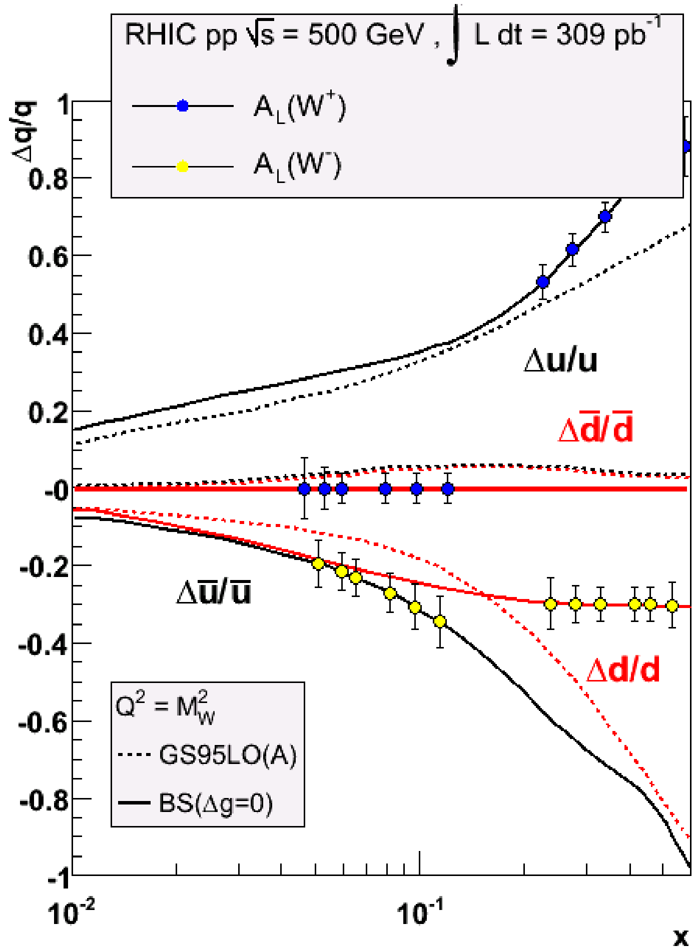


120-day run
75% polarization for a 5×10^{11} protons/spill
polarized solid NH3 target, 75% hydrogen
polarization and 0.15 dilution factor

Sea quark helicity distribution

- Polarized Drell-Yan
  - x: 0.25 – 0.5
- W production at RHIC
  - x: 0.05 – 0.1





reduction of uncertainties to determine the quak spin contribution ΔΣ and gluon spin contribution ΔG to the proton spin

