

Relational entanglement entropies in gauge theory and gravity

Philipp Höhn

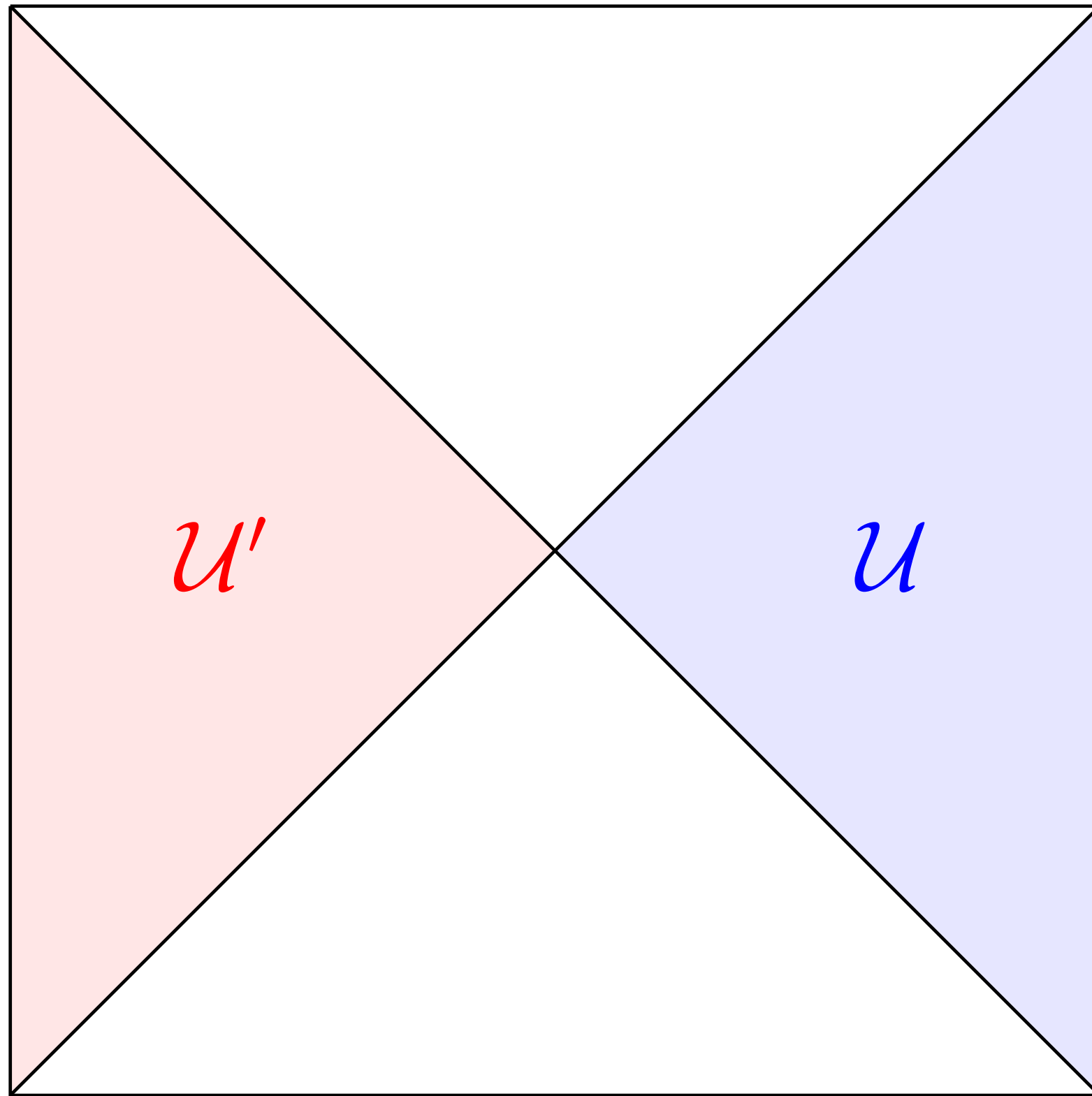
Okinawa Institute of Science and Technology



Concepts of Quantum and Spacetime
9 March 2026

based on: [De Vuyst, Eccles, PH, Kirklin, JHEP 07 \(2025\) 063; JHEP 07 \(2025\) 146](#)
[Araújo-Regado, PH, Sartini, arXiv:2506.23459](#)

Entanglement entropy in QFT



e.g. dS space

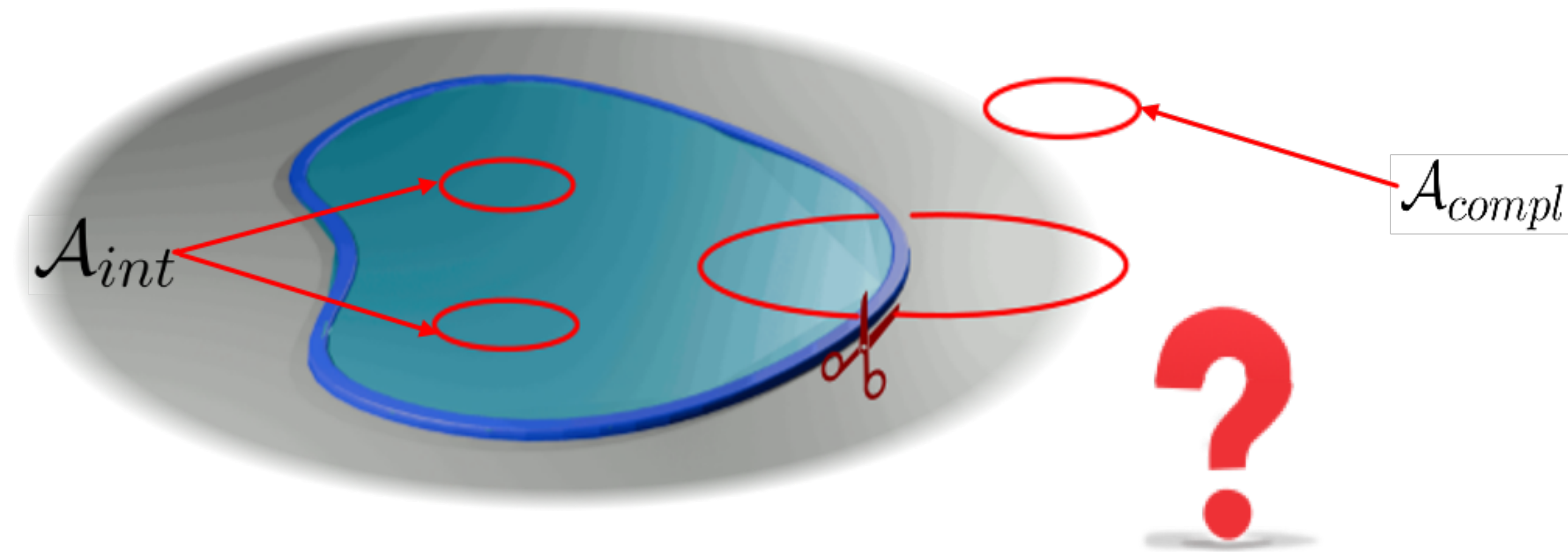
2 challenges:

- **UV divergences** (∞ many DoFs around entangling surface)
 \Rightarrow **usually: regulators** (but in tension with covariance/invariance)
- **Locality in the presence of gauge symmetry**
 \Rightarrow **what do we mean by subsystem?**

Subsystems in gauge systems

Gauge constraints nonlocal: $\mathcal{H}_{\text{phys}} \neq \mathcal{H}_{\text{int}} \otimes \mathcal{H}_{\text{compl}}$

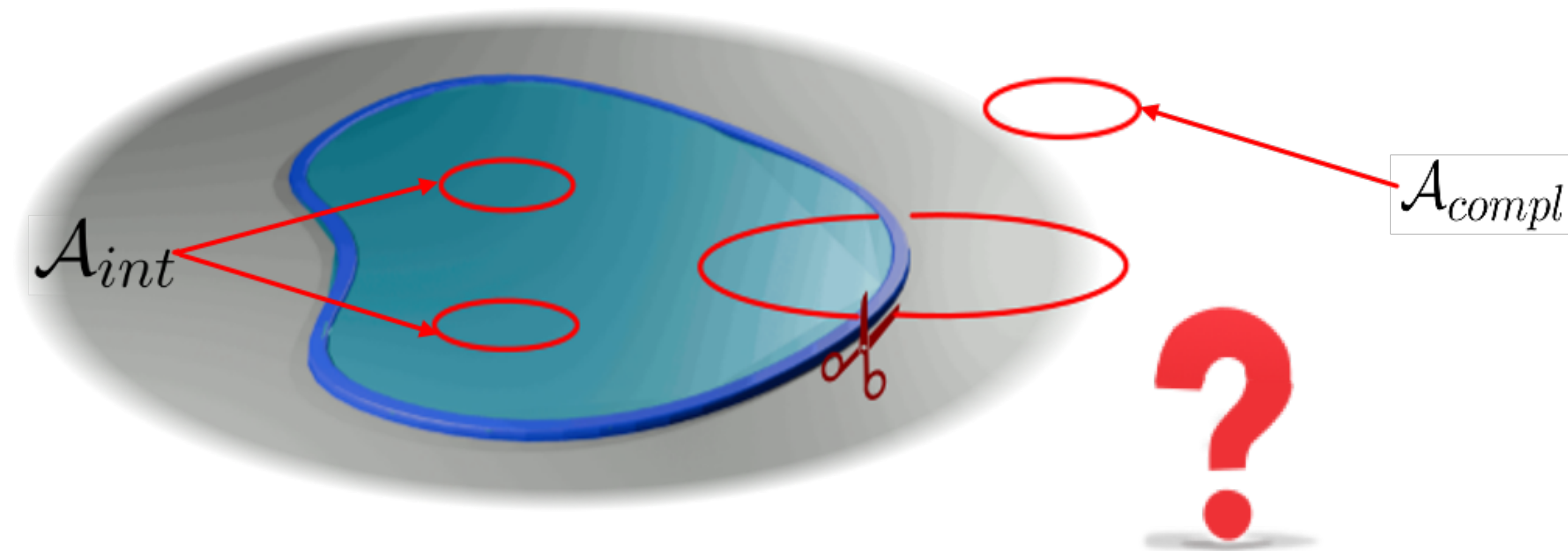
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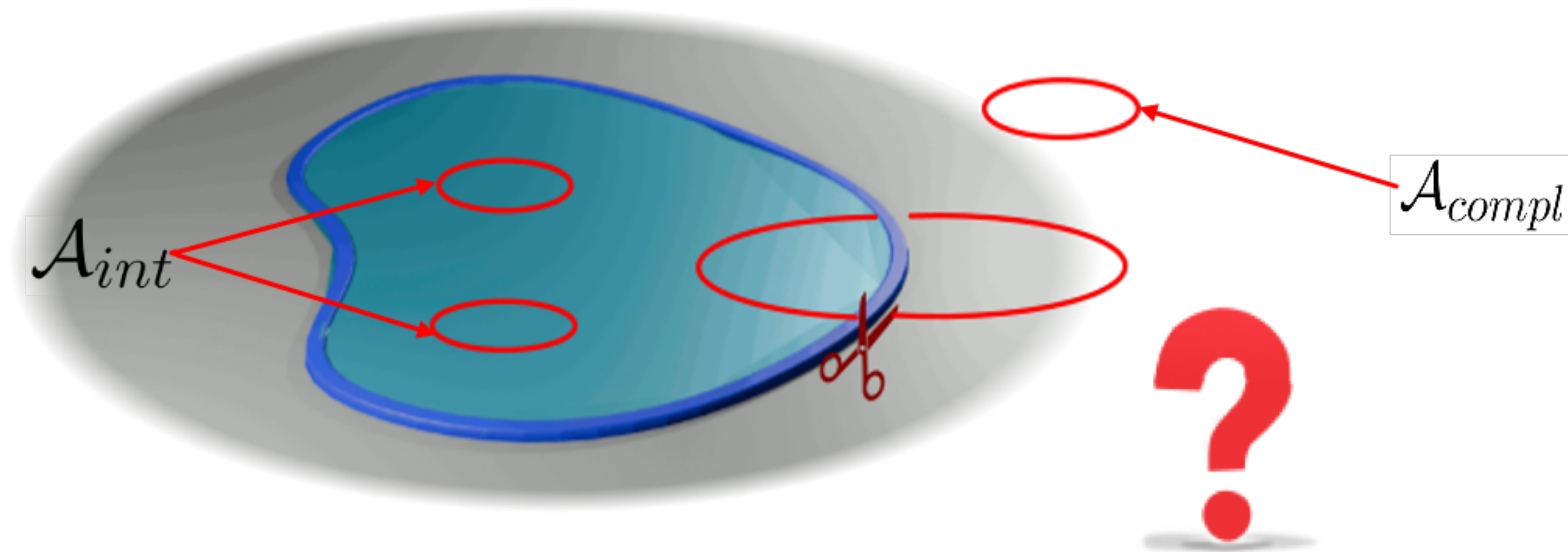
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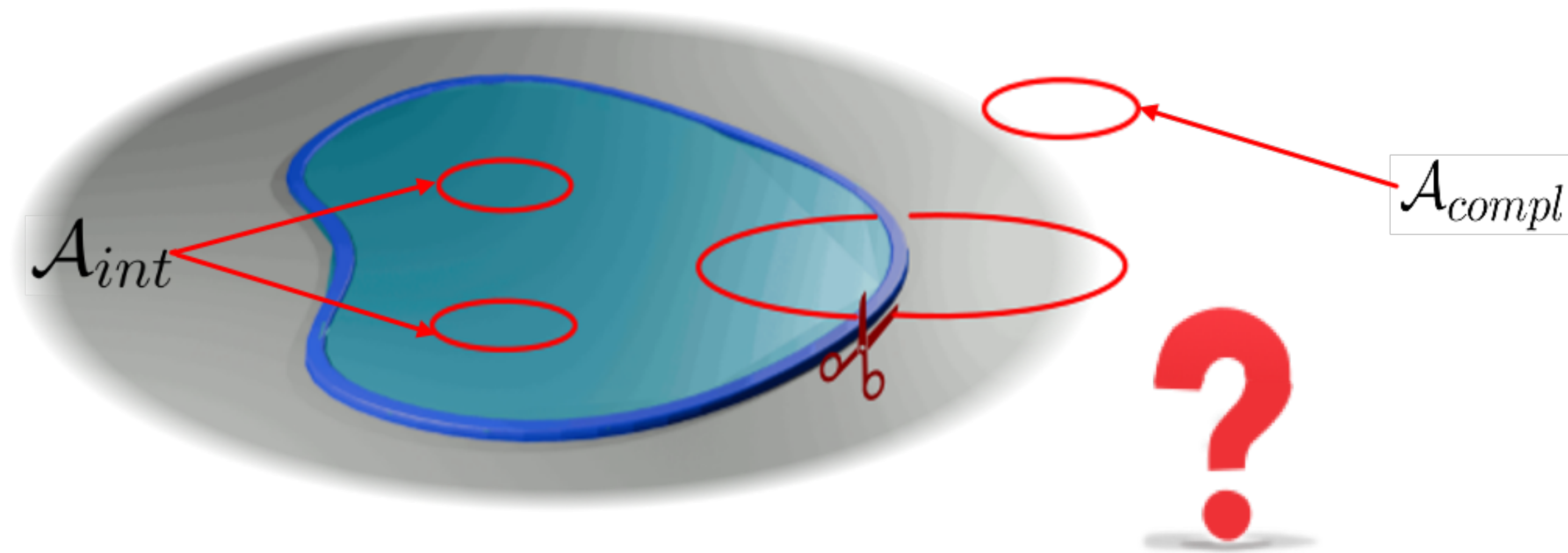


typically get center algebras: $\mathcal{A}_{\text{phys}} \supset \bigoplus_q \mathcal{A}_{\text{int}}^q \otimes \mathcal{A}_{\text{compl}}^q$

