

# Heterotic Type-II Duality

## ♣ Heterotic string vs Type-II string

### Heterotic string

perturbative gauge symmetry

No D-branes, RR fluxes

$$g_s^2 = g_{YM}^2$$

### Type II string

D-branes

gauge symmetry appears non-perturbatively

$$g_s = g_{YM}^2$$

How are they dual to each other?

Heterotic string on  $T^4 \iff$  Type IIA string on  $K_3$  at 6-dimension

### Hull-Townsend

Duality group:  $SO(20, 4)$

Moduli space:  $\frac{SO(20, 4)}{SO(20)SO(4)}$ ,      moduli of  $K_3$  surface:  $\frac{SO(19, 3)}{SO(19)SO(3)}$

20 hypermultiplets

$$H_2(K_3) = 22, 57 + 1 + 22 = 80 = 4 \times 20$$

24 vectors

Further compactification on  $P^1$  down to 4-dimension ( $\mathcal{N} = 2$  susy)

$$T^4 \times P^1 = T^2 \times P^1 \times T^2 = K_3 \times T^2$$

$$K'_3 \times P^1 = \text{CY with } K_3 \text{ fibration}$$

Heterotic on  $K_3 \times T^2 \iff$  type-IIA on CY with  $K_3$  fibration

gauge moduli of heterotic string is dual to geometric moduli of Calabi-Yau.

**Kachru-Vafa**

Consider standard embeddings of spin connection into gauge connection and series of breakings:

$$E_8 \implies E_7 \times SU(2) \implies E_6 \times SU(3) \implies \dots$$

## Example

$$E_8 \times E_8 \times U(1)^4 \implies E_7 \times E_8 \times U(1)^4$$

**vector multiplts :**  $7 + 8 + 4 = 19 = 18 + 1$

**hypermultiplets :**  $45 + 20 = 65$

**Denote as**  $(65, 19) = (M, N)$

## Type-IIA

$$b_{11} = N - 1, \quad b_{12} = M - 1$$

**Breaking of gauge symmetry  $K$  to  $G$**

$$K \supset G \times H$$

$H$  denotes the part broken due to instanton.

Matter from gauge multiplet

$$\text{adjoint } K = \sum_i (M_i, R_i)$$

$M_i, R_i$  are representations of  $G, H$ .

Index theorem

$$\begin{aligned} N_{M_i} &= -\frac{1}{2} \int_{K_3} \text{Tr}_{R_i} F^2 + \frac{1}{48} \dim R_i \int_{K_3} \text{Tr} R^2 \\ &= \dim R_i - \frac{1}{2} \int_{K_3} c_2(V) \times \text{index}(R_i) \end{aligned}$$

**Case  $G = E_7, H = SU(2)$**

$$\int_{K_3} c_2(V) = 24$$

$$248 = (133, 1) + (56, 2) + (1, 3)$$

$$N_{56} = 2 - \frac{1}{2} \cdot 24 \cdot 1 = -10$$

$$N_1 = 3 - \frac{1}{2} \cdot 24 \cdot 4 = -45$$

$$45(\text{gauge moduli}) + 20(K_3 \text{ moduli}) = 65(\text{CY moduli})$$

**Case  $G = E_6, H = SU(3)$**

$$248 = (78, 1) + (27, \bar{3}) + (\bar{27}, 3) + (1, 8)$$

$$N_{27} = 3 - \frac{1}{2} \cdot 24 \cdot 1 = -9$$

$$N_1 = 8 - \frac{1}{2} \cdot 24 \cdot 6 = -64$$

$$(M, N) = (84, 18)$$

These examples have large values of vector multiplets. It is not known if CY manifolds with such Hodge numbers actually exist.

### ♣ Simpler cases with N=3,4

Let us find examples of CY manifolds with small values of N.

#### N=3 case

Choose  $\tau = \rho$  so that  $T^2$  has enhanced gauge symmetry  $SU(2) \times U(1)^3$ .

Then distribute instantons as

$$\begin{array}{ccccccc}
 E_8 & \times & E_8 & \times & SU(2) & \times & U(1)^3 \\
 \uparrow & & \uparrow & & \uparrow & & \\
 10 & & 10 & & 4 & & = 24
 \end{array}$$

$$E_8 \supset E_7 \times SU(2)$$

$$N_{56} = 2 - \frac{1}{2} \cdot 10 \cdot 1 = -3, \quad N_1 = 3 - \frac{1}{2} \cdot 10 \cdot 4 = -17$$

$$SU(2) \supset SU(2)$$

$$N_1 = 3 - \frac{1}{2} \cdot 4 \cdot 4 = -5$$

**Neutral hypermultiplets:**

$$17 \times 2 + 5 + 20 = 59$$

**Higgsing of  $E_7 \times E_7$**

$$2 \times (3 \times 56 - 133) = 70$$

**Thus**

$$M = 59 + 70 = 129, \quad N = 3.$$

**There is a well-known CY manifold with Hodge numbers  $b_{11} = 2$ ,  $b_{12} = 128$ :**



$X_{12}^*[1, 1, 2, 2, 6]$  :

$$W = z_1^{12} + z_2^{12} + z_3^6 + z_4^6 + z_5^2 - \psi z_1 z_2 z_3 z_4 z_5 - \phi z_1^6 z_2^6 = 0$$