Spin physics at J-PARC

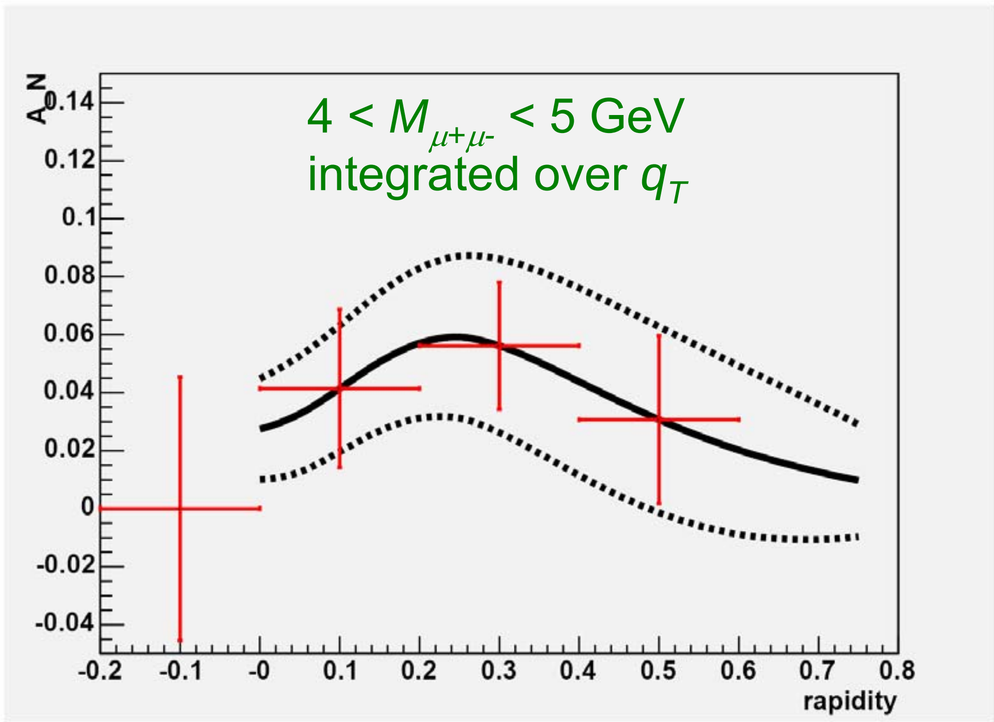
- 核子スピンに対する軌道角運動量の寄与
  - ハドロン衝突実験から直接軌道角運動量の寄与に結びつく理論はまだない
  - しかし、偏極核子中の横方向の運動に関わる効果は軌道角運動量と関係するはず
    - Sivers 効果
    - higher-twist 効果
  - A_{LL} 測定も high-x_F で軌道角運動量に sensitive…
    - Feng Yuan らの $L_z = 1$ に対する計算… [Harut Avakian, Feng Yuan, et al. arXiv:0705.1553.]

Polarized Drell-Yan

- A_N measurement
  - Ji, Qiu, Vogelsang, and Yuan
    - PRD 73, 094017 (2006)
    - sensitive to higher-twist effect at high $q_T \sim Q$
    - Sivers effect and higher-twist effect provide the same description of SSA on Drell-Yan at moderate q_T: $\Lambda_{QCD} << q_T << Q$
    - similar for SSA on semi-inclusive DIS: hep-ph/0604128
  - Sivers function in Drell-Yan is expected to have a sign opposite to that in DIS
    - test of QCD between e+p data and p+p data

Polarized Drell-Yan

- A_N measurement



Theory calculation by Ji, Qiu, Vogelsang and Yuan based on Sivers function fit of HERMES data (Vogelsang and Yuan: PRD 72, 054028 (2005))

1000 fb^{-1} (120-day run), 75% polarization, no dilution factor

Polarized Drell-Yan

- A_{TT} measurement
  - $h_1(x)$: transversity

$$A_{TT} = \hat{a}_{TT} \cdot \frac{\sum_q e_q^2 (\bar{h}_{1q}(x_1) h_{1q}(x_2) + (1 \leftrightarrow 2))}{\sum_q e_q^2 (\bar{f}_{1q}(x_1) f_{1q}(x_2) + (1 \leftrightarrow 2))}$$

$$\hat{a}_{TT} = \frac{\sin^2\theta \cos(2\phi - \phi_{S_1} - \phi_{S_2})}{1 + \cos^2\theta}$$

- SSA measurement, $\sin(\phi+\phi_S)$ term
  - $h_1(x)$: transversity
  - $h_1^{\perp(1)}(x)$: Boer-Mulders function (1st moment of)

$$\hat{A} = -\frac{1}{2} \frac{\sum_q e_q^2 (\bar{h}_{1q}^{\perp(1)}(x_1) h_{1q}(x_2) + (1 \leftrightarrow 2))}{\sum_q e_q^2 (\bar{f}_{1q}(x_1) f_{1q}(x_2) + (1 \leftrightarrow 2))}$$

Unpolarized Drell-Yan

- Boer-Mulders function $h_1^{\perp}(x, k_T^2)$
  - angular distribution of unpolarized Drell-Yan