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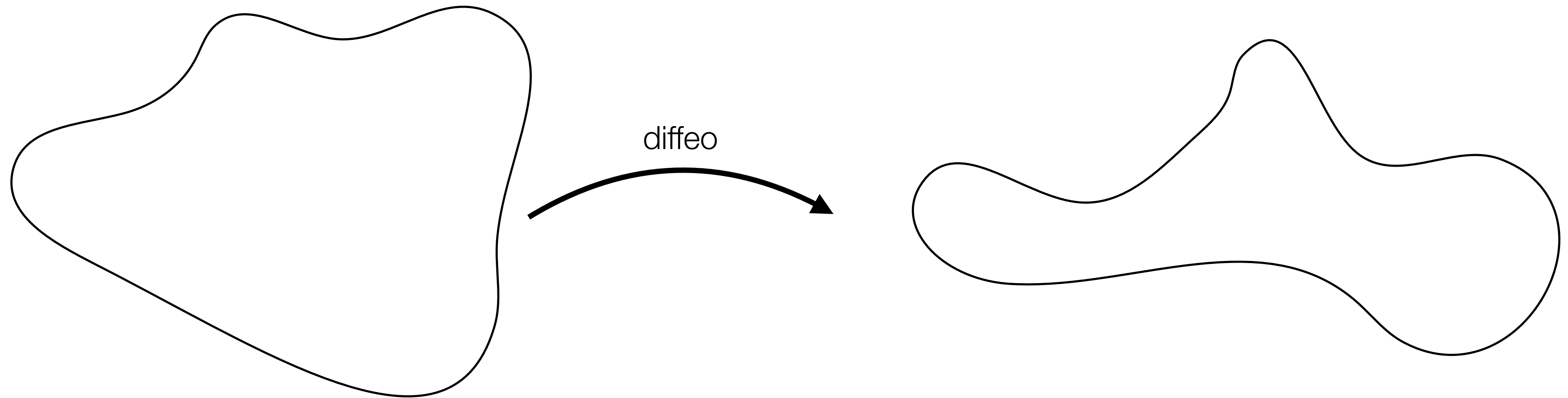
\Rightarrow vN entropy non-distillable (doesn't measure entanglement)

$$S_{\text{vN}}(\rho_{\text{int}}) = \sum_q p_q S_{\text{vN}}(\rho^q) + H(\{p_q\}) + \langle \log d_q \rangle$$

Shannon \swarrow

Subsystems in gravity

⇒ situation aggravated: gauge diffeos move naive definitions of subregions around

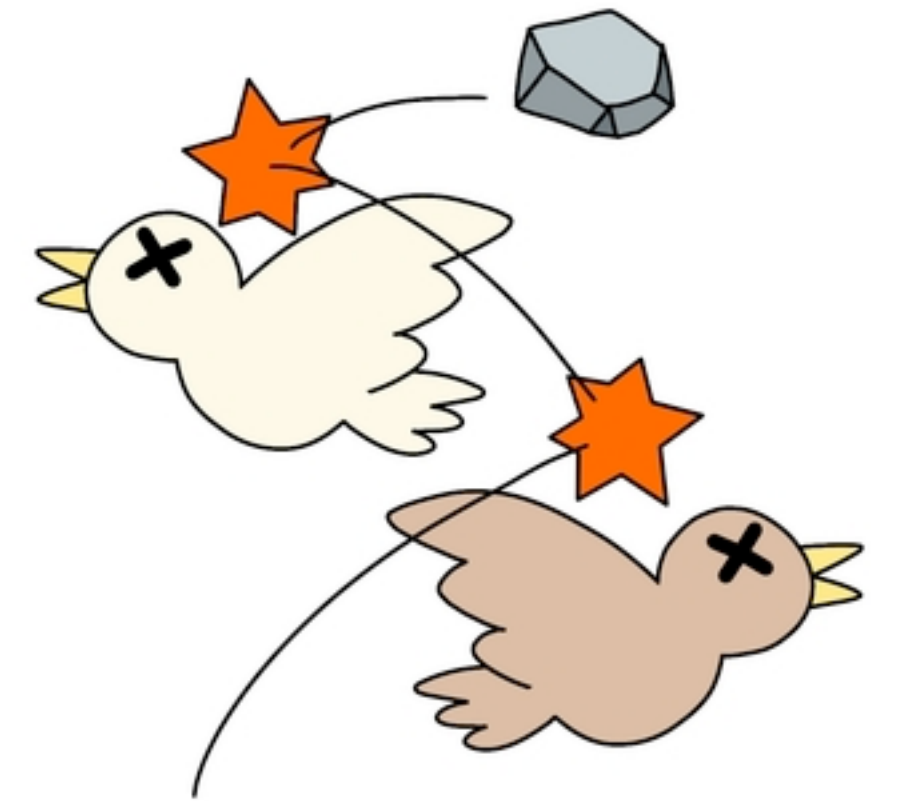


+ also need to worry about what is “in” and “out”

Quantum reference frames ...

... offer a new perspective on both challenges

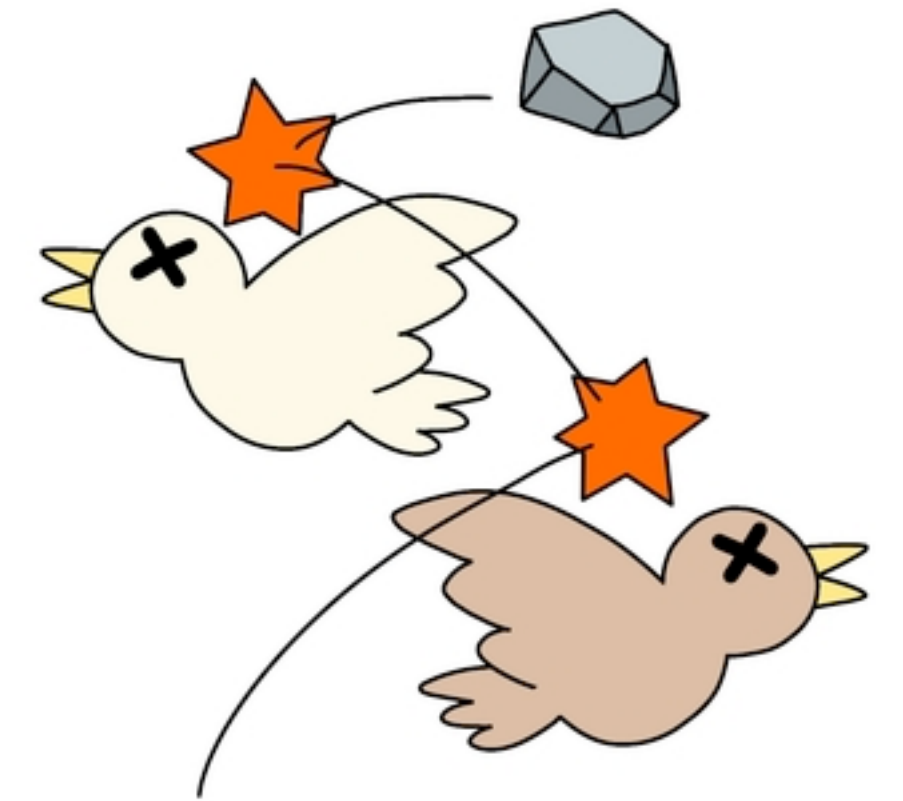
- **intrinsic regularization when describing regional data relative to a QRF** (type conversion of the corresp. vN algebra)
- **QRF perspective \simeq gauge-inv. subsystem decomposition**
 \Rightarrow **many distinct definitions of subsystem associated with same subregion**



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- **intrinsic regularization when describing regional data relative to a QRF** (type conversion of the corresp. vN algebra)
- **QRF perspective \simeq gauge-inv. subsystem decomposition**
 - \Rightarrow many distinct definitions of subsystem associated with same subregion
 - \Rightarrow embrace that diversity:
entropy and algebra hierarchy + some factor algebras and distillable definitions



Quantum reference frames

[PH, Kotecha, Mele '13; de la Hamette, Galley, PH, Loveridge, Müller '21; De Vuyst, Eccles, PH, Kirklin '24; Araújo-Regado, PH, Sartini '25;...]

Quantum reference frames and crossed products

[PH, Kotecha, Mele '13; de la Hamette, Galley, PH, Loveridge, Müller '21; De Vuyst, Eccles, PH, Kirklin '24; Araújo-Regado, PH, Sartini '25;...]

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e.g., tetrad valued in the Lorentz group, particle position translation-group-valued,...

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- **nonideal QRF:** fuzzy orientations

$$\text{e.g. } \mathcal{H}_R \subsetneq L^2(G)$$

Physics from the inside

$$\mathcal{H}_R$$

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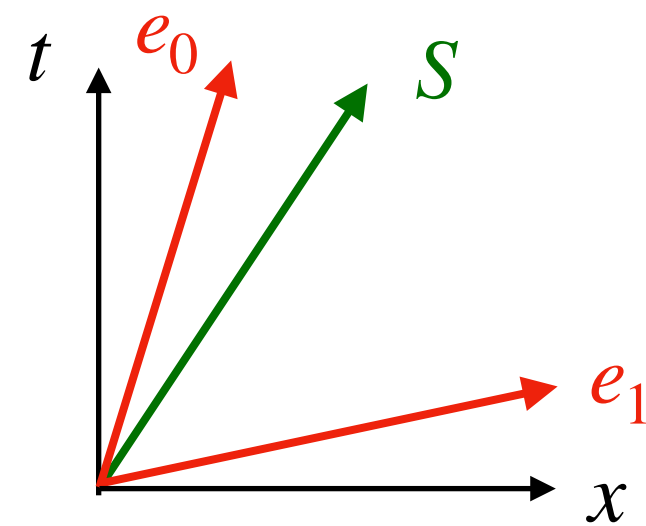
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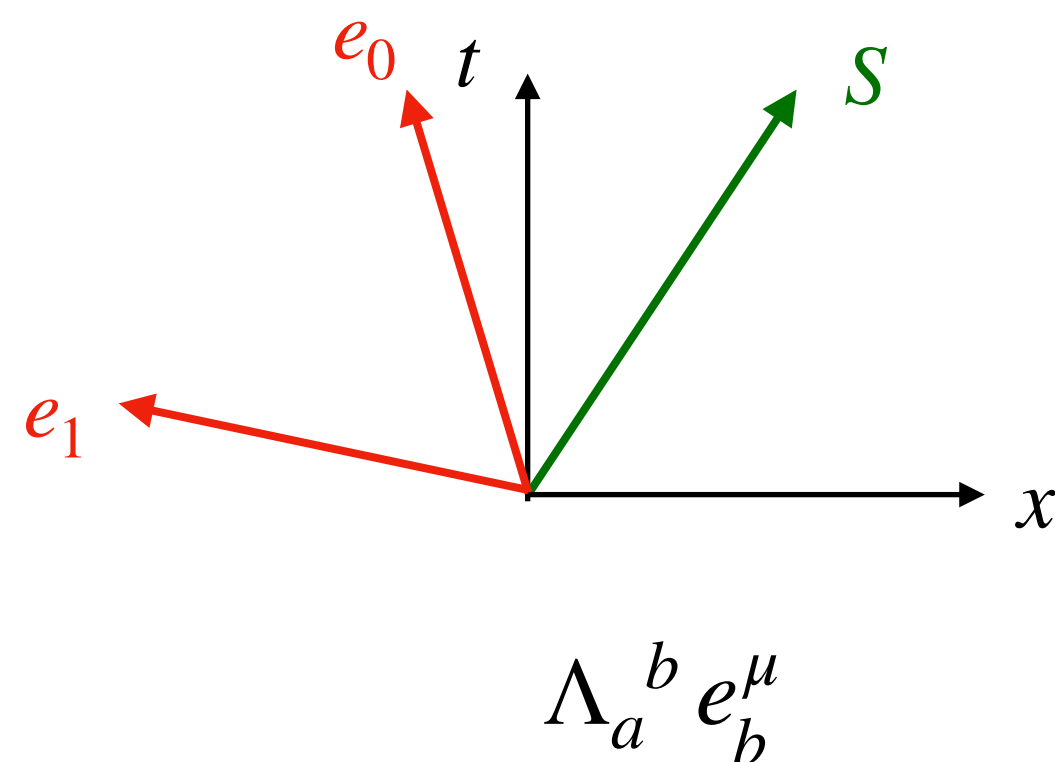
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measures f_S conditional on R being in orientation g

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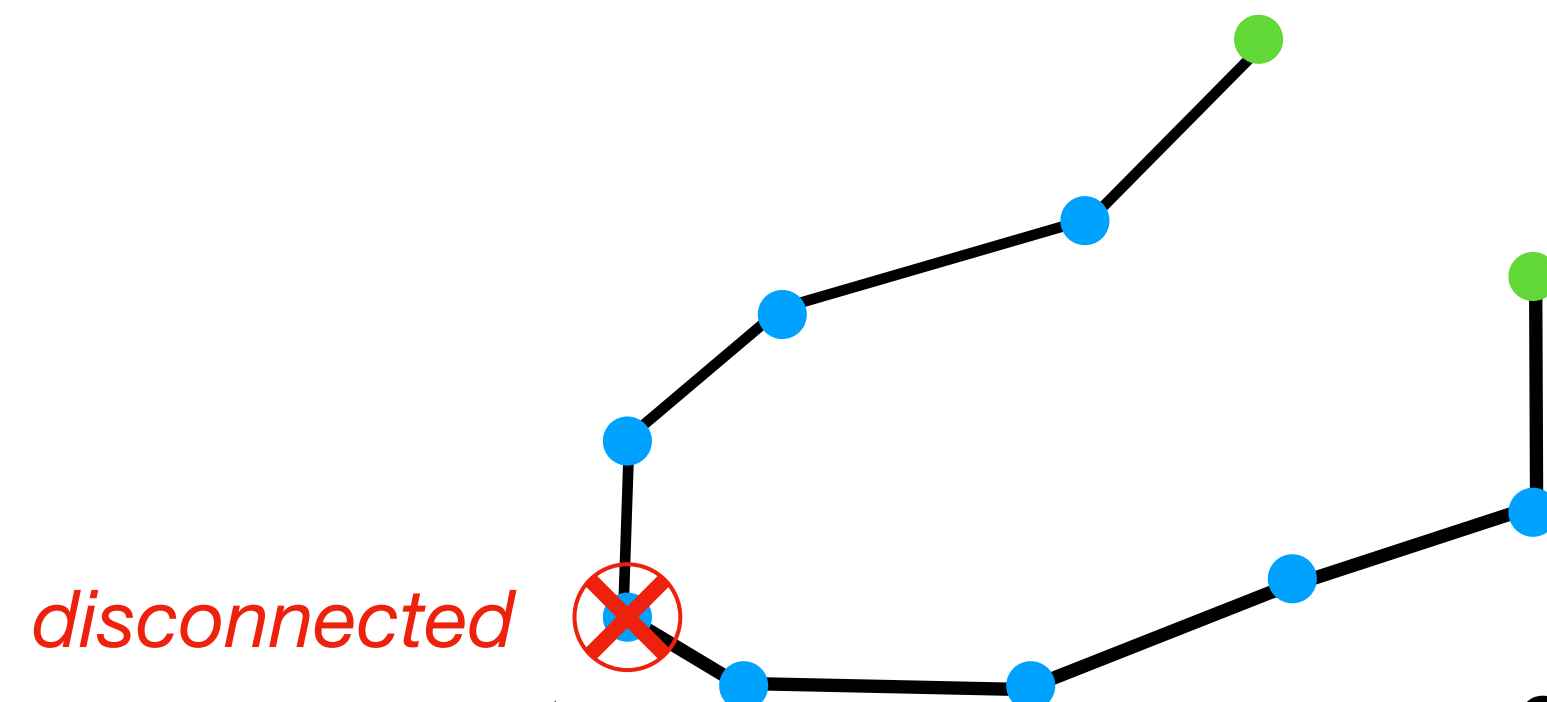
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2. lattice gauge theory