**N=4 case**

$$\begin{array}{ccc} E_8 & \times & E_8 & \times & U(1)^4 \\ \uparrow & & \uparrow & & \\ 12 & & 12 & & = 24 \end{array}$$

$$N_{56} = -4, \quad N_1 = -21$$

**Neutral hypermultiplets: =  $2 \times 21 + 20 = 62$**

**$E_7 \times E_7$  Higgsing**

$$2 \times (4 \times 56 - 133) = 182$$

$$62+182=244.$$

Thus we look for CY with  $b_{11}=3$  and  $b_{12}=243$ . Again there is a well-known example of CY manifold with these Hodge numbers.

$$X_{24}^*[1, 1, 2, 8, 12] :$$

$$W = \frac{B}{24}(x_1^{24} + x_2^{24}) - \frac{\psi_2}{12}(x_1x_2)^{12} + \frac{1}{12}x_3^{12} + \frac{1}{3}x_4^3 + \frac{1}{2}x_5^2 = 0$$

$$- \psi_0x_1x_2x_3x_4x_5 - \frac{1}{6}\psi_1(x_1x_2x_3)^6$$

This space also has a K3 fibration. By a change of variable  $x_0 = x_1x_2$ ,  $\zeta = (x_1/x_2)^{12}$  it is rewritten as

$$W = \frac{B'}{12}x_0^{12} + \frac{1}{12}x_3^{12} + \frac{1}{3}x_4^3 + \frac{1}{2}x_5^2 - \psi_0x_0x_3x_4x_5 - \frac{1}{6}\psi_1(x_0x_3)^6$$

$$B' = \frac{1}{2}\left(B\zeta + \frac{B}{\zeta}\right) - \psi_2$$

**Discriminant is given by**

$$\Delta_{CY} = B^2(B^2 - (\psi_1^2 + \psi_2)^2)(B^2 - ((\psi_1 + \psi_0^6)^2 + \psi_2)^2)(B^2 - \psi_2^2)$$

$\Downarrow$   
**Decoupling**

$\searrow \quad \swarrow$   
**SU(3)**

$\Downarrow$   
**LCS**

***SU(3)* gauge theory limit is taken as**

$$B = \epsilon\Lambda^3, \quad \psi_0^6\psi_1 = \epsilon u^{3/2}, \quad \psi_1^2 + \psi_2 = \epsilon(v - u^{3/2}), \quad \epsilon \rightarrow 0$$

**By a suitable redefinition of variables we obtain an  $A_2$  singularity fibered over  $P^1$**

$$W = \epsilon \left[ \frac{1}{12} \left( \zeta + \frac{\Lambda^6}{\zeta} \right) + \frac{y_3^2}{2} + \frac{y_4^2}{2} + \frac{y_5^3}{3} - \frac{u}{4} y_5 - \frac{v}{12} \right]$$

***K*<sub>3</sub> fibration**

Define heterotic (axion-)dilaton as

$$S = \frac{1}{g^2} - \frac{i\theta}{8\pi^2}$$

Heterotic prepotential is then given by

$$F = F_0(S, T^i) + F_1(T^i) + \dots, \quad F_0 = -\gamma_{ij} T^i T^j S$$

$$K = -\log(S + \bar{S}) - \log \gamma_{ij} (T^i + \bar{T}^i)(T^j + \bar{T}^j)$$

Type II side

$$B + iJ = \sum_{\alpha=1}^{h^{11}} (B + iJ)_\alpha e_\alpha$$

where  $\{e_\alpha\}$  is a basis of  $H_2(X; \mathbb{Z})$ . Divisor  $D_\alpha$  is Poincare dual to  $e_\alpha$ .

Type-II prepotential has a form

$$F_0 = -\frac{i}{6} \sum_{\alpha, \beta, \gamma} \langle D_\alpha \cdot D_\beta \cdot D_\gamma \rangle t_{\alpha\beta\gamma} + \text{world-sheet instantons}$$

**Heterotic dilaton maps to some divisor  $D_S$  and Kähler parameter  $(B + iJ)_S = 4\pi i S$  of type IIA side.  $B \rightarrow B + 1$  corresponds to  $\theta \rightarrow \theta + 2\pi$ . Then we have**

$$D_S \cdot D_S \cdot D_S = 0, \quad D_S \cdot D_S \cdot D_i = 0 \implies D_S \cdot D_S = 0$$

**Then according to a mathematical theorem CY has a  $K_3$  fibration over  $P^1$ . Size of  $P^1$  is  $\sim 1/g^2$ .**

## Problems

**More example of dual pairs**

**Moduli stabilization**

**Understanding of gauge  $\leftrightarrow$  geometry moduli**

**(Perturbative?) identification of type-II D-branes**

**Phenomenology**