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Wigner Function Shapelets and its application

S. Arai, arXiv: 2602.01141, PRD submitted

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Background : Shapes of astronomical bodies tells rich information

- Images of astronomical bodies contain information of fundamental physics
- Galaxy shapes are essential to probe cosmology/gravity

cf. S.Arai et.al 2023 +

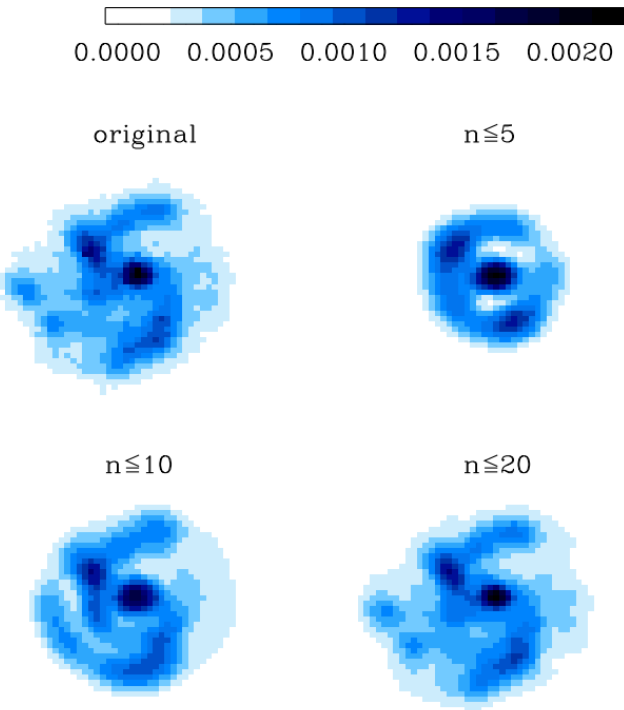
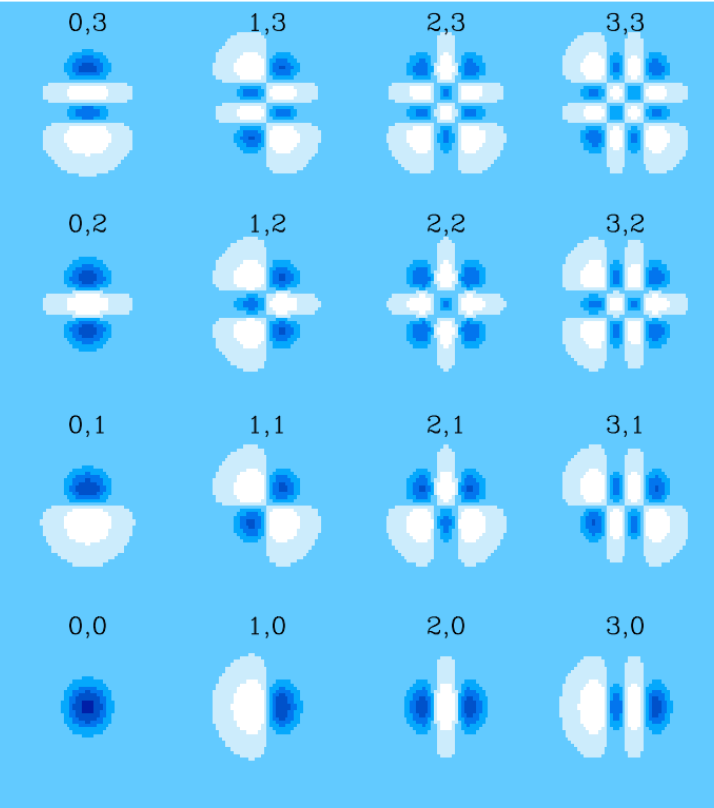
- Galaxy evolution imprinted in its shape
- Shape distortions by gravitational lensing
- Higher shape moments of galaxies can exhibit 3D parity violation, higher-spin moments etc.
- Machine-learning Simulation-based inference



©Galaxy Zoo from Euclid

2D representation of images: Shapelets

A.Refregier 2003, R.Massey and A.Refregier 2004 +

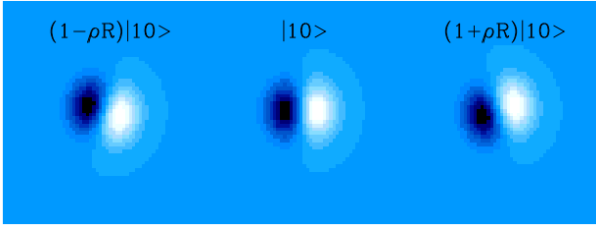


$\psi(\mathbf{x})$ image field

$$\psi(\mathbf{x}) = \sum_{nm} \psi_{nm} \Psi_{nm}(\mathbf{x})$$

- Hermite/Laguerre-Gaussian modes are adopted

-grav. lensing follows selection rule ($m \rightarrow m \pm 2$)



A.Refregier 2003 Fig.2, 3

Motivation: Why Wigner Function?