some root vertex in the lattice

2 Wilson lines define G -valued QRF R at green nodes

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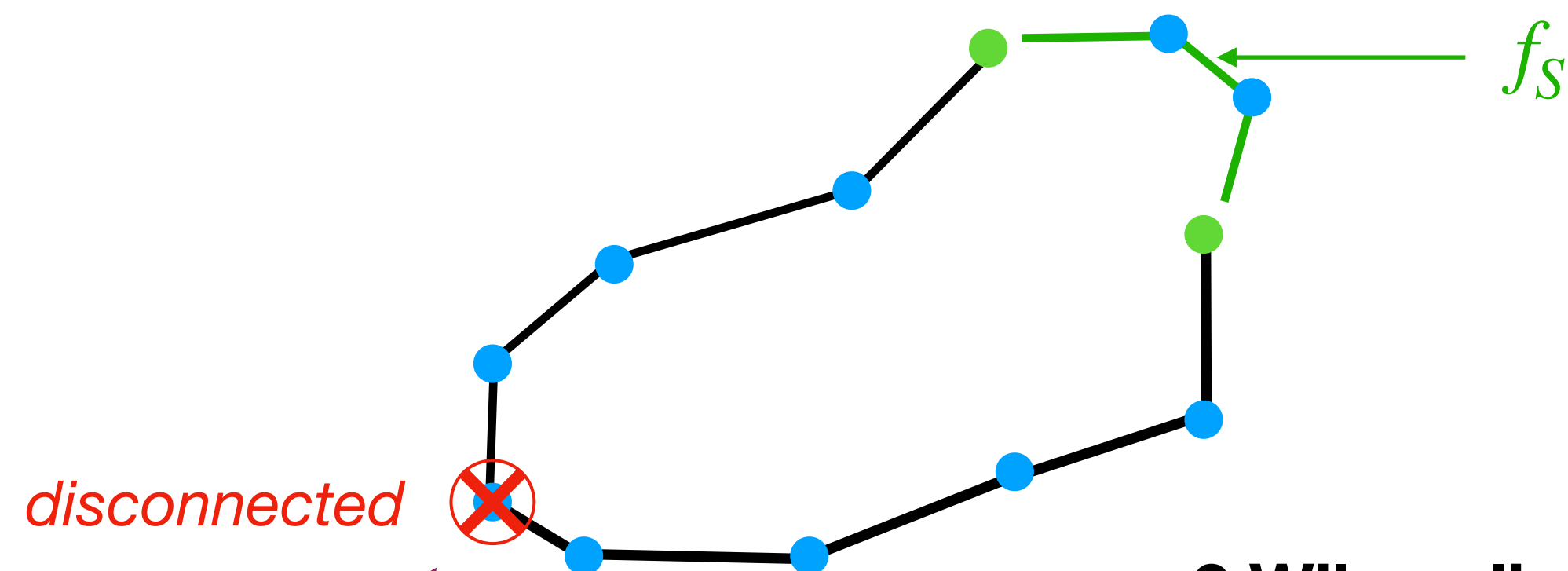
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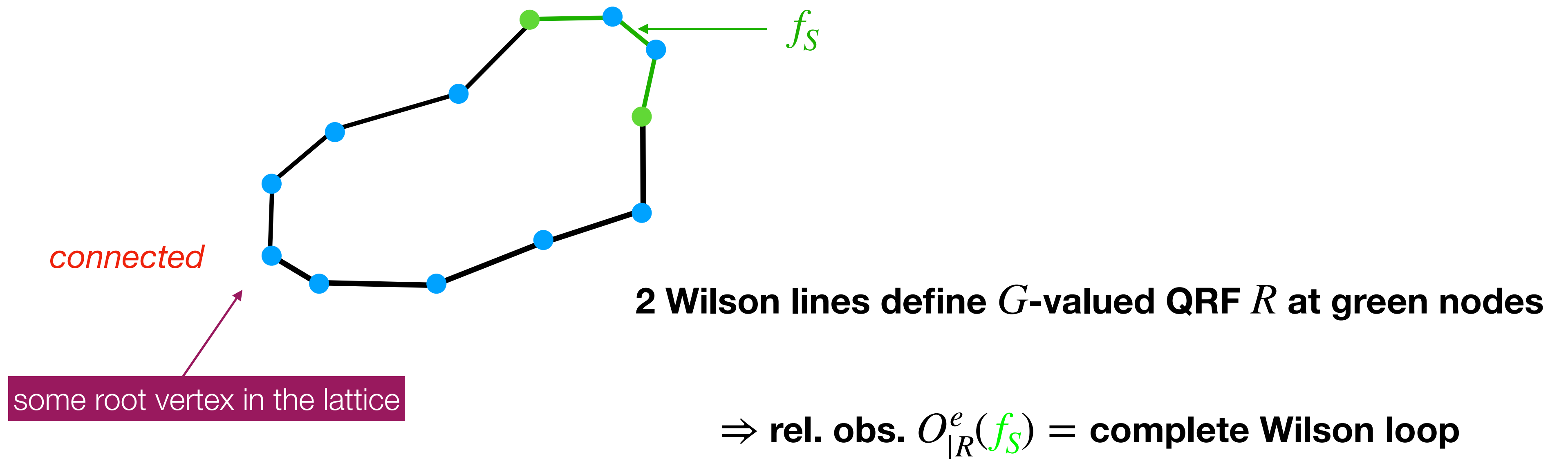
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Crossed products

Crossed product of system algebra \mathcal{A}_S with reorientations of an **ideal QRF**

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so far: everything on $\mathcal{H}_{\text{kin}} = \mathcal{H}_R \otimes \mathcal{H}_S$

⇒ to compute entropies, proceed to physical Hilbert space: $\mathcal{H}_{\text{phys}} = \Pi(\mathcal{H}_{\text{kin}})$

$$\Pi = \int_G dg U_R(g) \otimes U_S(g)$$

Subsystem relativity

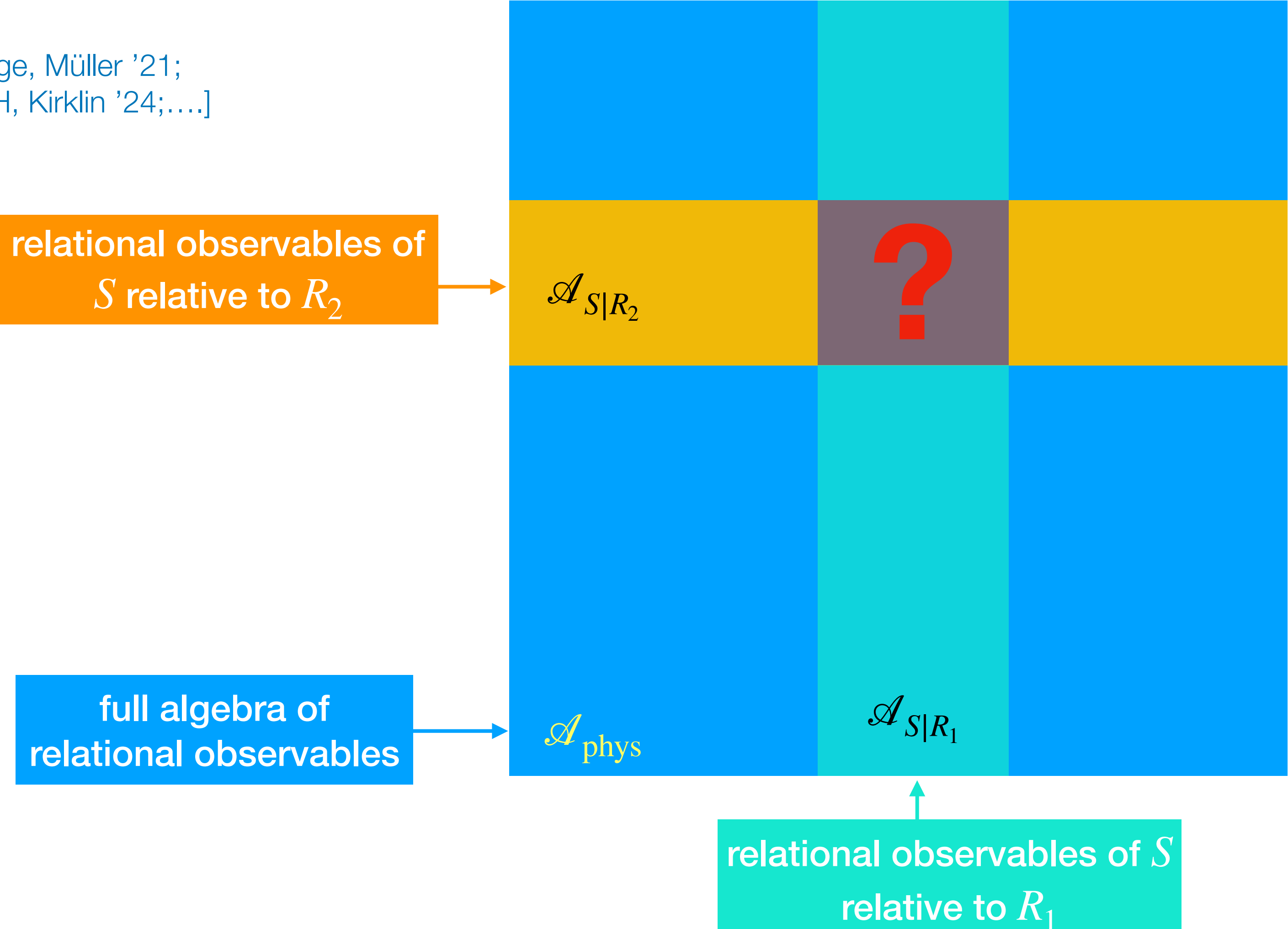
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[Ahmad, Galley, PH, Lock, Smith '21; de la Hamette, Galley, PH, Loveridge, Müller '21; Castro-Ruiz, Oreshkov '21; PH, Kotecha, Mele '23; De Vuyst, Eccles, PH, Kirklin '24;...]

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relational observables of S relative to R_2

$\mathcal{A}_{S|R_2}$

all rel. observables describing S that are inv. under *both* R_1 - & R_2 -reorientations \Rightarrow all internal S -relations

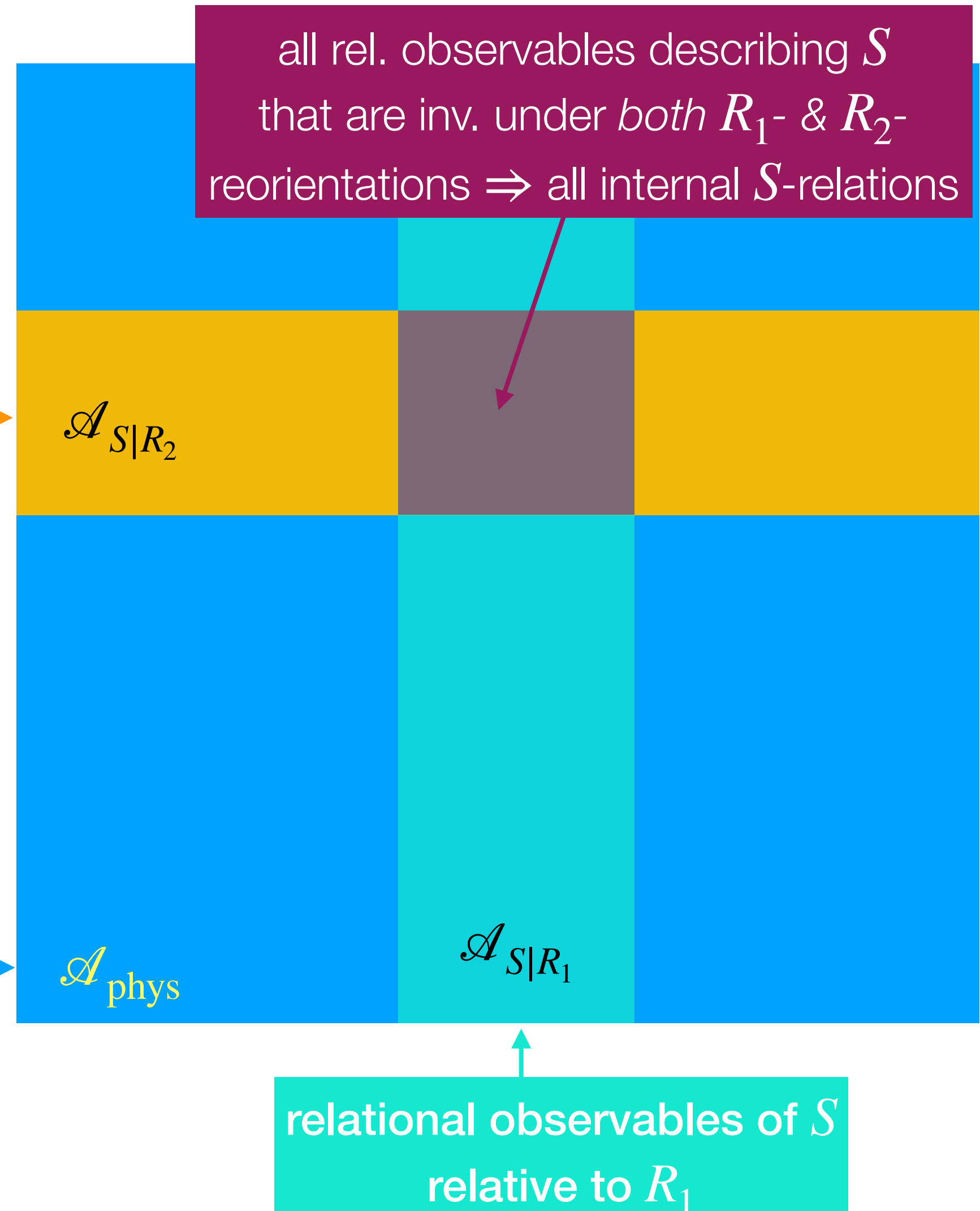
“frames R_1 and R_2 mean different gauge inv. DoFs when they refer to subsystem S ”

full algebra of relational observables

$\mathcal{A}_{\text{phys}}$

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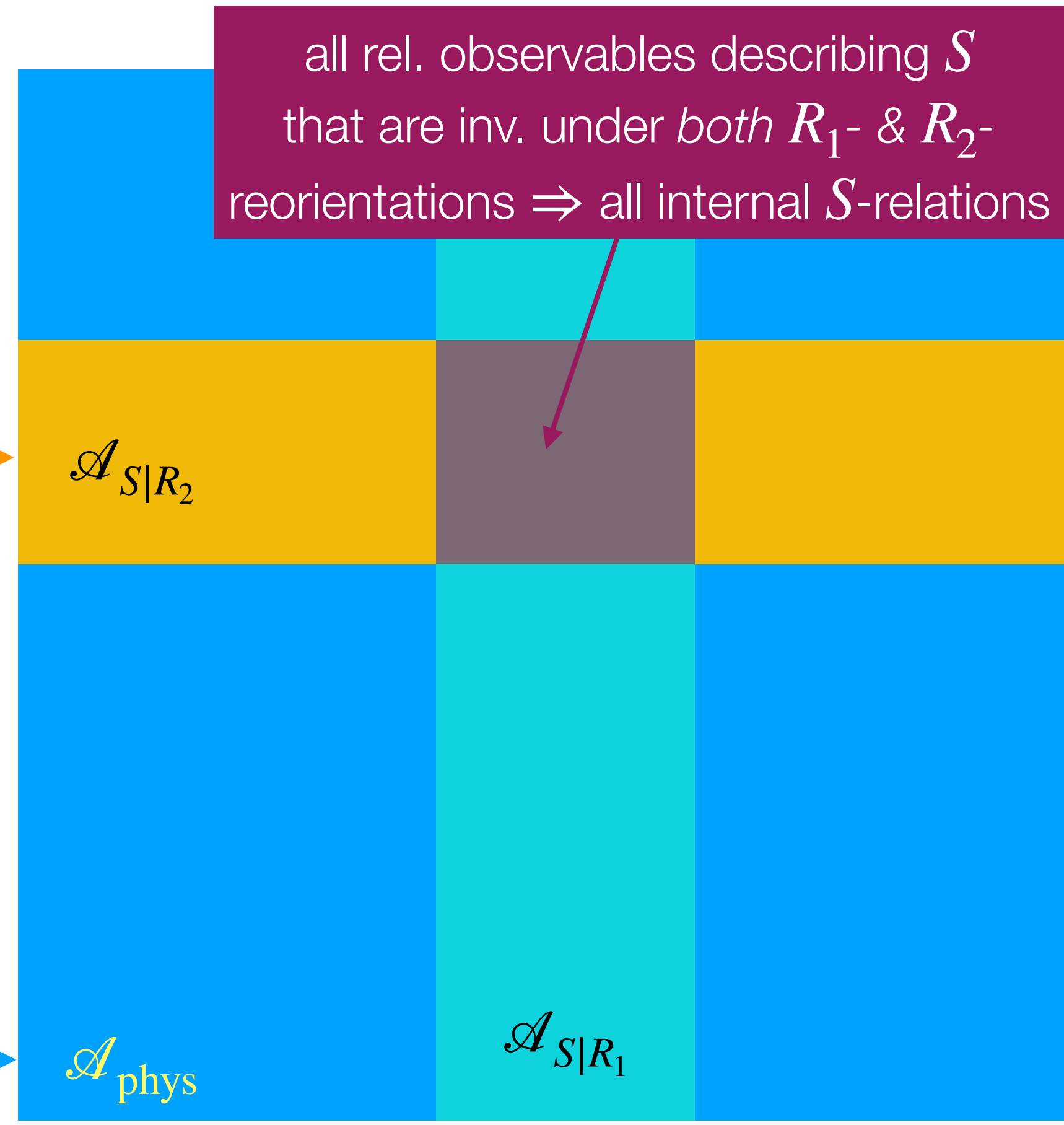
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full algebra of relational observables