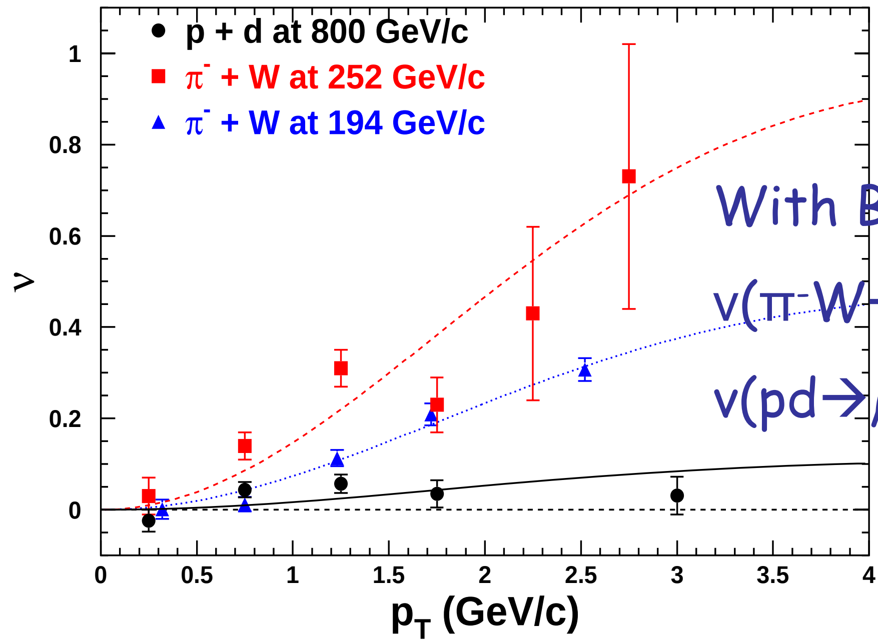
$$\left(\frac{1}{\sigma}\right)\left(\frac{d\sigma}{d\Omega}\right)=\left[\frac{3}{4\pi}\right]\left[1+\lambda\cos^2\theta+\mu\sin 2\theta\cos\phi+\frac{\nu}{2}\sin^2\theta\cos 2\phi\right]$$

  - correlation between transverse quark spin and quark transverse momentum

$$N(\phi)\propto h_1^{\perp q}(x_1,k_\perp^2)\cdot\frac{(\hat{P}\times\vec{k}_\perp)\cdot\vec{S}_q}{M}\cdot h_1^{\perp\bar{q}}(x_2,\bar{k}_\perp^2)\cdot\frac{(\hat{P}\times\vec{\bar{k}}_\perp)\cdot\vec{S}_{\bar{q}}}{M}$$

Unpolarized Drell-Yan

- Boer-Mulders function by unpol. Drell-Yan
  - Lam-Tung relation $1-\lambda=2\nu$
    - reflect the spin-1/2 nature of quarks
  - violation of the Lam-Tung relation suggests non-perturbative origin $\nu \neq 0, 1-\lambda \neq 2\nu$



With Boer-Mulders function $h_1^{\perp}$:

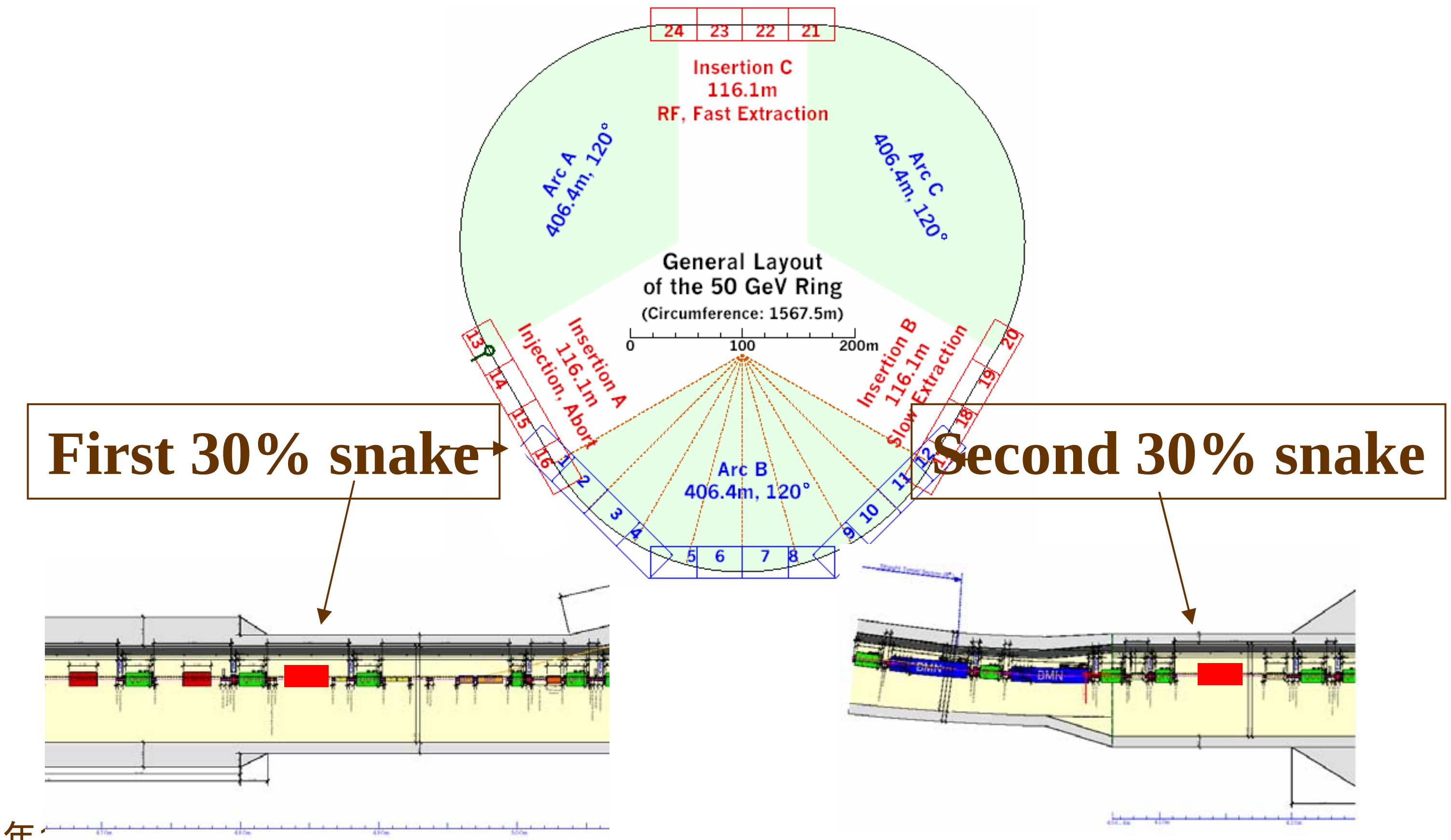
$\nu(\pi^- W \to \mu^+\mu^- X)$ ~ valence $h_1^{\perp}(\pi)$ * valence $h_1^{\perp}(p)$