- Shapelets in config/Fourier sp. are mathematically well-defined but...
 - ? How does it relate to image propagation/processing
e.g. radiative transfer, point spreading, sensed on CCD...

What this study reveals is

“Wigner Function Shapelets” organises

- ✓ Wigner transport is a counterpart of transfer eq. , given a resolution limit of a telescope
- ✓ group-theoretic representation of observables
- ✓ Local signatures in phase space reflects the nature of galaxy morphology

Intro. Wigner Function

$$W(\mathbf{x}, \mathbf{p})[\psi] = (2\pi\lambda)^{-2} \int d^2\xi \psi(\mathbf{x} + \xi/2) \psi^*(\mathbf{x} - \xi/2) \exp(-i\mathbf{p} \cdot \xi/\lambda)$$

H.Weyl 1927, E. Winger 1932

$\psi(\mathbf{x})$: Image field

Parity invariant

Symplectic covariant

Exp. val.

$$\langle \hat{A} \rangle = \text{Tr}(\hat{\rho} \hat{A}) = \int d^2x d^2p A_W(\mathbf{x}, \mathbf{p}) W(\mathbf{x}, \mathbf{p})$$

“Weyl symbol”

$$A_W(\mathbf{x}, \mathbf{p}) = \int d^2\xi A(\mathbf{x} + \xi/2, \mathbf{x} - \xi/2) \exp(-i\mathbf{p} \cdot \xi/\lambda)$$

Normalisation

$$1 = \text{Tr}(\hat{\rho}) = \int d^2x d^2p W(\mathbf{x}, \mathbf{p})$$

“quasi”probability dist (\because it can be negative)

Intensity

$$I(\mathbf{x}) = \int d^2p W(\mathbf{x}, \mathbf{p})$$

Power spectrum

$$F(\mathbf{p}/\lambda) = \int d^2x W(\mathbf{x}, \mathbf{p})$$

Intro. Cross Wigner Function

$$W(\mathbf{x}, \mathbf{p})[\psi_a, \psi_b] = (2\pi\lambda)^{-2} \int d^2\xi \psi_a(\mathbf{x} + \boldsymbol{\xi}/2) \psi_b^*(\mathbf{x} - \boldsymbol{\xi}/2) \exp(-i\mathbf{p} \cdot \boldsymbol{\xi}/\lambda)$$

★ a, b: polarisation (I, Q, U, V), colours (ugriz), quantum number

Parity inv. iff. $a = b$

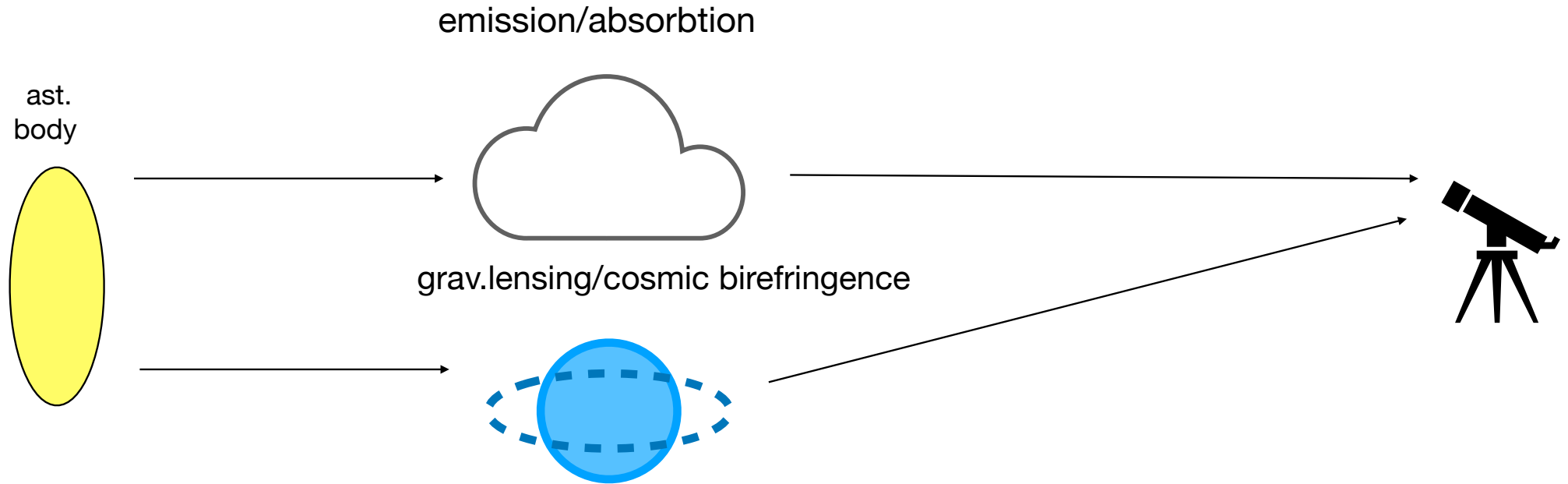
Symplectic cov.