⇒ correlations, entropies, thermodynamic properties, ... become QRF-dependent

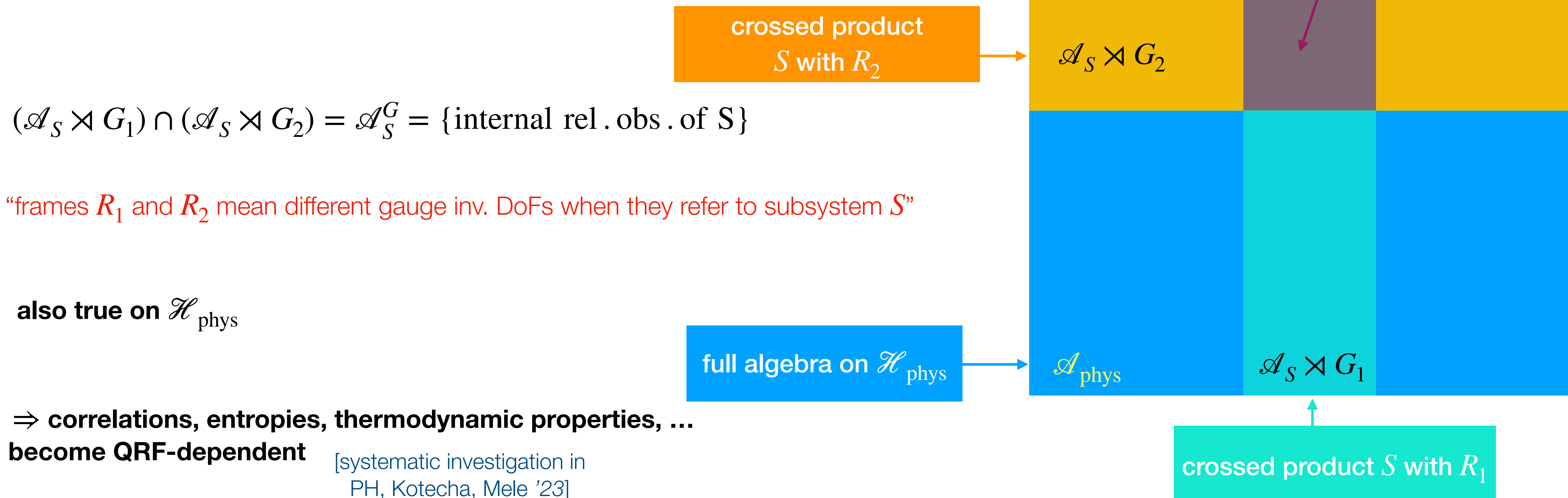
[systematic investigation in PH, Kotecha, Mele '23]

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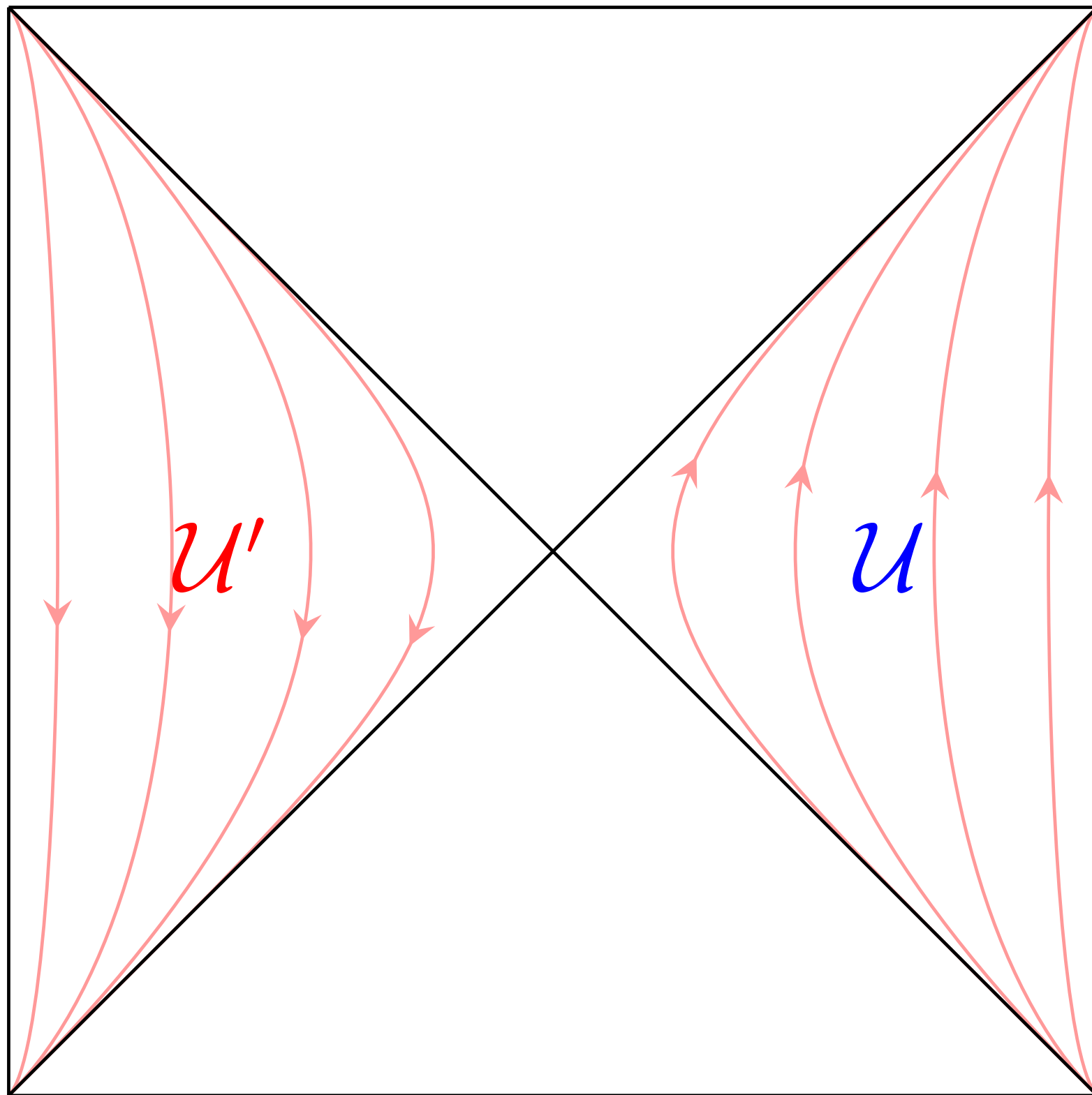
Relational entanglement entropies in perturbative QG

[Chandrasekaran, Longo, Penington, Witten '22; Jensen, Sorce, Speranza '23; Kudler-Flam, Leutheusser, Satishchandran '23; ...
De Vuyst, Eccles, PH, Kirklín '24 I & II]

Regional algebras in perturbative QG

consider two static patches in dS^* filled with matter and gravitons

\Rightarrow what is the EE of the right patch \mathcal{U} in some global state?



* works the same way for any Killing horizon

Regional algebras in perturbative QG

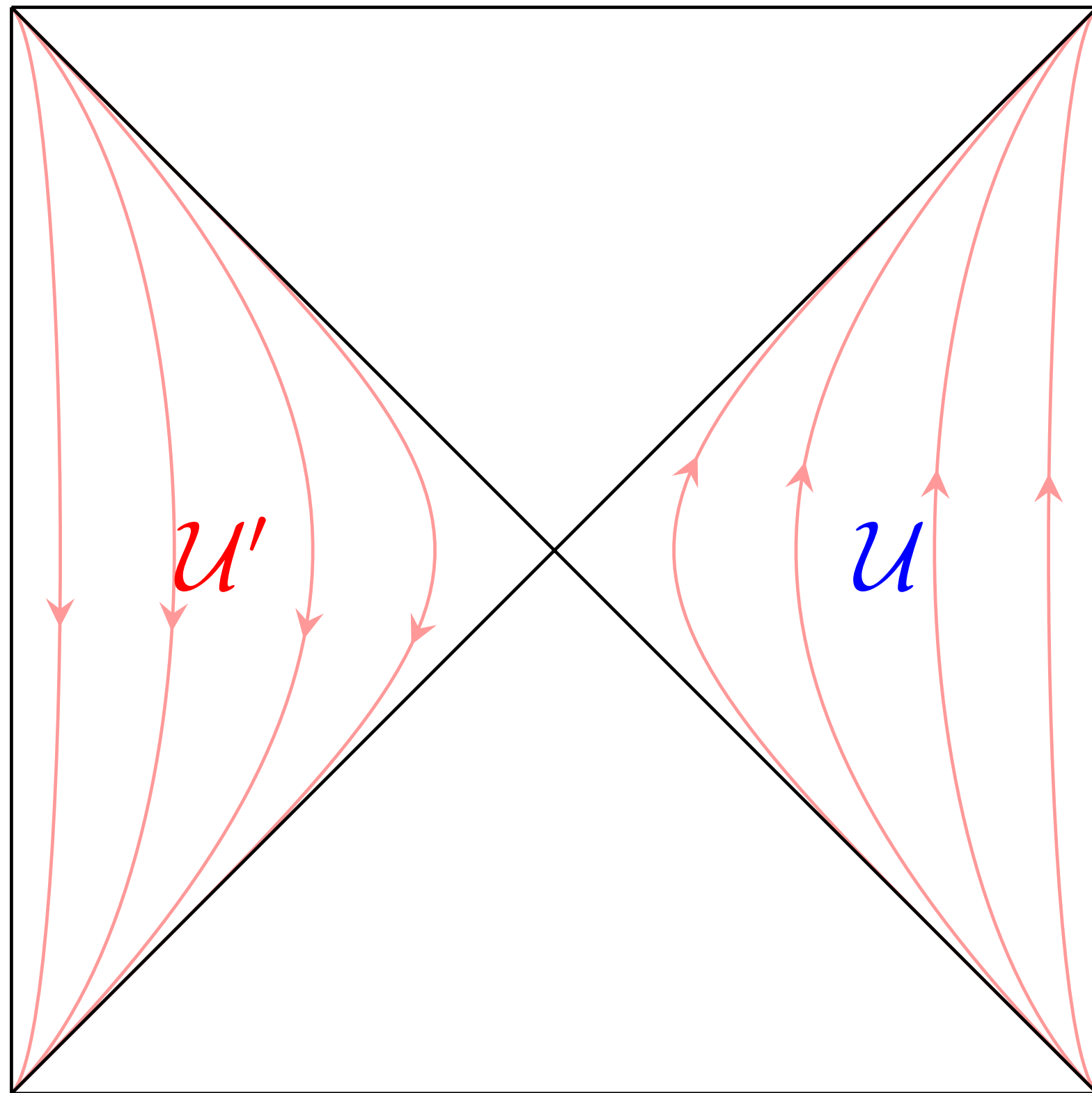
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2 challenges:

- regional algebras $\mathcal{A}_{\mathcal{U}}, \mathcal{A}_{\mathcal{U}'}$ are type III factors [Araki '64; Longo '82; Fredenhagen '85]

\Rightarrow no trace, no density operators, no entropies



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The need for “observers”

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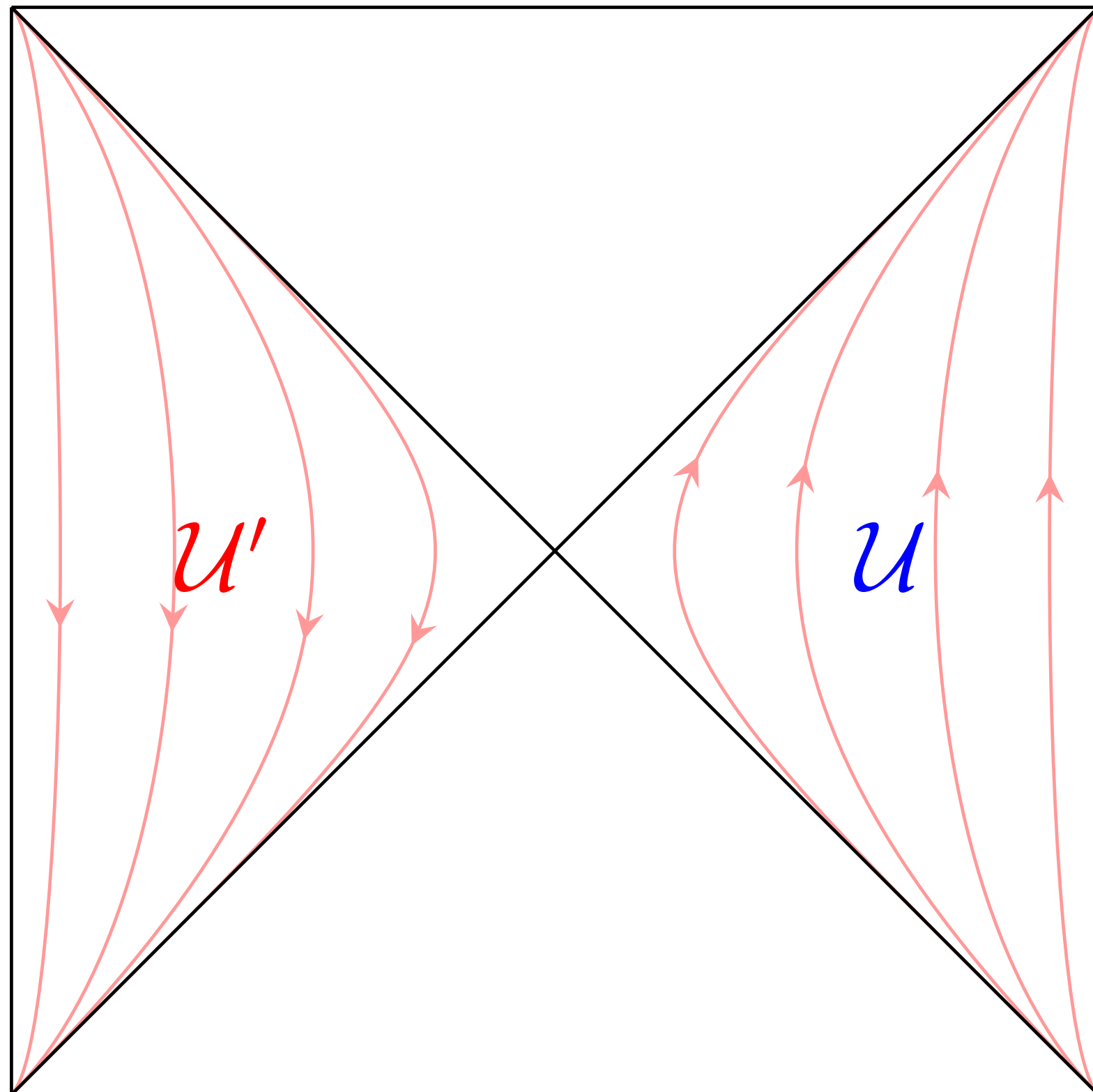
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- isometries gauge in dS , so impose boost Hamiltonian H (mod. Ham. of the vacuum) as gauge constraint [Chandrasekaran et al. '22]

