

# A Quantum Mechanical Model of Black Hole

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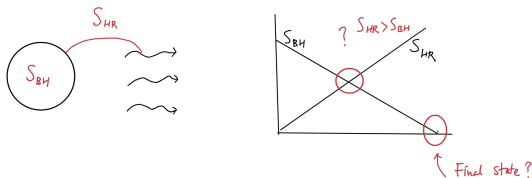
# Fundamental problems of quantum black hole

Black hole provides an ideal (only!?) window to quantum gravity!

- Bekenstein-Hawking entropy:

Microscopic origin of  $S_{\text{B-H}} = \frac{A}{4G\hbar}$ ?

- Information problem:

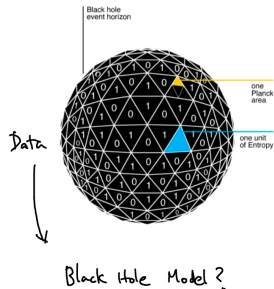
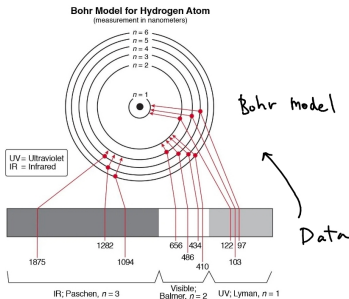


- BH chaos: most efficient quantum information scrambler  
(Susskind, Maldacena-Shenker-Stanford, Shenker et al)

$$\tau_s \sim T^{-1} \log S$$

Resolution of these problems require understanding of the quantum physics of black hole.

- my understanding: Horizon and its infinite redshift is the origin to ALL of the black hole problems.
- We want to construct a quantum mechanical model of horizon and learn about quantum black hole.  
This is like Bohr model (data: spectral line. learnt: nucleus, eigenstates, ...)



# Outline

1 Part I: Model+ black hole checks

2 Part II: Hawking radiation

# Quantum mechanical model of black hole

- We consider the large  $N$  quantum mechanics of 3d quantum space

$$L = \text{tr} \left[ \frac{1}{2a_0^2 M_P} \dot{X}^{a2} + \frac{M_P}{N^2} ([X^a, X^b]^2 + 4X^{a2}) + i\psi^\dagger \psi - a_2 \frac{M_P}{N^2} \psi^\dagger \sigma^a X^a \psi \right]$$

where  $a = 1, 2, 3$ ,  $X_{mn}^a, \psi_{mn}, \psi_{mn}^\dagger$   $N \times N$  traceless matrices.

- Remarks:
  - 1 Like BFSS but no SUSY.
  - 2 tachyonic mass term instead of Chern-Simons mass term: important for capturing the instability associated with the infinite red shift of horizon.
  - 3  $a_0, a_2$  coefficients to be fixed.

# Quantum Schwarzschild as Fuzzy Sphere

- The EOM admits static solution given by  $K \times K$  matrices obeying

$$[X^a, X^b] = i\epsilon^{abc} X^c, \quad \sum_a X^{a2} = \frac{K^2 - 1}{4} \mathbf{1},$$

- Introduce dimensional coordinates  $Y^a = 2l_P X^a$ , the fuzzy sphere solution becomes

$$[Y^a, Y^b] = \frac{2iR}{\sqrt{K^2 - 1}} \epsilon_{abc} Y^c, \quad \sum_a Y^{a2} = R^2 \mathbf{1},$$

where  $R = Kl_P$  is the radius of the fuzzy sphere for  $K$  large.

## Fermionic states

- Quantizing the fermions over the fuzzy sphere background, we obtain  $2N^2$  fermionic oscillators

$$H_F = \frac{a_2 M_P}{2N} \sum_{p,k=1}^N (\xi_k^{p\dagger} \xi_k^p + \chi_k^{p\dagger} \chi_k^p).$$

- For a half filled Fermi sea, we reproduce precisely the Schwarzschild radius

$$R = \frac{M}{2G}$$

for static black hole.

# Microstate counting

- Let us consider the microstates counting. For the half-filled Fermi sea, there is a degeneracy of

$$\Omega_0 = \binom{2N^2}{N^2} = 2^{2N^2}$$

in the leading large  $N$  limit. These microstates give rises to the entropy  $S = \log_2 \Omega_0$ :

$$S = 2N^2 = \frac{A}{4G},$$

which is precisely the Bekenstein-Hawking entropy.