$\nu(pd \to \mu^+\mu^- X)$ ~ valence $h_1^{\perp}(p)$ * sea $h_1^{\perp}(p)$

L.Y. Zhu, J.C. Peng, P. Reimer et al., hep-ex/0609005

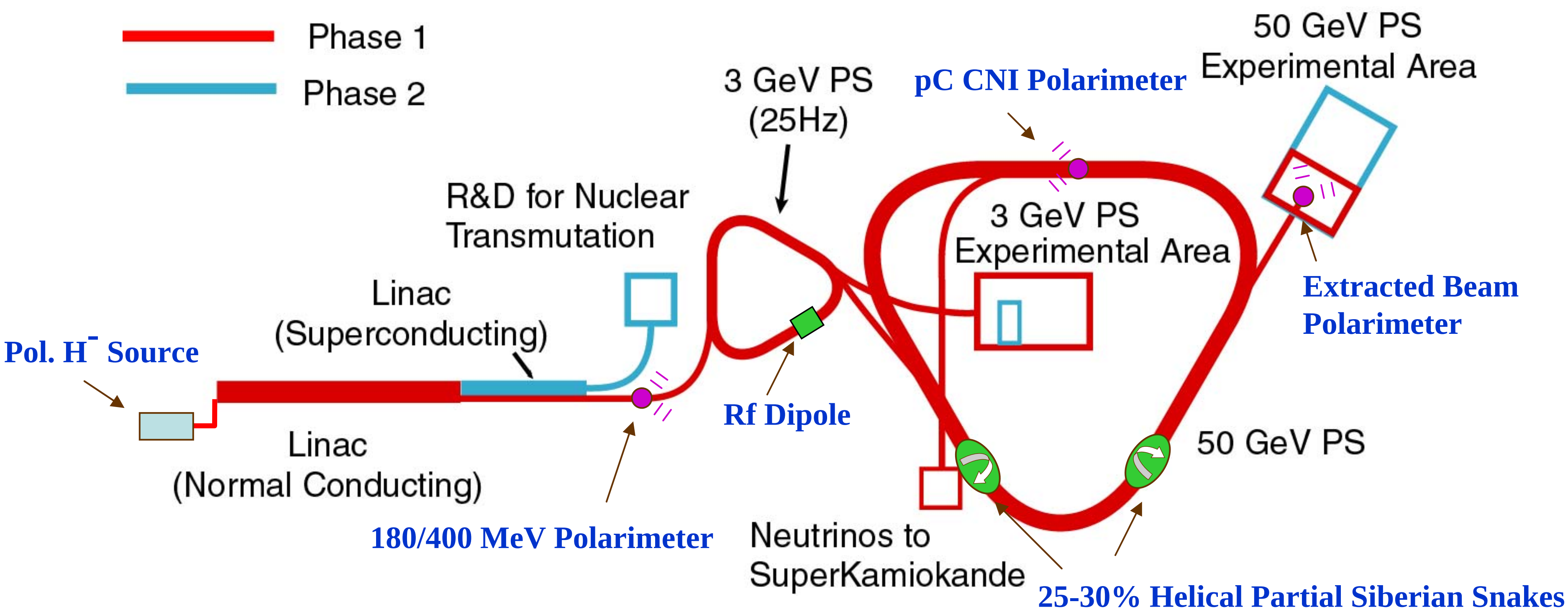
Towards the goal

- 30 GeV $\rightarrow$ 50 GeV
- unpolarized $\rightarrow$ polarized target $\rightarrow$ polarized beam
  - polarized beam study by BNL & KEK groups
  - possible locations of partial snakes in MR



Polarized proton acceleration at J-PARC

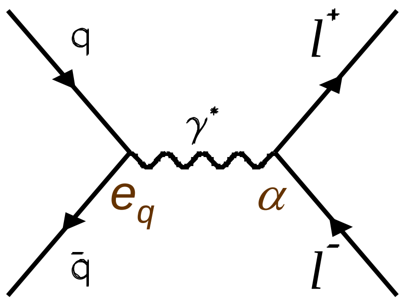
– 50 GeV polarized protons for slow extracted beam primary fixed target experiments
– low intensity (~ 10^{12} ppp), low emittance (10 π mm mrad) beams

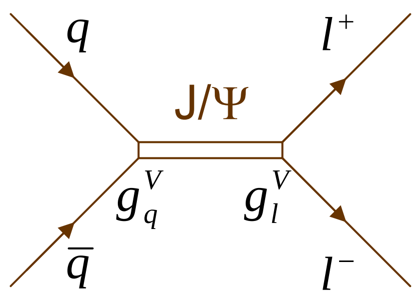


Thomas Roser (BNL), et al.