⇒ acts ergodically on $\mathcal{A}_{\mathcal{U}}, \mathcal{A}_{\mathcal{U}'}$

$$\mathcal{A}_{\mathcal{U}}^H = \mathbb{C}$$

⇒ no interesting **regional** diff-inv. observables (need additional DoFs)



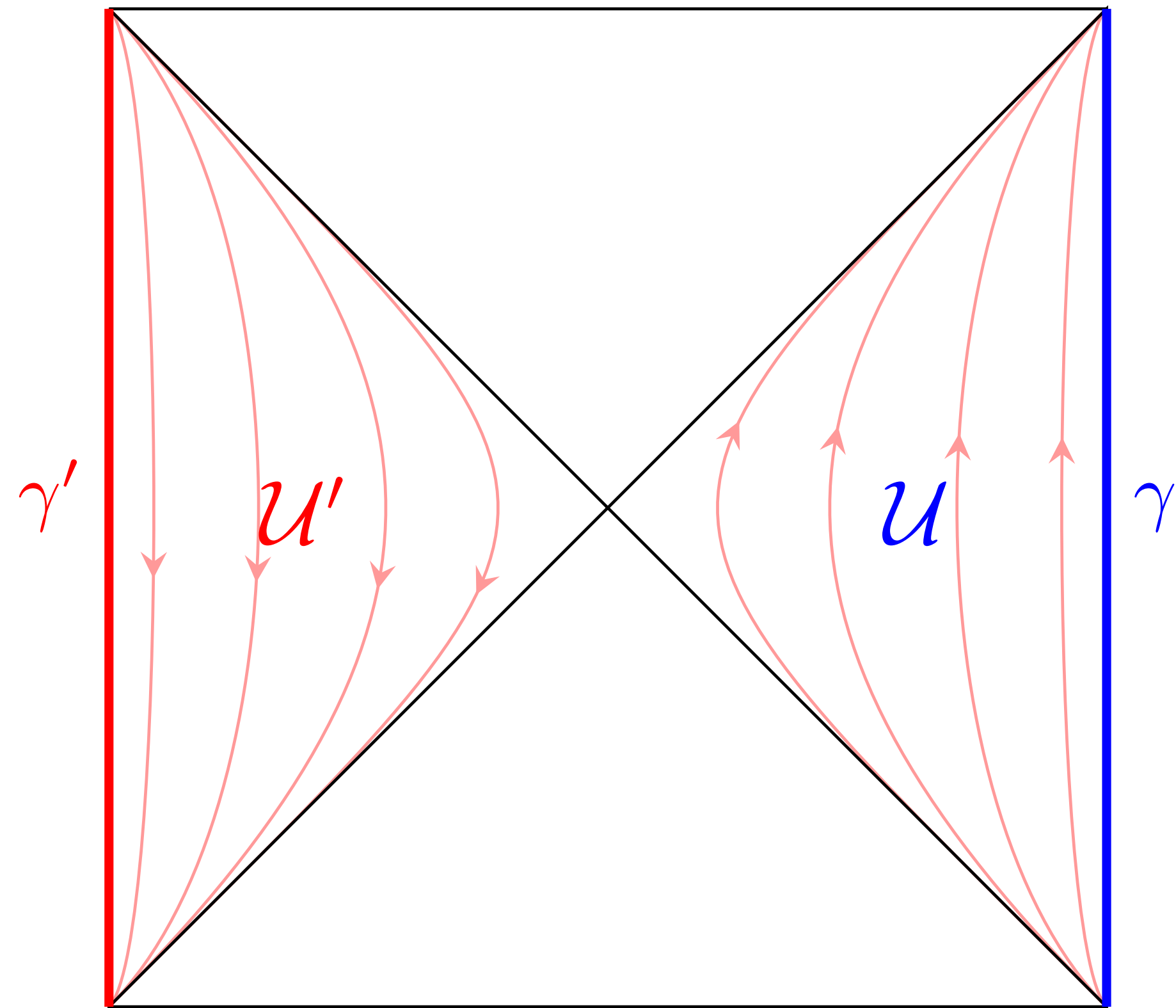
Observers are QRFs

[Chandrasekaran, Longo, Penington, Witten (CLPW) '22]

remedy: introduce observers with **ideal** clock in each static patch

$$C = H + p_1 - p_2$$

(perturbative QG $G_N \rightarrow 0$)



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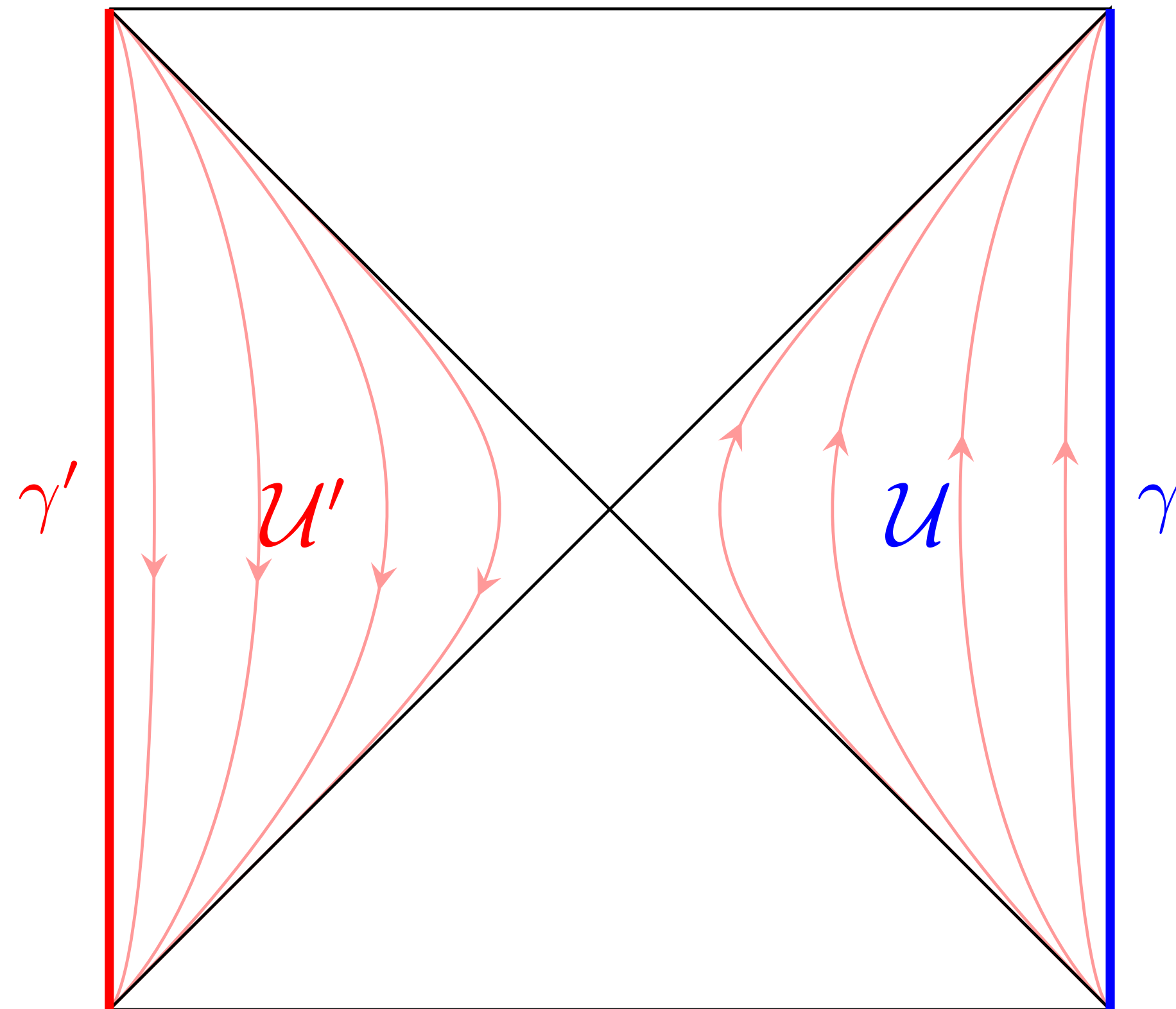
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inv. regional algebra becomes non-trivial crossed product

$$\mathcal{A}_{\mathcal{U}}^H \longrightarrow (\mathcal{A}_{\mathcal{U}} \otimes \mathcal{B}(L^2(\mathbb{R}))^C = \langle e^{iTH} a e^{-iTH}, e^{-isp_1} \mid a \in \mathcal{A}_{\mathcal{U}}, s \in \mathbb{R} \rangle''$$

crossed product = relational observables + clock reorientations



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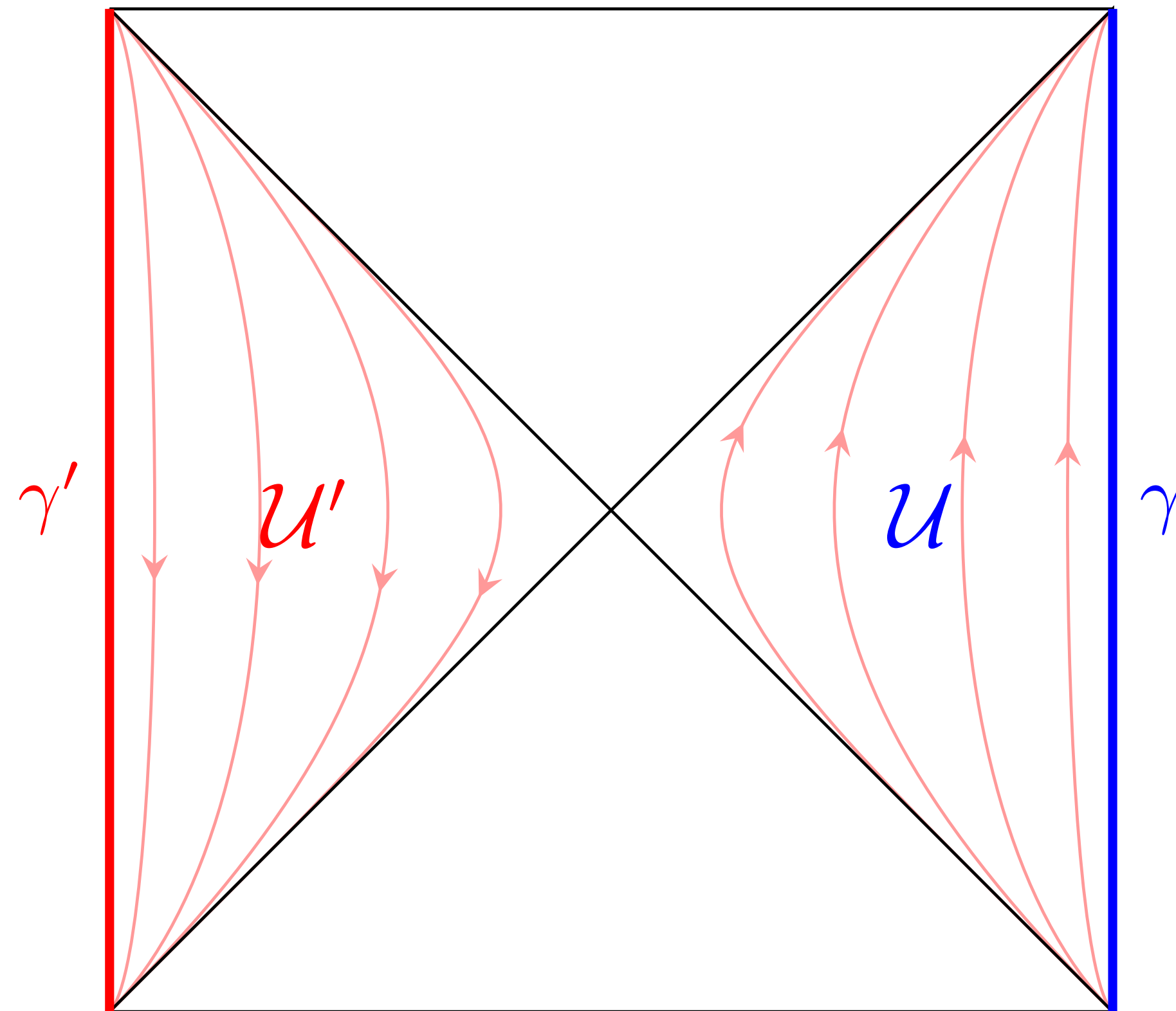
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$$\mathcal{A}_{\mathcal{U}}^H \longrightarrow (\mathcal{A}_{\mathcal{U}} \otimes \mathcal{B}(L^2(\mathbb{R}))^C = \langle e^{iTH} a e^{-iTH}, e^{-isp_1} \mid a \in \mathcal{A}_{\mathcal{U}}, s \in \mathbb{R} \rangle''$$

crossed product = relational observables + clock reorientations

crux:

regional QFT (incl. graviton) algebra $\mathcal{A}_{\mathcal{U}}$ is a type III factor, but crossed product is a type II factor (possesses trace, density operators and entropies)



Observers are QRFs

[Chandrasekaran, Longo, Penington, Witten (CLPW) '22]

remedy: introduce observers with **ideal** clock in each static patch

$$C = H + p_1 - p_2 \quad (\text{perturbative QG } G_N \rightarrow 0)$$

inv. regional algebra becomes non-trivial crossed product

$$\mathcal{A}_{\mathcal{U}}^H \longrightarrow (\mathcal{A}_{\mathcal{U}} \otimes \mathcal{B}(L^2(\mathbb{R}))^C = \langle e^{iTH} a e^{-iTH}, e^{-isp_1} \mid a \in \mathcal{A}_{\mathcal{U}}, s \in \mathbb{R} \rangle''$$

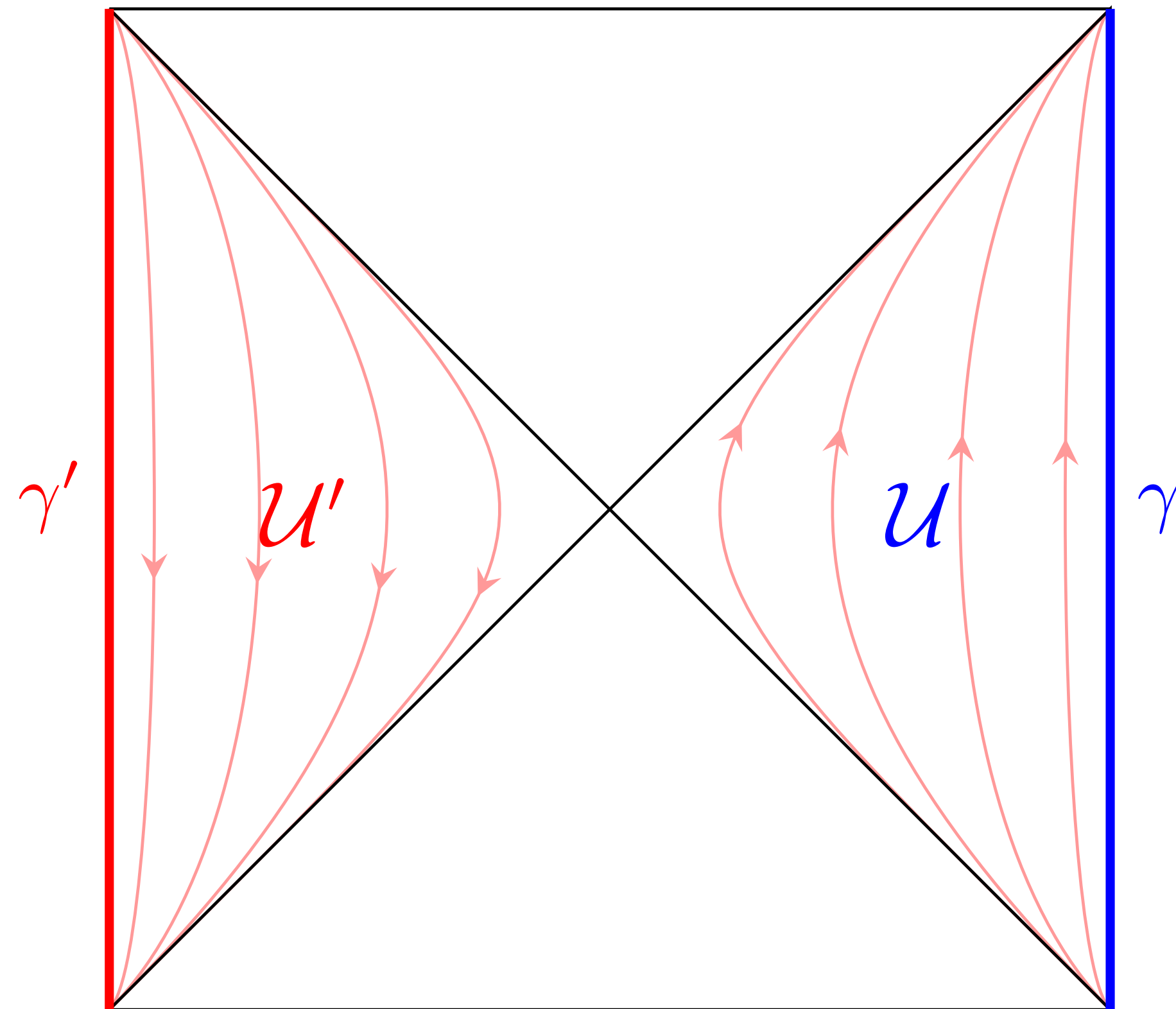
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crux:

regional QFT (incl. graviton) algebra $\mathcal{A}_{\mathcal{U}}$ is a type III factor, but crossed product is a type II factor (possesses trace, density operators and entropies)

\Rightarrow description **relative to clock QRF** yields well-defined EE (“intrinsic regularization”)

(better behaved for non-ideal clocks)



Density operators and entropy

[CLPW '22; Jensen, Sorce, Speranza '23; ...]