# Kerr black hole

- A similar analysis can be performed for rotating black hole

$$ds^2 = - \left( 1 - \frac{2M}{\rho^2} \right) dt^2 - \frac{4Mar \sin^2 \theta}{\rho^2} dt d\phi + \frac{\Sigma}{\rho^2} \sin^2 \theta d\phi^2 + \frac{\rho^2}{\Delta} dr^2 + \rho^2 d\theta^2,$$

$$\rho^2 := r^2 + a^2 \cos^2 \theta, \Delta := r^2 - 2Mr + a^2, \Sigma := (r^2 + a^2)^2 - a^2 \Delta \sin^2 \theta.$$

- The Kerr BH has a horizon at  $r_+ = M + \sqrt{M^2 - a^2}$ . Viewed in the Cartesian coordinates, the horizon is an elliptical surface

$$\frac{x^2 + y^2}{r_+^2 + a^2} + \frac{z^2}{r_+^2} = 1.$$

- Bekenstein-Hawking entropy, mass and angular momentum

$$S = \frac{A}{4G} = \frac{\pi(r_+^2 + a^2)}{G}, \quad M = \frac{r_+^2 + a^2}{2Gr_+}, \quad J = a \frac{r_+^2 + a^2}{2Gr_+}.$$

- We showed that the rotating fuzzy sphere solution of the QM reproduces precisely all of these properties of Kerr black hole.
- In particular, matching of the angular momentum requires

$$a_0 = \frac{\pi}{3}.$$

This will be important later.

# Outline

- 1 Part I: Model+ black hole checks
- 2 Part II: Hawking radiation

- Semi-classical QFT computation of Hawking radiation gives the decay rate of BH

$$\Gamma = -\frac{1}{M} \frac{dM}{dt} = -\frac{Q}{G^2 M^3}.$$

- It gives also the Hawking temperature

$$T_H = \frac{1}{8\pi GM}.$$

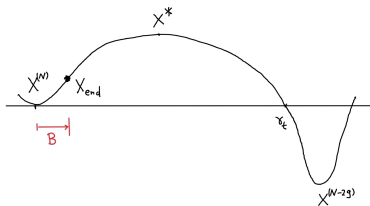
Our model of BH provides a natural explanation of this process in terms of QM tunneling of fuzzy spheres. Both results are reproduced.

## Monopole channel and decay rate

- We have a collection of fuzzy sphere vacua

$$E = \frac{bNM_P}{2}$$

of different sizes (rank  $N$ ). Decay can occur by tunneling.



- However since there is an excess of Fermi states in the Fermi sea of  $X^{(N)}$ , it cannot decay to the small fuzzy sphere with a smaller Fermi sea unless the decay path has the right number of fermi zero modes to soak up the excess fermi states.

- The criteria is so strong that the tunneling path is uniquely determined to be

$$X_a = X_a^{(N)} - \gamma(X_a^{(N)} - X_a^{(N')}), \quad 0 \leq \gamma \leq 1.$$

- The difference of two fuzzy spheres

$$A_a^{(n)} := X_a^{(N)} - X_a^{(N-n)},$$

is a gauge field which has the property that the Dirac equation on  $S_N^2$  has exactly  $n$  zero modes,

- Moreover in the large  $N$  limit,

$$A = A_a dx^a = \frac{n}{2}(1 - \cos \theta)d\varphi$$

i.e. precisely the Dirac monopole connection on  $S^2$  with charge  $g = n/2$ .

- For a quantum mechanics

$$L = \frac{1}{2}m\dot{x}^2 - V(x),$$

the tunneling rate is given by (Coleman)

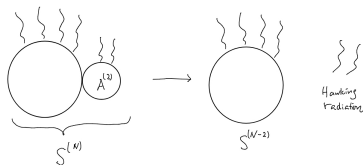
$$\Gamma = Ke^{-S_b}, \quad K = \frac{1}{2} \sqrt{\frac{S[x_b] - S[x_0]}{2\pi\hbar m}} \sqrt{\frac{\text{Det}S''[x_0]}{\text{Det}'S''[x_b]}}.$$

- computing the path integral, we obtain the tunneling rate

$$\Gamma \sim \frac{1}{N^3} \sim -\frac{1}{G^2M^3}.$$

Up to a numerical constant, this reproduces the semi-classical decay rate of black hole.

# Hawking temperature



- Physically as the monopole is nucleated and leave  $X^{(N)}$ , it leaves behind the smaller fuzzy sphere. At the same time, the released fermi states appear as the Hawking radiation.
- We find the emission probability for a Hawking quanta with energy  $\omega$  to be

$$P(\omega) = e^{-\omega/T_H},$$

This is a thermal distribution with the Hawking temperature  $T_H$ !

- With a real time formulation of the tunneling process, one can go beyond the probabilistic description by determining the full wave function of the multi-patite Hawking radiation.

## Interesting applications:

- other properties of Hawking radiation: Page curve? information?
- final state? PBH? remanants?
- memory burden?
- spacetime uncertainty principle from matrix model?
- black hole singularity?
- chaos?
- metric? forces between black hole ?
- holography?

Thank you for listening!