Moyal/Plancherel identity

$$\langle W_{cd} | W_{ab} \rangle_{\text{HS}} = \frac{1}{(2\pi\lambda)^2} \langle \psi_a | \psi_c \rangle \langle \psi_b | \psi_d \rangle$$

$\langle \psi_a | \psi_b \rangle \propto \delta_{ab} \rightarrow$ Hilbert-Schmidt product becomes orthogonal

Intro. Wigner transport



Wigner func.

$$W = W(\mathbf{x}, \mathbf{p}, s)$$

$$\partial_s W = \{H_W, W\}_{\text{Moyal Product}}$$

Specific intensity

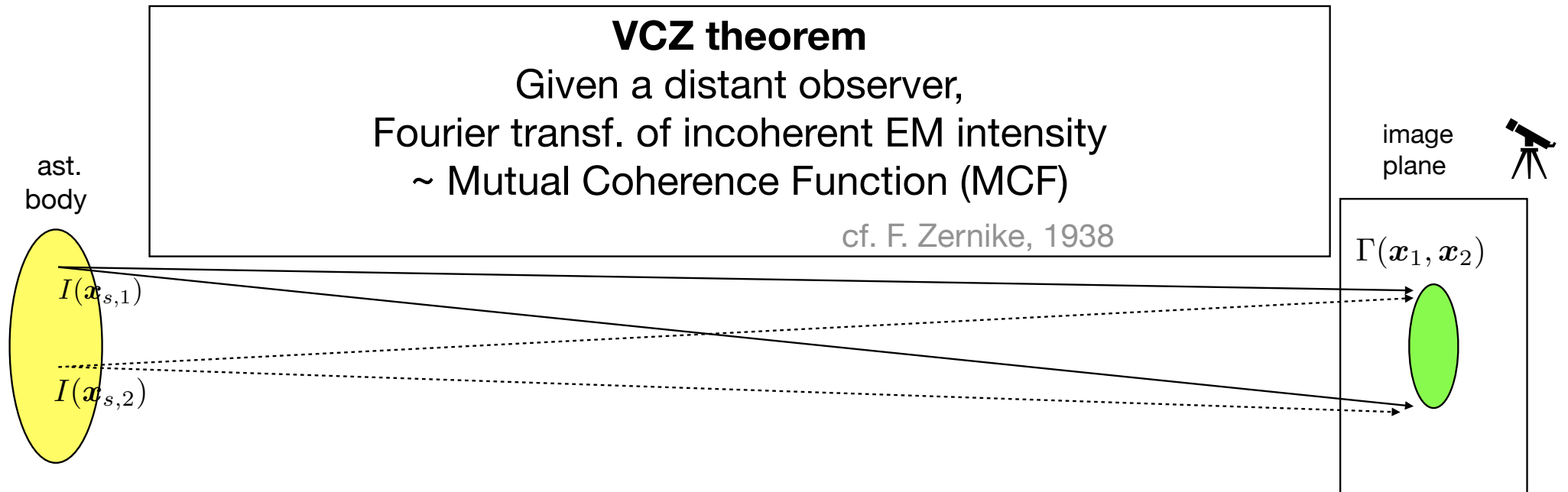
$$L = L(\mathbf{x}, \mathbf{k}, s)$$

$$\lambda \rightarrow 0 \quad \mathbf{k} \equiv \mathbf{p}/\lambda$$

$$\partial_s L = \{H, L\}_{\text{Poisson Bracket}}$$

✓ Wigner transport naturally corresponds to radiative transf.