- **type II, so trace, density operators and entropies exist**

⇒ here: crossed product is a **factor**, so grav. EE distillable

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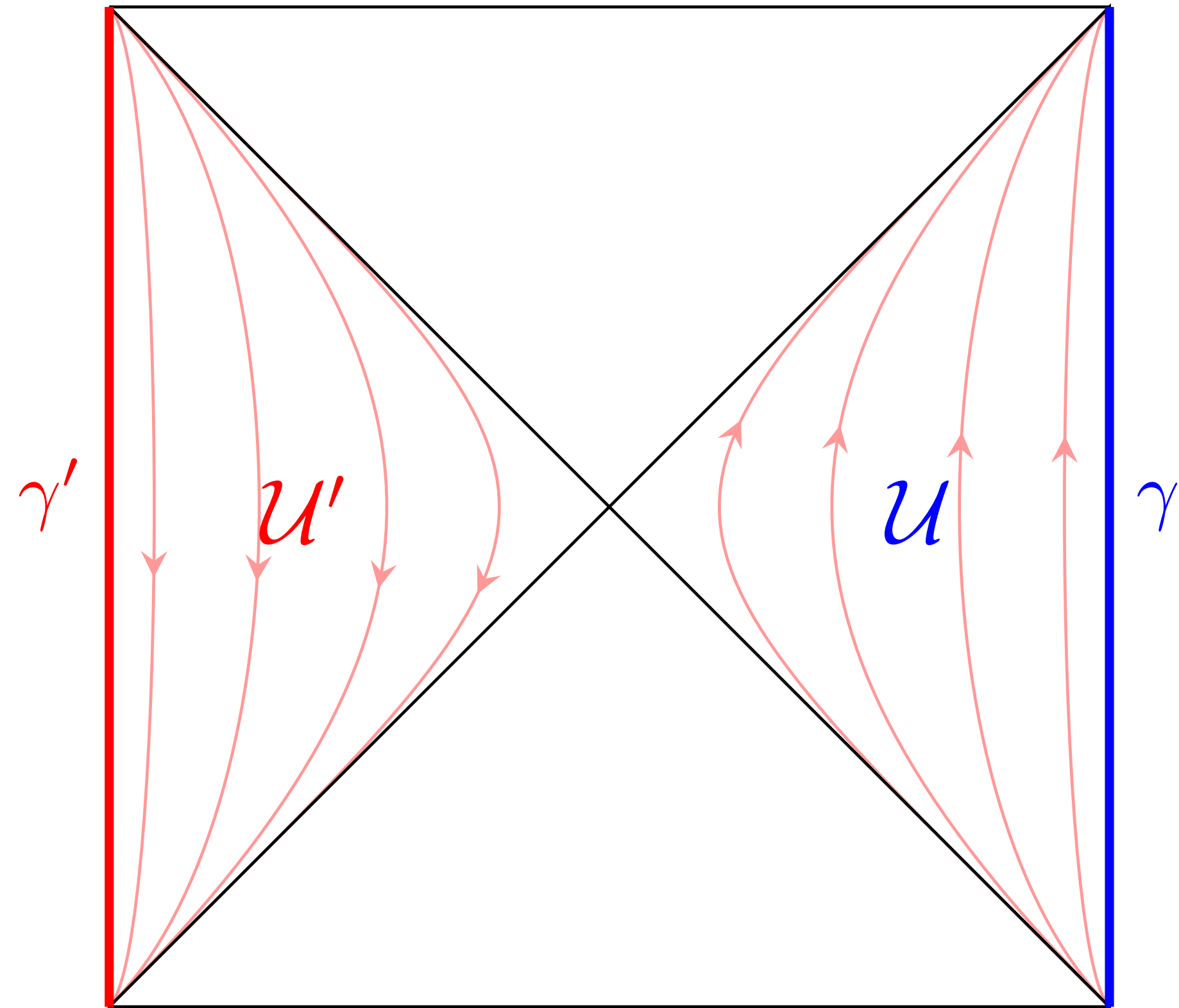
⇒ here: crossed product is a **factor**, so grav. EE distillable

⇒ for certain states reproduces generalized entropy

$$S_{\mathcal{U}} = \frac{A_{\partial\mathcal{U}}}{4G} + S_{\text{QFT}}$$

These gravitational EEs are relational

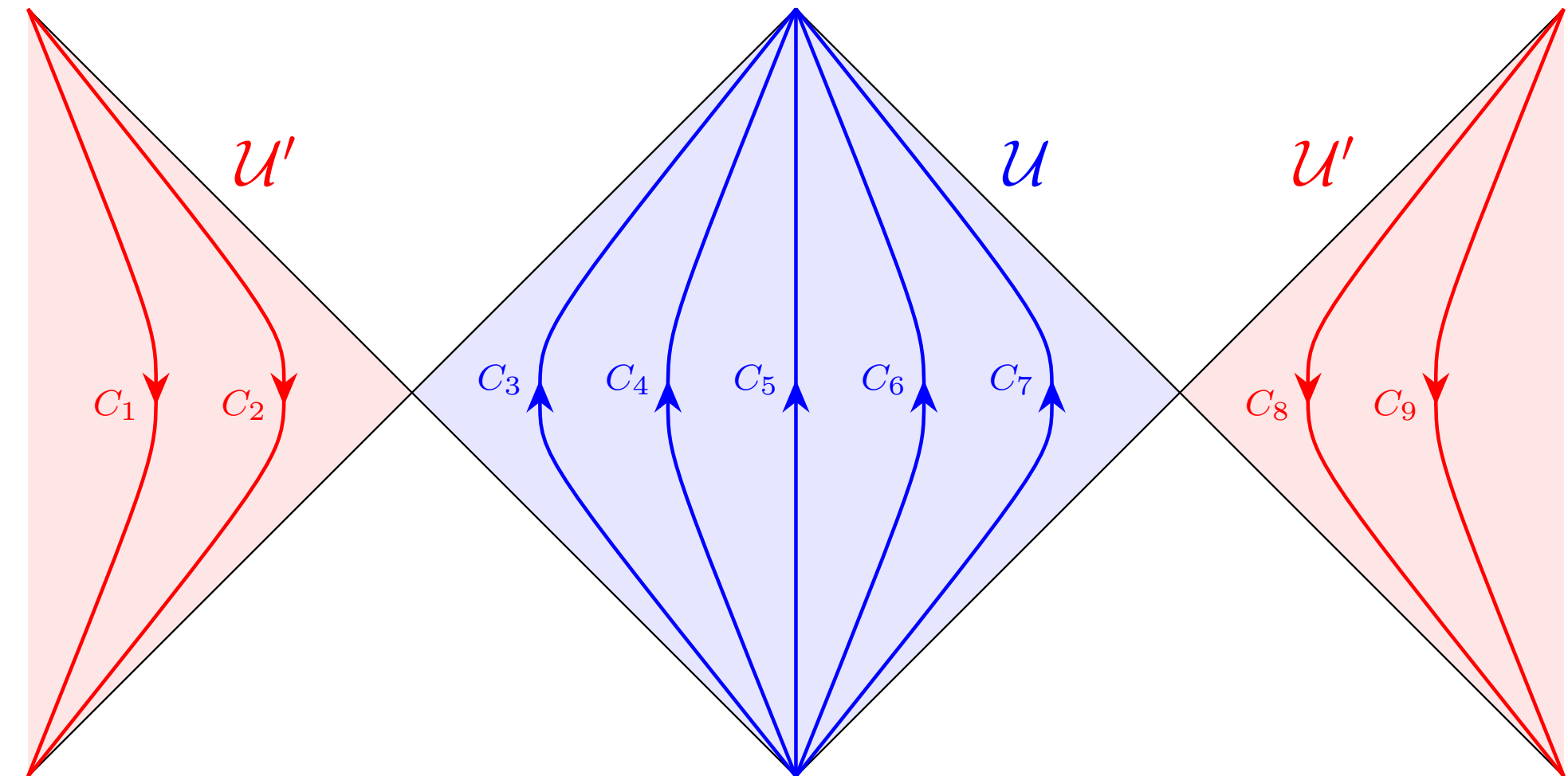
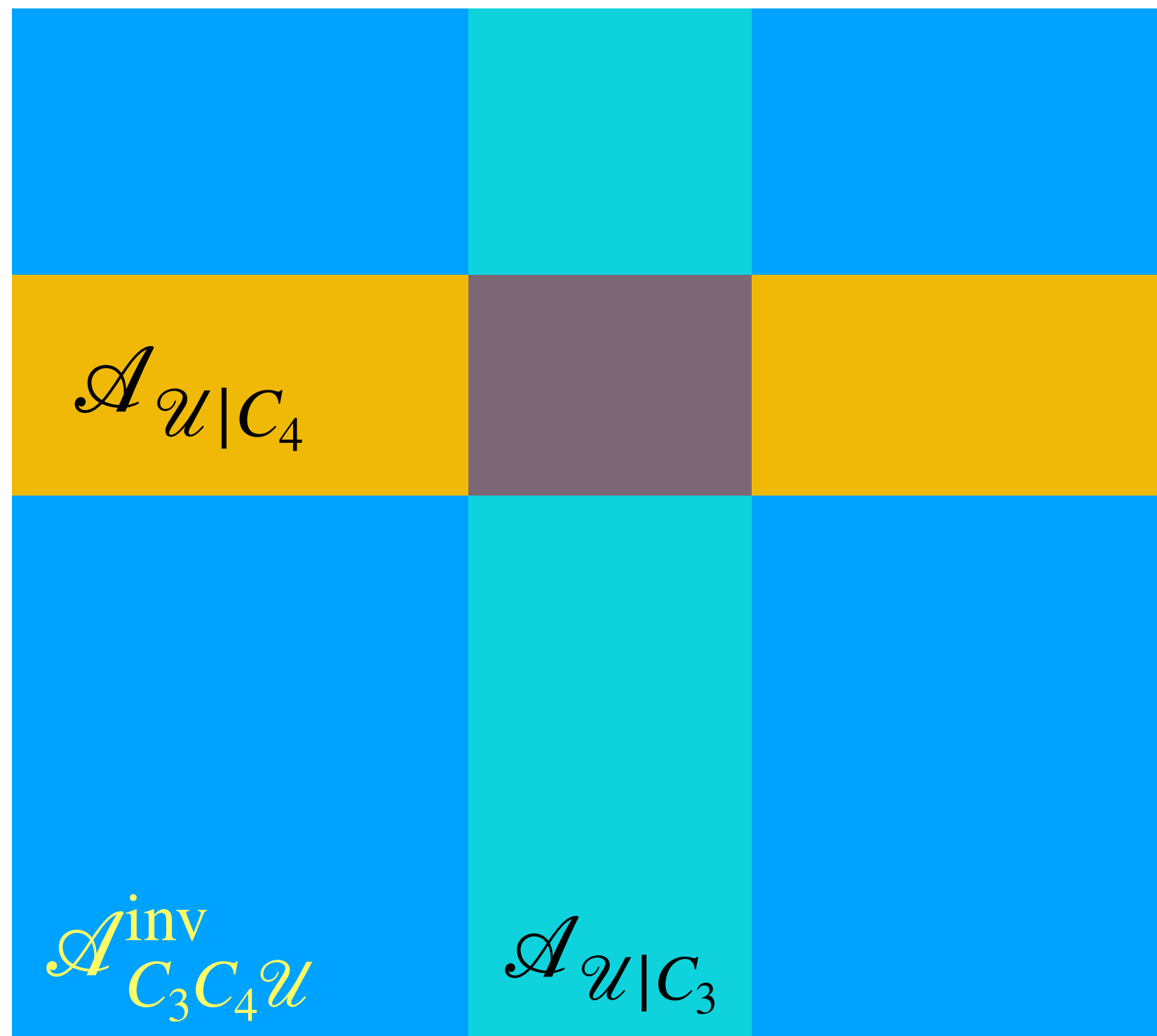
... describing regional QFT DoFs relative to chosen QRF/observer