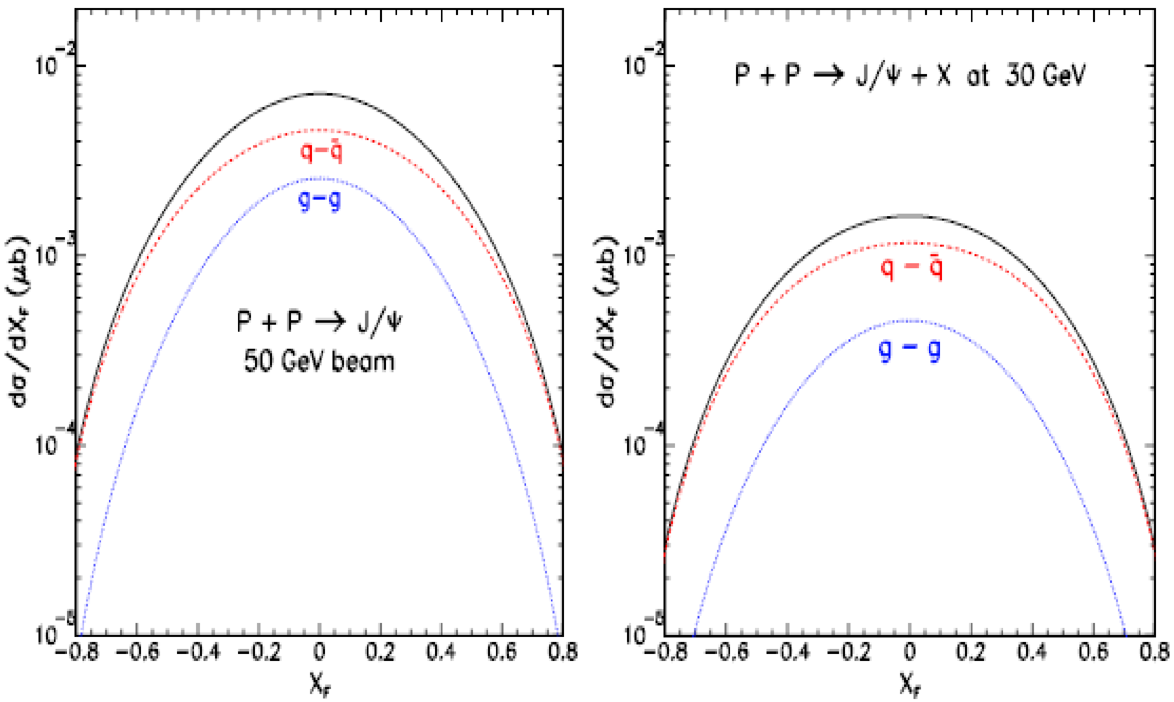
Physics at 30 GeV

- J/ψ
  - gluon fusion or quark-pair annihilation
  - quark-pair annihilation dominant
    - must be confirmed experimentally…
    - similar physics topics as Drell-Yan process