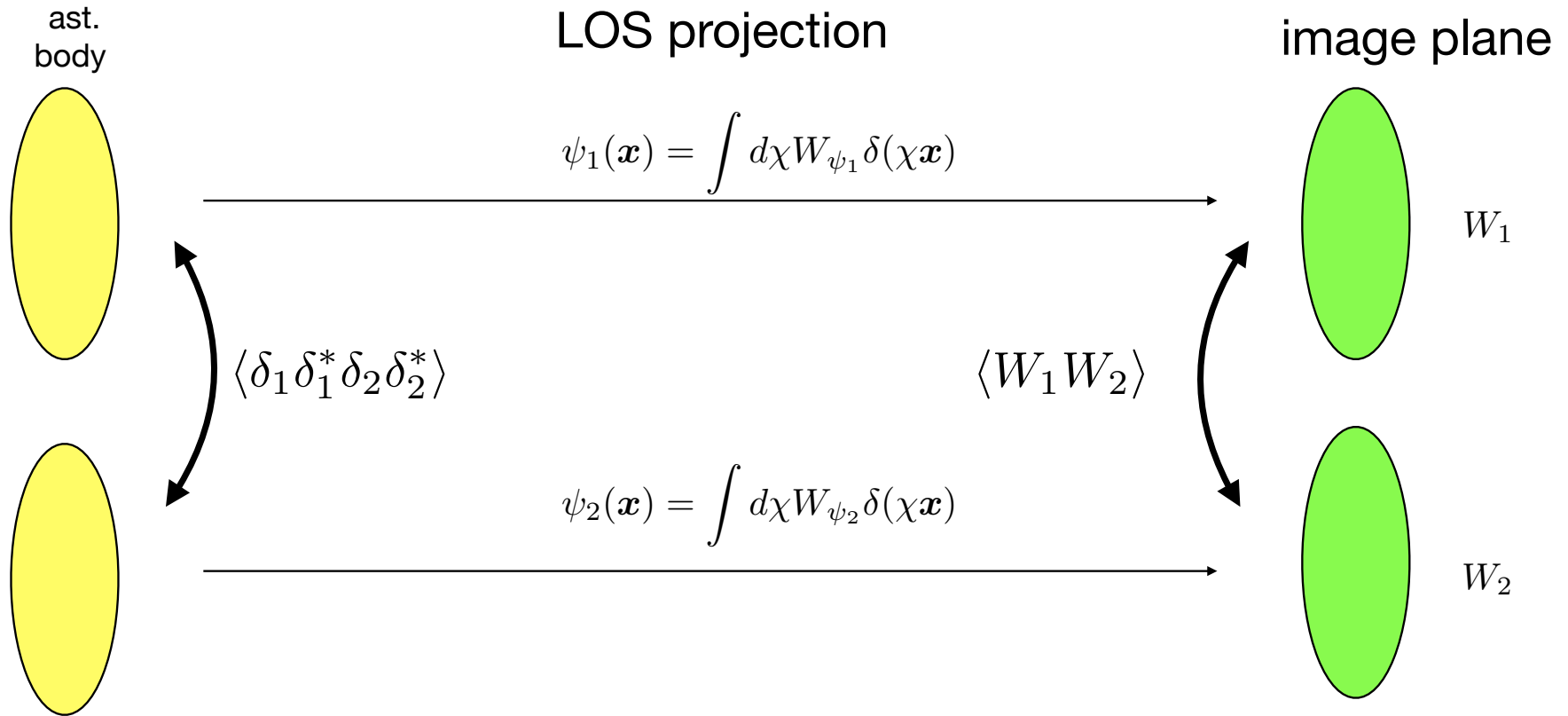
Wigner func. and Van-Cittert Zernike (VCZ) theorem



$$W(\mathbf{x}, \mathbf{p}) = \int d^2\xi \Gamma(\mathbf{x} + \xi/2, \mathbf{x} - \xi/2) \exp(-i\mathbf{p} \cdot \xi/\lambda)$$

✓ Wigner function is Fourier inverse transf. of MCF

e.g. Wigner func. and cosmo. 4pt func.



✓ Wigner func. at different sky position \sim 4pt (collapsed limit)

Wigner Function Shapelets

R.Simons and G.Agarwal 2004 +

Point : Cross Wigner func. for Laguerre-Gaussian modes preserves SU(2)

$$W_{jm}^{\text{LG}}(Q_0, Q_2) \propto (-1)^{2j} e^{-4Q_0} L_{j+m}(4(Q_0 + Q_2)) L_{j-m}(4(Q_0 - Q_2))$$

$$Q_0 = \frac{1}{2} \left[\frac{|\mathbf{x}|^2}{w^2} + \frac{w^2 |\mathbf{p}|^2}{4\lambda^2} \right] \quad Q_2 = \frac{1}{2\lambda} \mathbf{x}^T J \mathbf{p} \quad J = \begin{pmatrix} 0 & 1 \\ -1 & 0 \end{pmatrix}$$

Cross Wigner func. is analytically diagonalised by Swinger bosonic oscillator

$$\begin{aligned} & W[\Psi_{js}^{\text{LG}}, \Psi_{j's'}^{\text{LG}}](Q_0, Q_2; \phi_u, \phi_v) \\ &= \frac{(-1)^{n_{u<} + n_{v<}}}{\pi^2 \lambda^2} \sqrt{\frac{n_{u<}! n_{v<}!}{n_{u>}! n_{v>}!}} (\sqrt{2} |u|)^{|\Delta_u|} (\sqrt{2} |v|)^{|\Delta_v|} \\ & \quad \times e^{-2(|u|^2 + |v|^2)} e^{i(\Delta_u \phi_u + \Delta_v \phi_v)} \\ & \quad \times L_{n_{u<}}^{(|\Delta_u|)}(4|u|^2) L_{n_{v<}}^{(|\Delta_v|)}(4|v|^2), \end{aligned}$$