Gravitational entropy is observer-dependent

- populating spacetime with arbitrarily many observers/QRFs

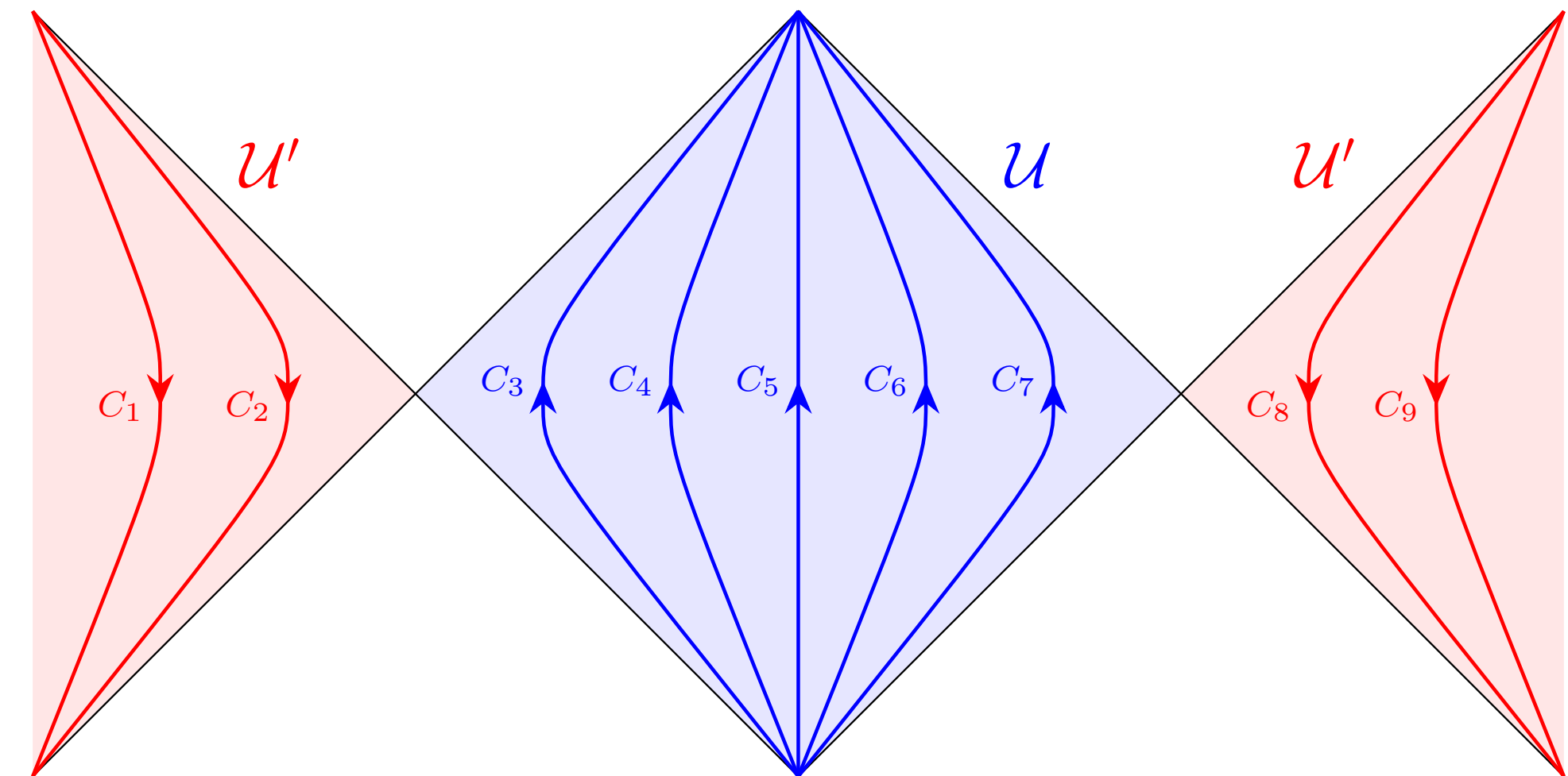
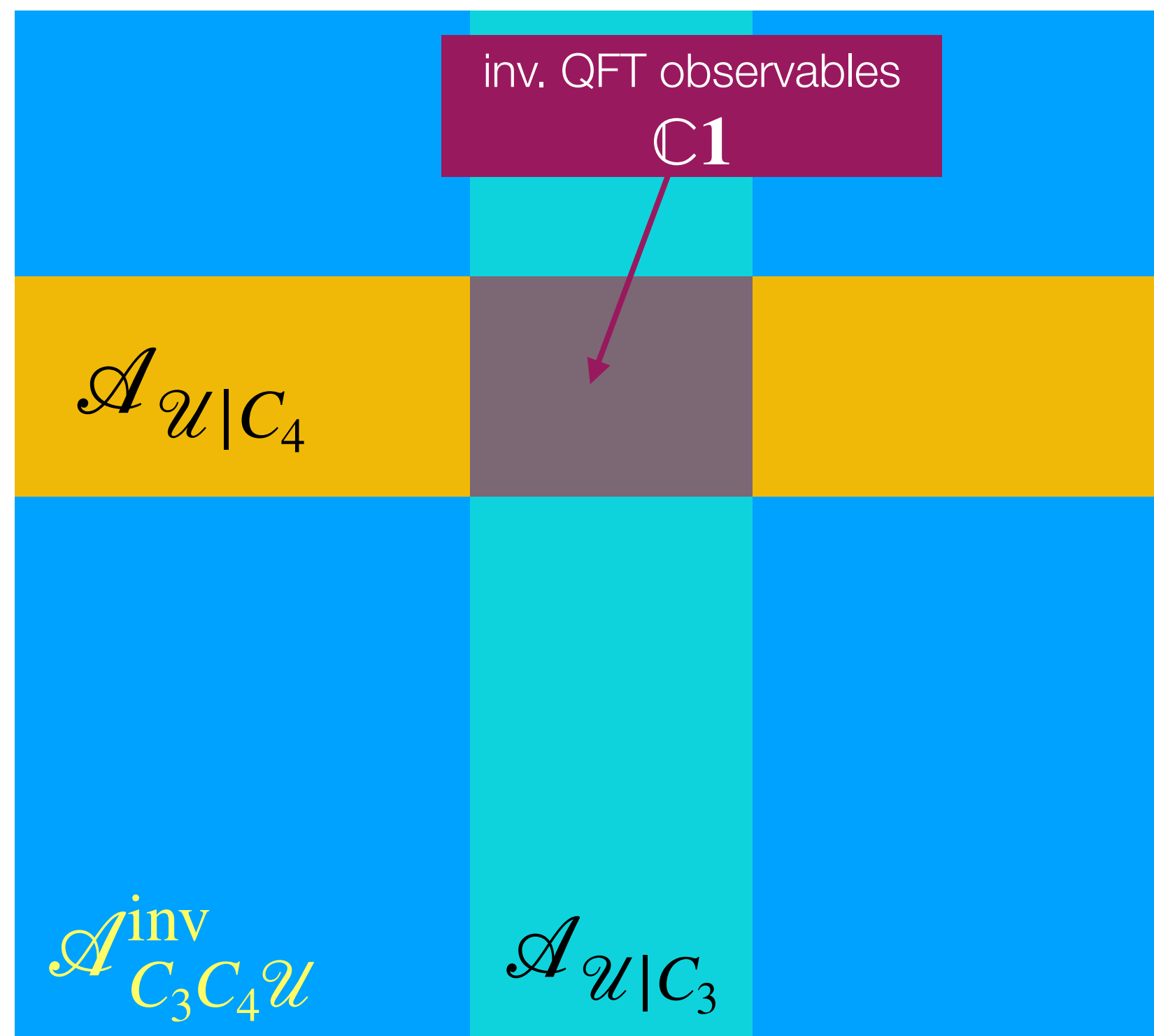
kinematical QFT (incl. graviton) subsystem \mathcal{A}_U



Gravitational entropy is observer-dependent

- populating spacetime with arbitrarily many observers/QRFs

kinematical QFT (incl. graviton) subsystem $\mathcal{A}_{\mathcal{U}}$



\Rightarrow full regional/horizon entropy invariant,
but QFT contributions are not

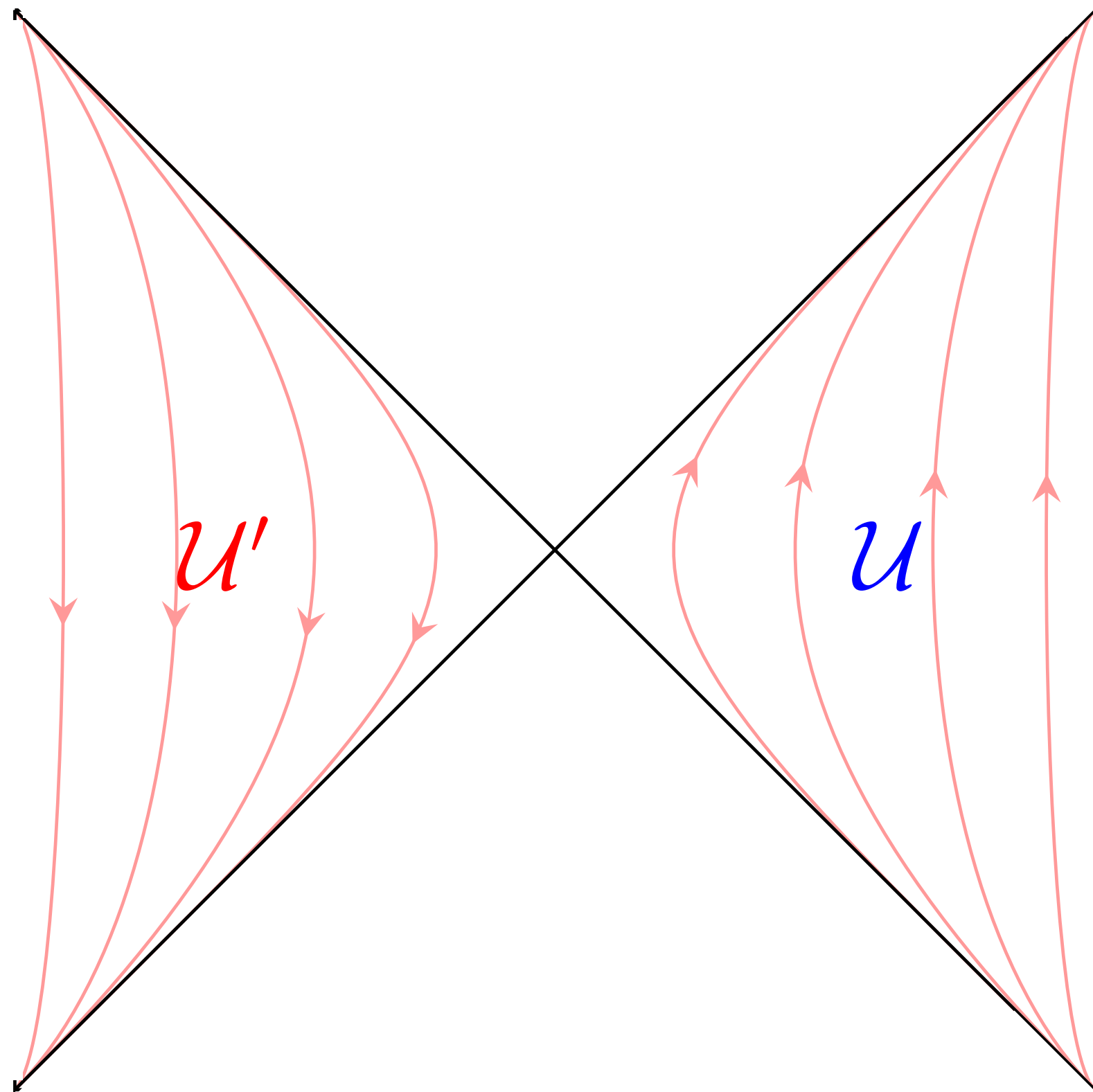
can have $S_{\mathcal{U}|C_3} \neq S_{\mathcal{U}|C_4}$

\Rightarrow many different gravitational entropies per region if multiple observers/QRFs

Type transition beyond gravity

[Fewster, Janssen, Loveridge, Rejzner, Waldron, CMP '25]

no multiple QRFs, no subsystem relativity



so far gravity, but similar construction applies also in QFT
more generally without imposing constraints on states
(within QRF-based measurement scheme in AQFT)

⇒ need Killing horizon

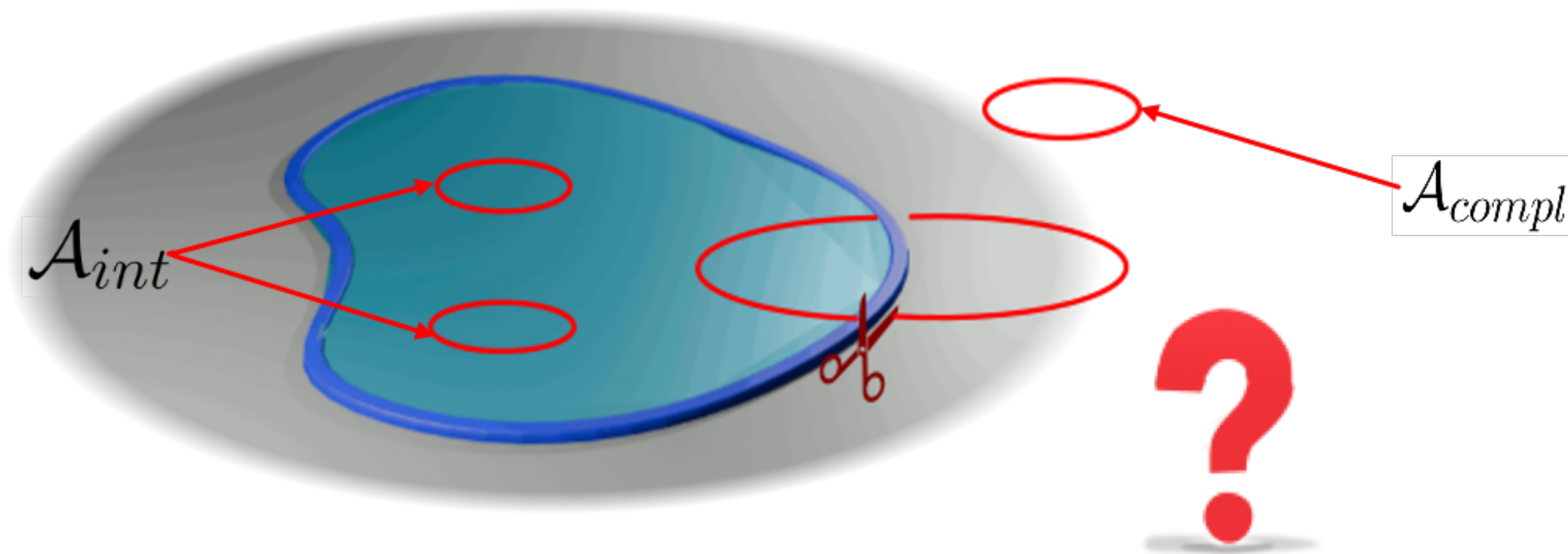
⇒ more generally: describing QFT relative to QRFs can lead to
better behaved entropies

e.g. scalar QFT across Rindler horizon in Minkowski

Relational entanglement entropies in gauge theory

[Araújo-Regado, PH, Sartini, 2506.23459]