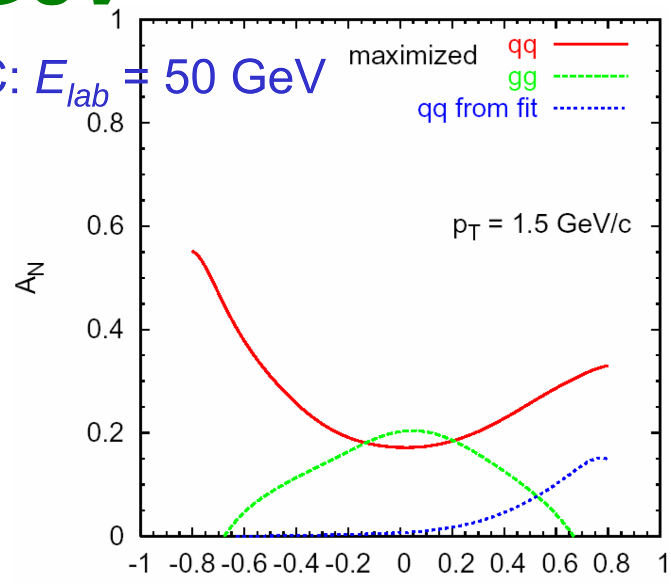
calculations by color-evaporation model



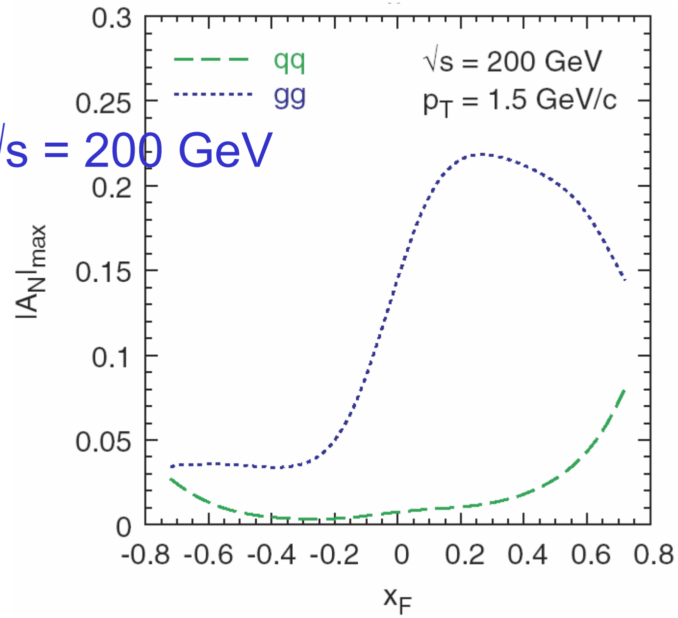
Physics at 30 GeV

- SSA measurement of open charm production
  - no single-spin transfer to the final state
  - sensitive to initial state effect: Sivers effect
  - collider energies: gluon-fusion dominant
    - sensitive to gluon Sivers effect
  - fixed-target energies: quark-pair annihilation dominant
    - sinsitive to quark Sivers effect

J-PARC: E_{lab} = 50 GeV



RHIC: √s = 200 GeV



M. Anselmino, U. D'Alesio, F. Murgia, et al.

SSA measurements of neutron

– large asymmetry found at RHIC

– production mechanism of neutron

  - one-pion exchange dominant – spin flip

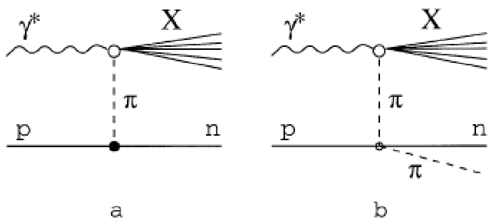


Fig. 1. One-pion-exchange diagrams contributing to the deep-inelastic scattering $ep \rightarrow e'nX$ without (**a**) and with (**b**) production of an additional pion in the proton fragmentation region

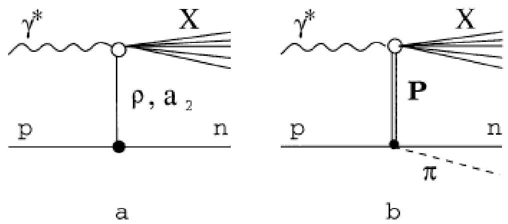


Fig. 2. The ρ- and a_2-Reggeon (**a**) and the Pomeron (**b**) exchange contributions to the deep-inelastic scattering $ep \rightarrow e'nX$

  - figures for study in DIS exps. (HERA), please replace virtual photon with proton to apply them for pp reaction

– asymmetry measurement

  - sensitive to interference between spin-flip term (one-pion exchange) and non-spinflip term (other reggeon exchanges)

– relation to the meson-cloud model? $p \rightarrow n\pi^+$

– neutron-tagged measurement of Drell-Yan?

  - production rate
  - asymmetry measurement

まとめ

- J-PARC dimuon 実験のスピン物理メニュー
  - 縦偏極 A_{LL} of Drell-Yan
    - sea-quark 偏極のフレーバー非対称性
  - 横偏極 A_N of Drell-Yan
    - Sivers関数（sin-term）
    - transversity分布関数 & Boer-Mulders関数（cos-term）
  - 横偏極 A_{TT} of Drell-Yan
    - transversity分布関数
  - その他（30-GeV）
    - J/Ψ: Drell-Yan と同様のメニュー?
    - A_N of open charm: Sivers関数
    - neutron-tagged Drell-Yan