★ Hilbert-Schmidt space w. orthogonal and complete basis

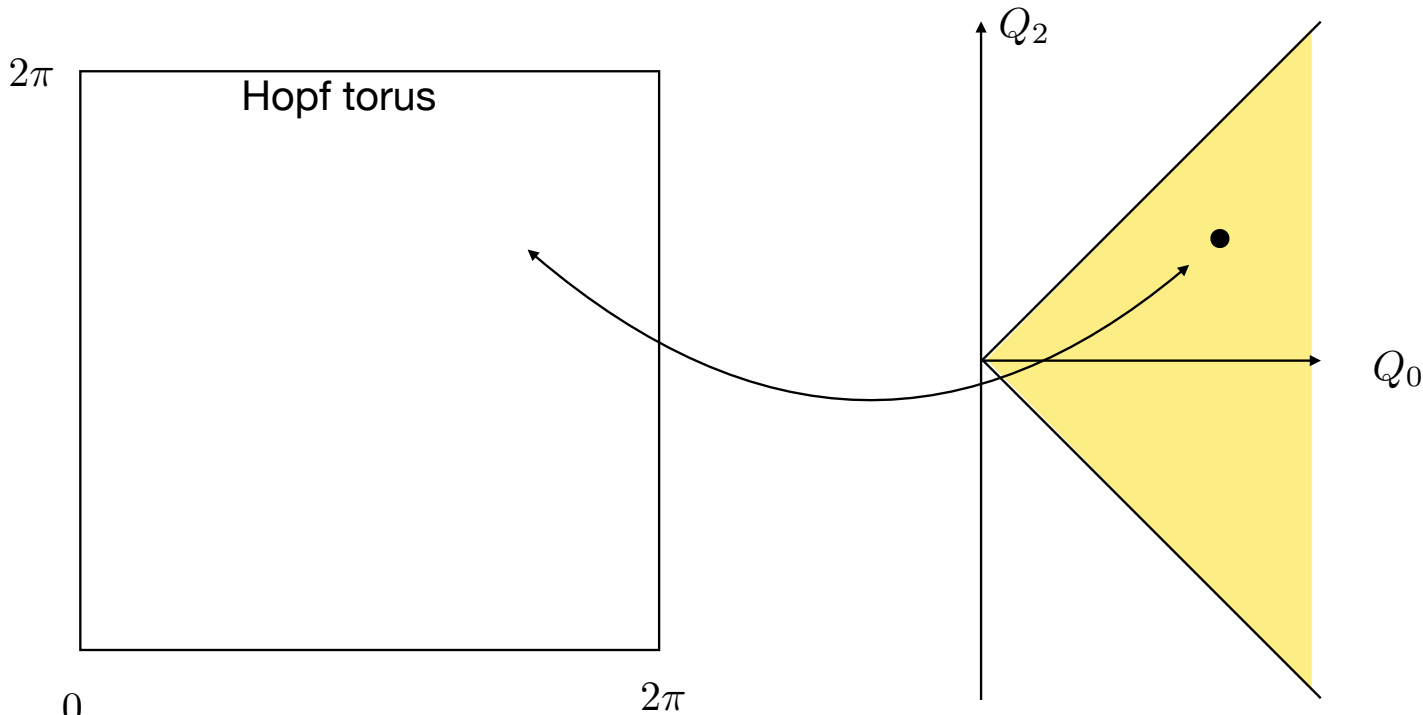
Compression of Wigner Function Shapelets to 2D plane

WFS is complex in 4D phase space that prevents me from intuitive analyses

★ WFS is further characterised by "Hopf Fibration" H. Hopf 1931+

$$SU(2) \supset U(1) \times U(1)$$

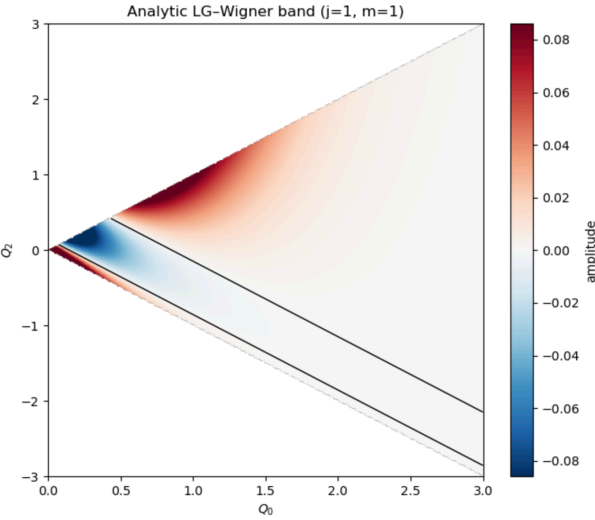
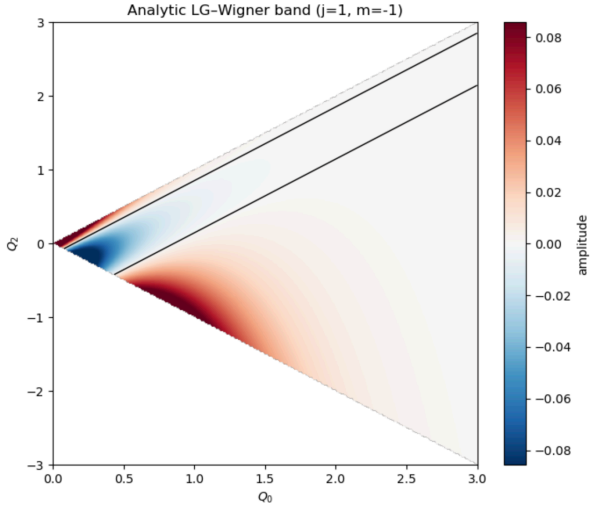
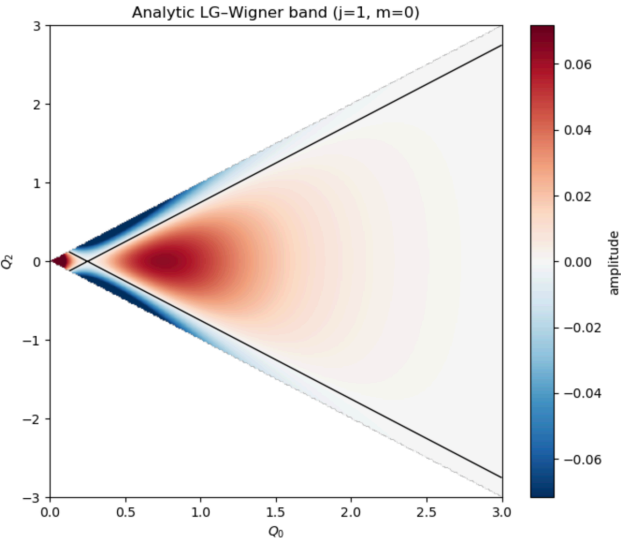
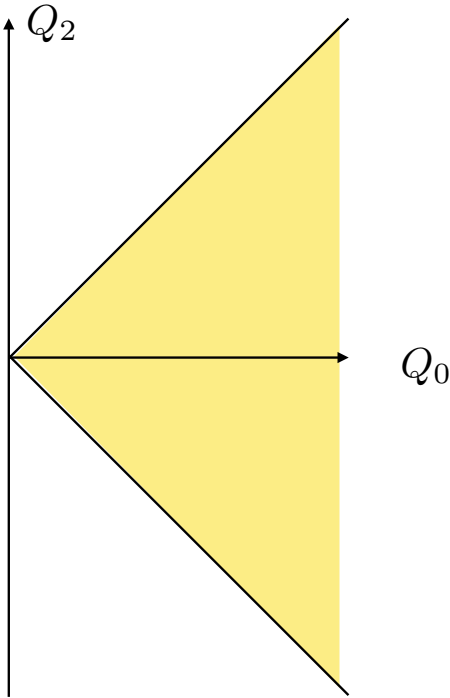
$$u = \sqrt{Q_0 + Q_2} e(\phi_u), v = \sqrt{Q_0 - Q_2} e(\phi_v) \quad e(\phi) = {}^T(\cos(\phi), \sin(\phi))$$



✓ group-theoretic operation preserves information during compression

Result 1: Wigner Function Shapelets "band structure"

$$W_{jm}^{LG}(Q_0, Q_2) \propto (-1)^{2j} e^{-4Q_0} L_{j+m}(4(Q_0 + Q_2)) L_{j-m}(4(Q_0 - Q_2))$$

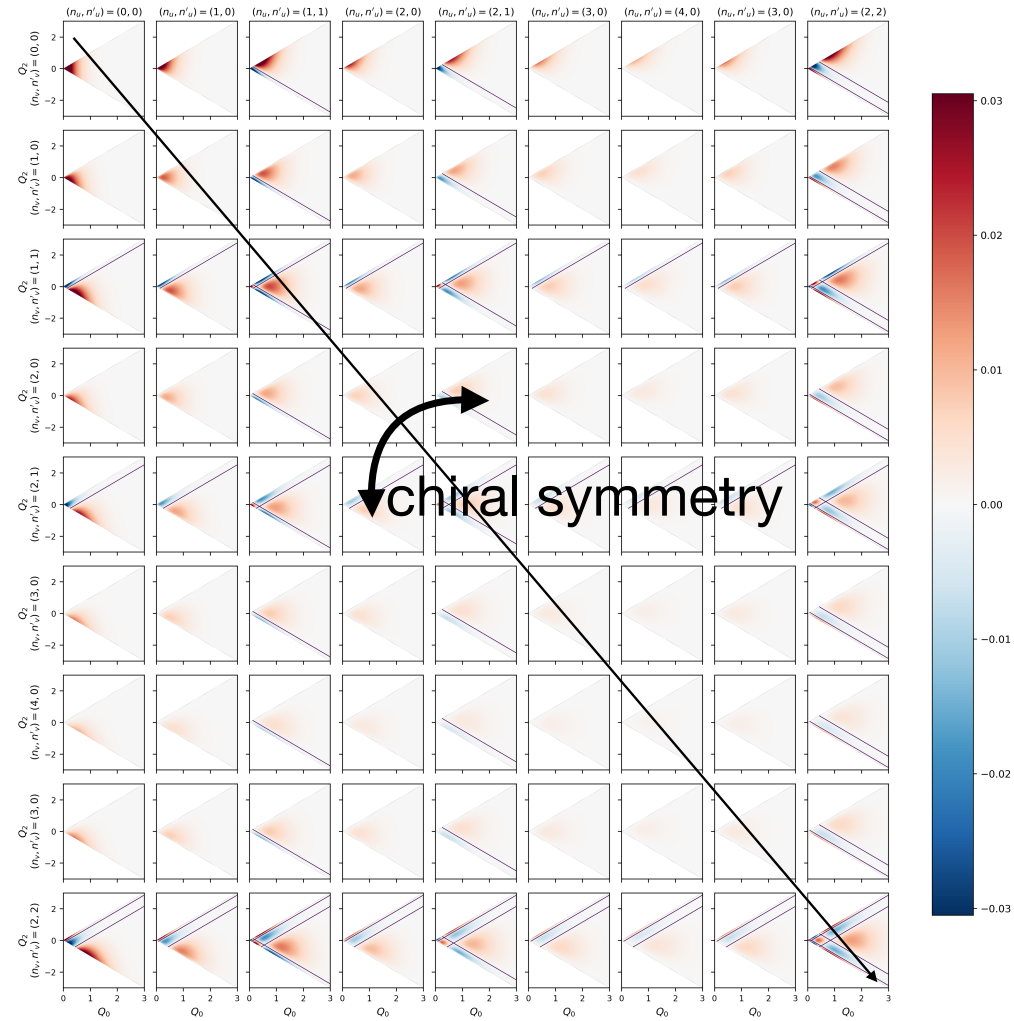


$Q_2 \rightarrow -Q_2$

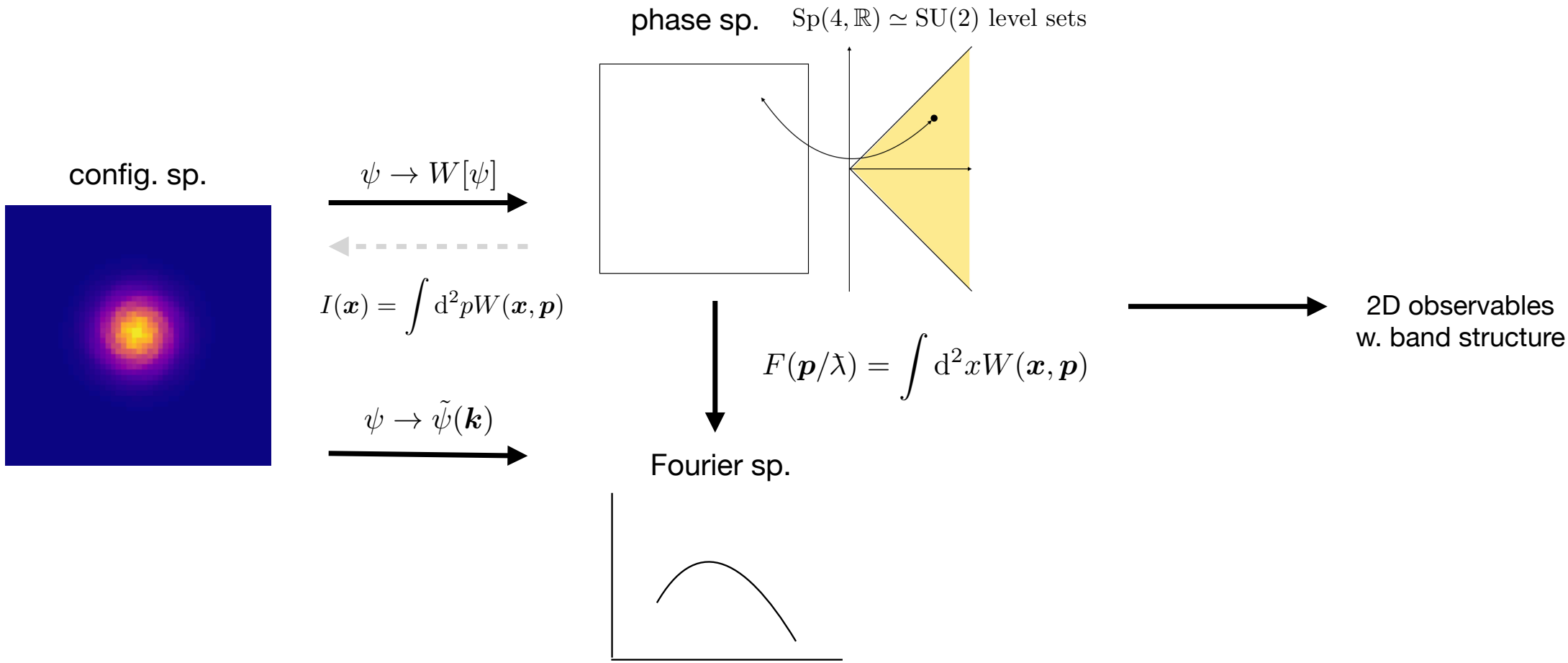
- ★ WFS exhibits +/-/0 in geometry i.e. "band structure"
 - ★ Wigner is everywhere positive iff. it is a Gaussian state
- * single image

WFS rep.

$$\mathcal{W}_{kl}(Q_0, Q_2) = \int \frac{d\phi_u}{2\pi} \frac{d\phi_v}{2\pi} W \exp(-ik\phi_u - il\phi_v)$$

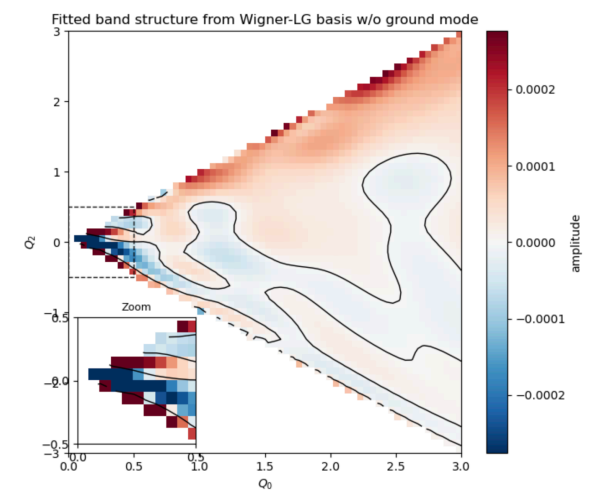
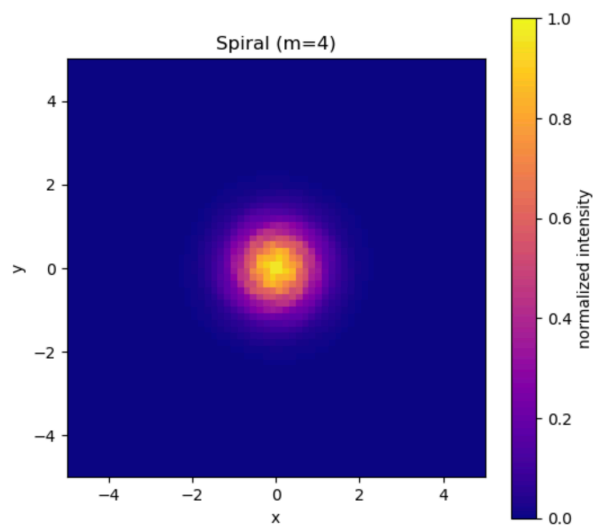
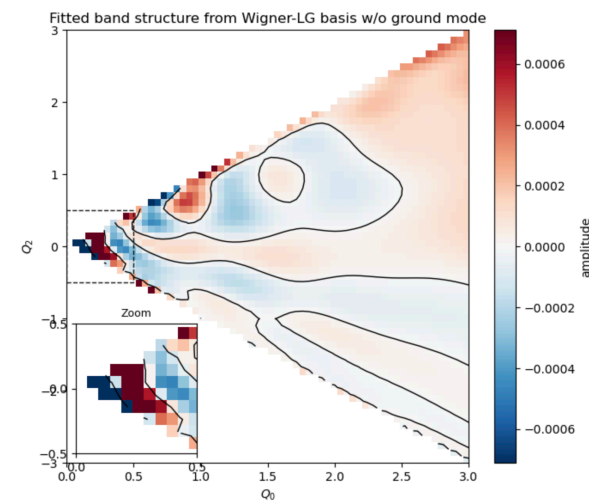
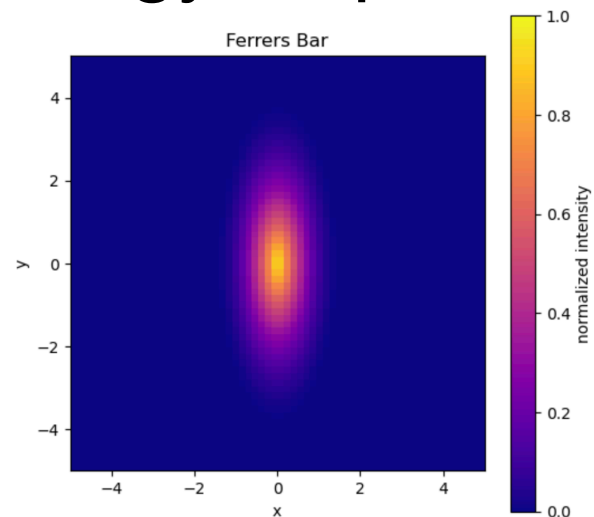


Q. How to use WFS?



Result 2 : Galaxy morphology response

$$W = \sum_{j=0, \frac{1}{2}, 1, \dots} \sum_{m=-j}^j W_{jm} W_{jm}^{\text{LG}} + \text{cross Wigner}$$



Takeaways

- We extend the conventional shapelets, employing Wigner function

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prospects :

- Mock data challenges (Point spread function, model bias, deblending...)
- app. for extended fields (CMB, LSS, 21cm BT etc.) by Wigner function

“The great advances in science usually result from new tools rather than from new doctolines...”
F. Dyson, NYBR essay 1995