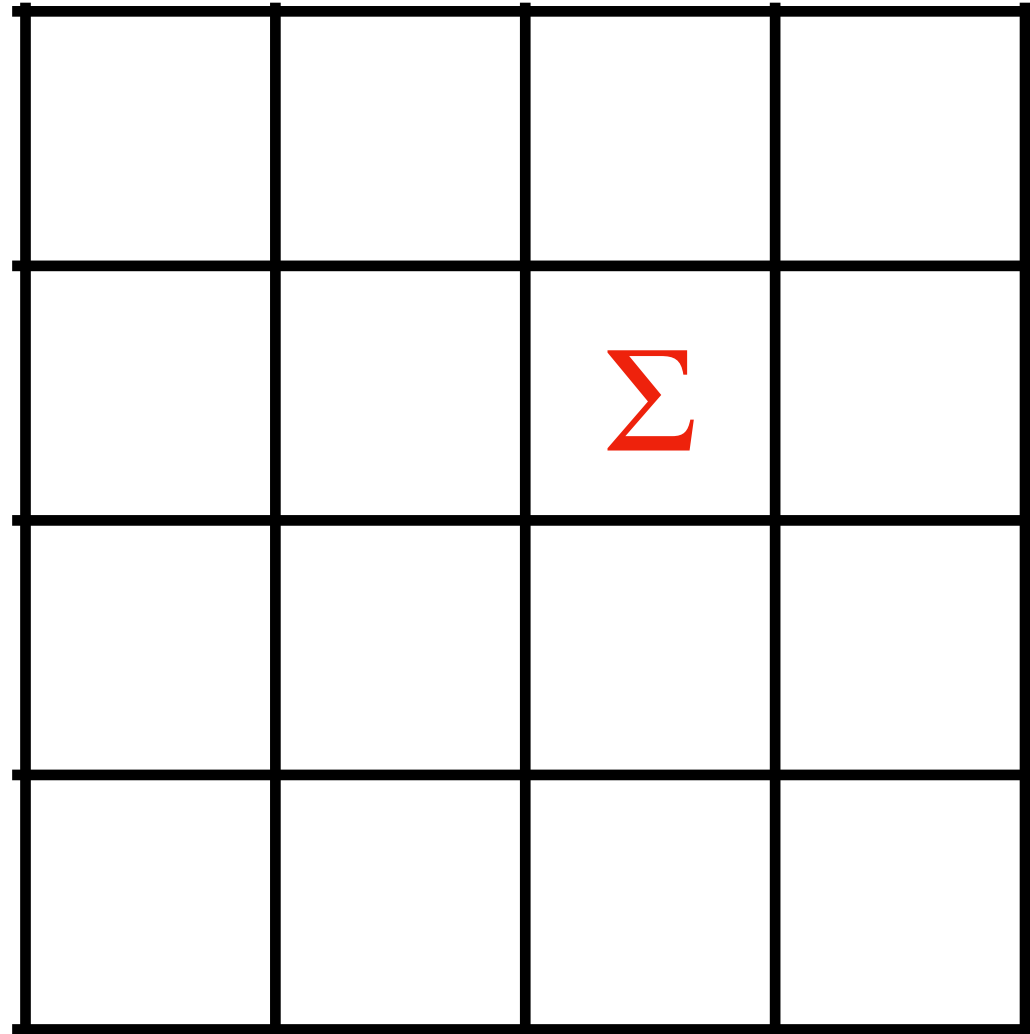
Standard regional center algebras



Standard regional center algebras

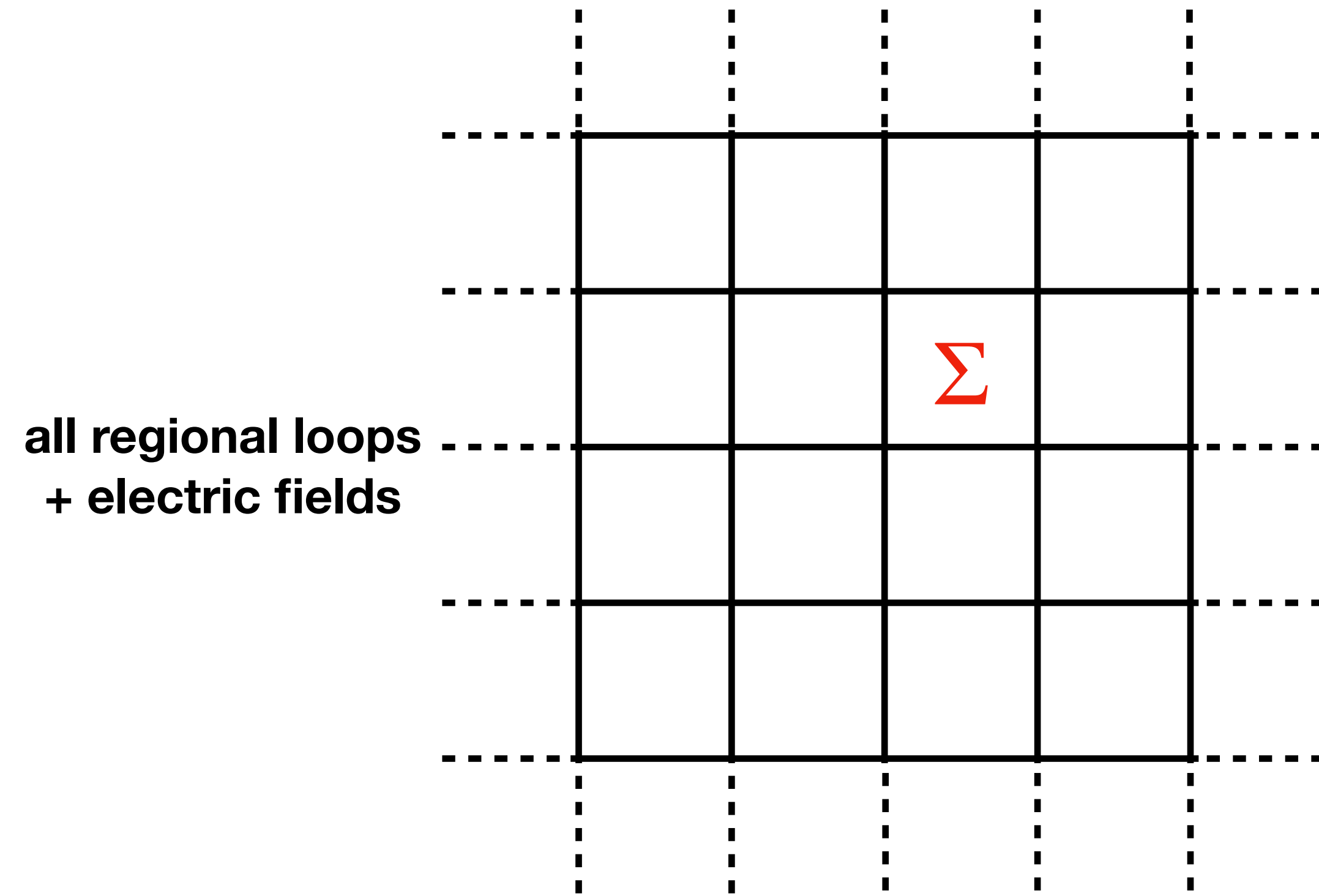
$$\mathcal{A}_{\text{phys}} \supset \bigoplus_q \mathcal{A}_{\text{int}}^q \otimes \mathcal{A}_{\text{compl}}^q$$

all regional loops
+ electric fields



Standard regional center algebras

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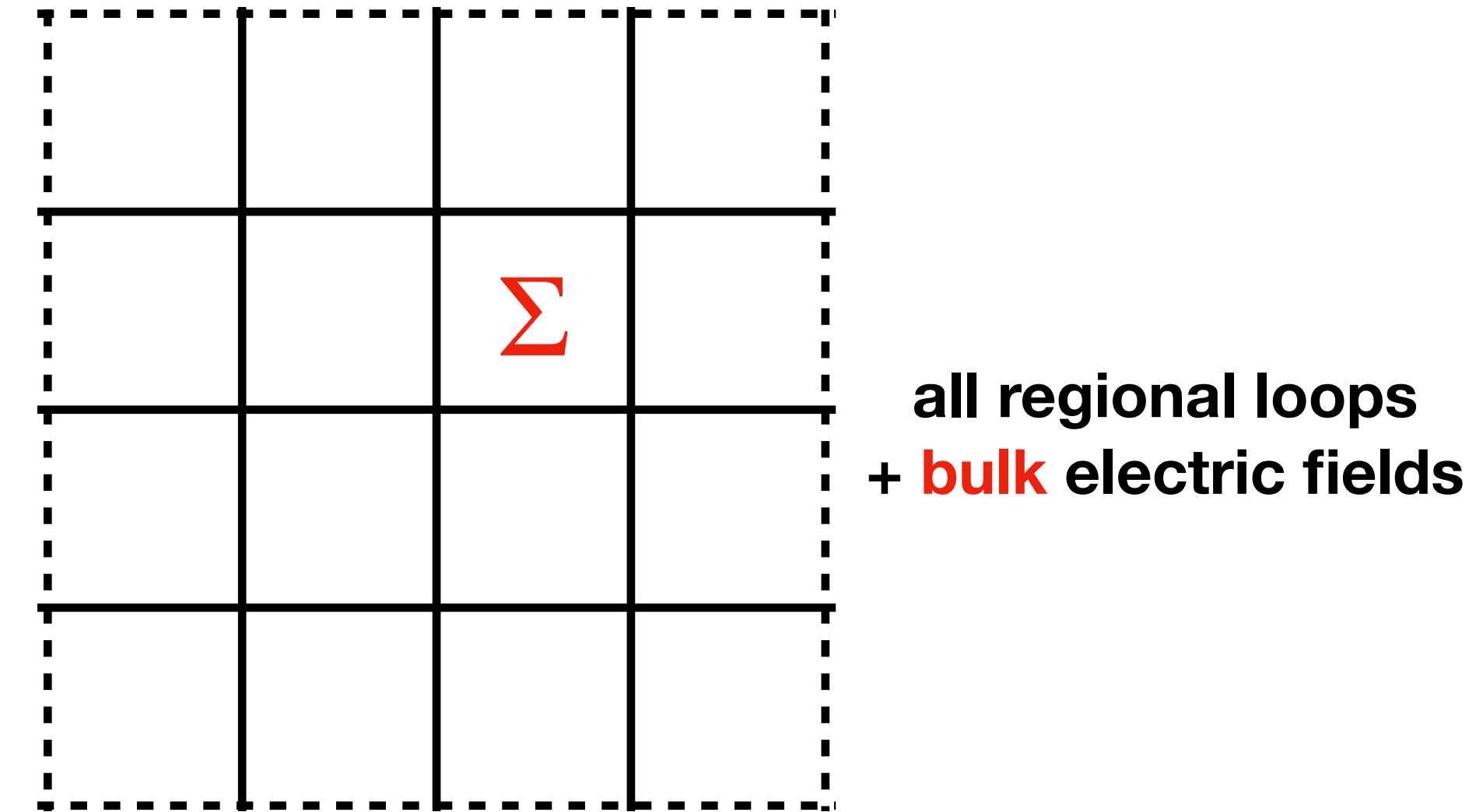
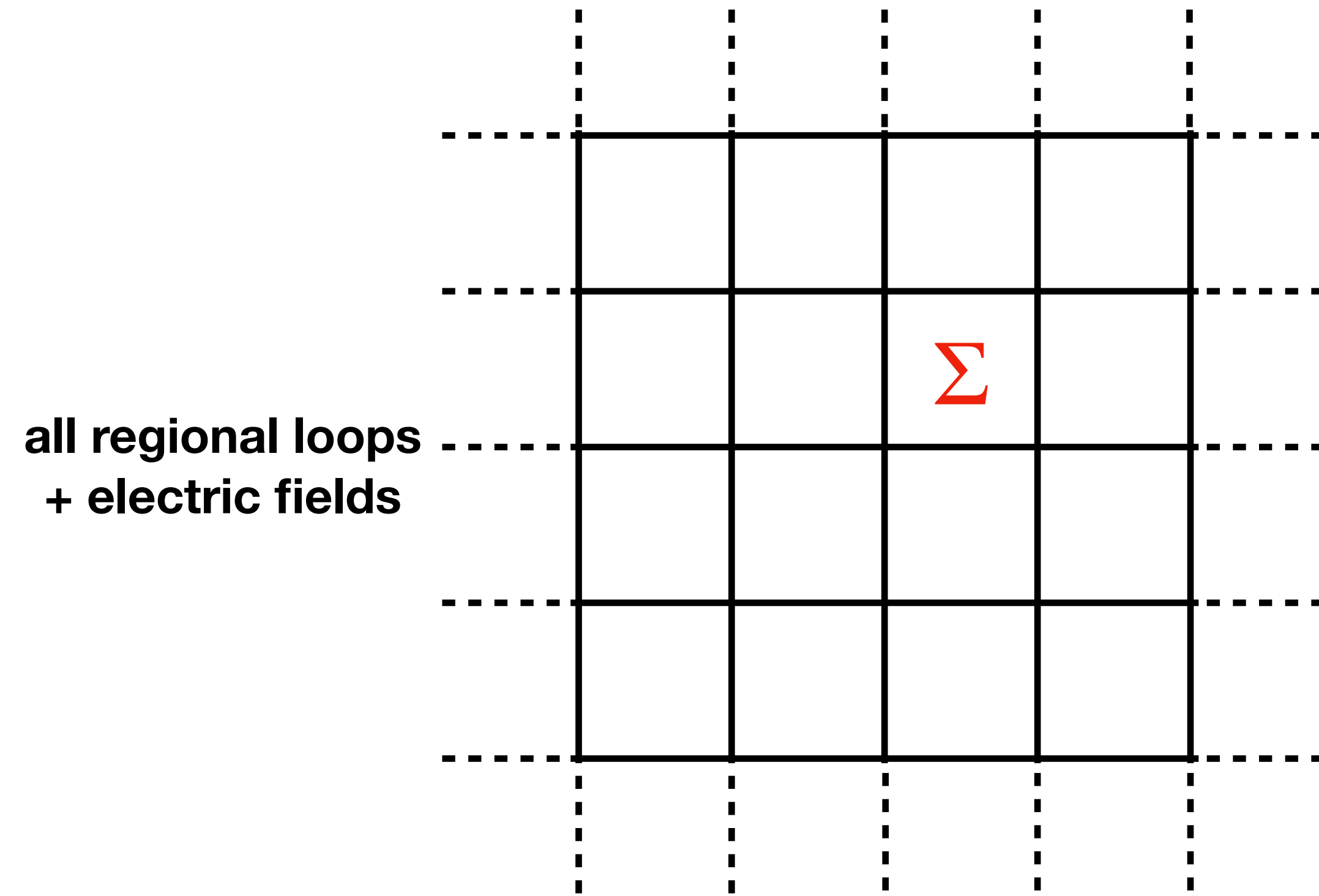
⇒ by Gauss: get (Casimirs of) normal electric flux as center

[Casini, Huerta, Rosbal '13; Soni, Trivedi '15; Delcamp, Dittrich, Riello '16]

$$\mathcal{A}_E$$

Standard regional center algebras

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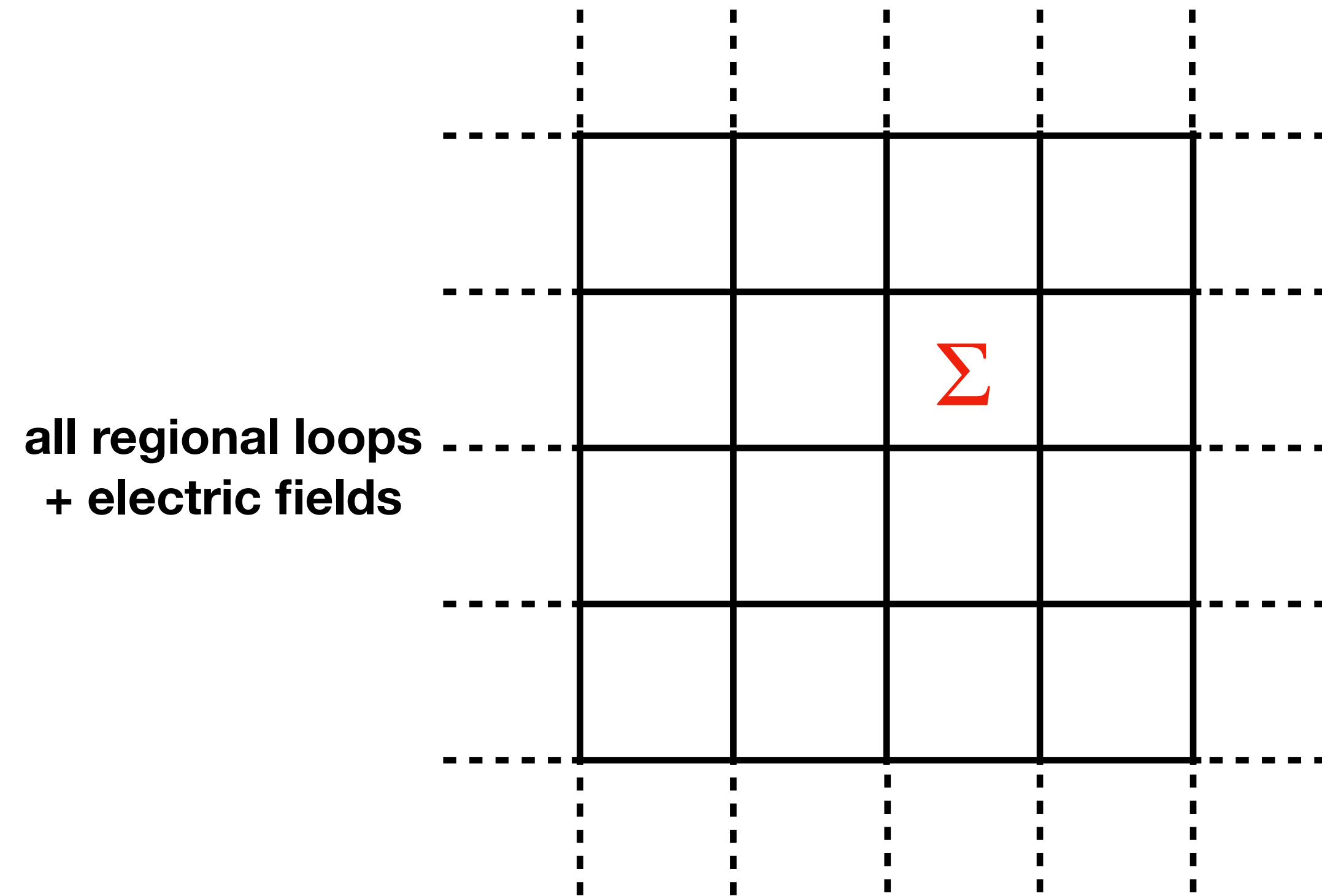
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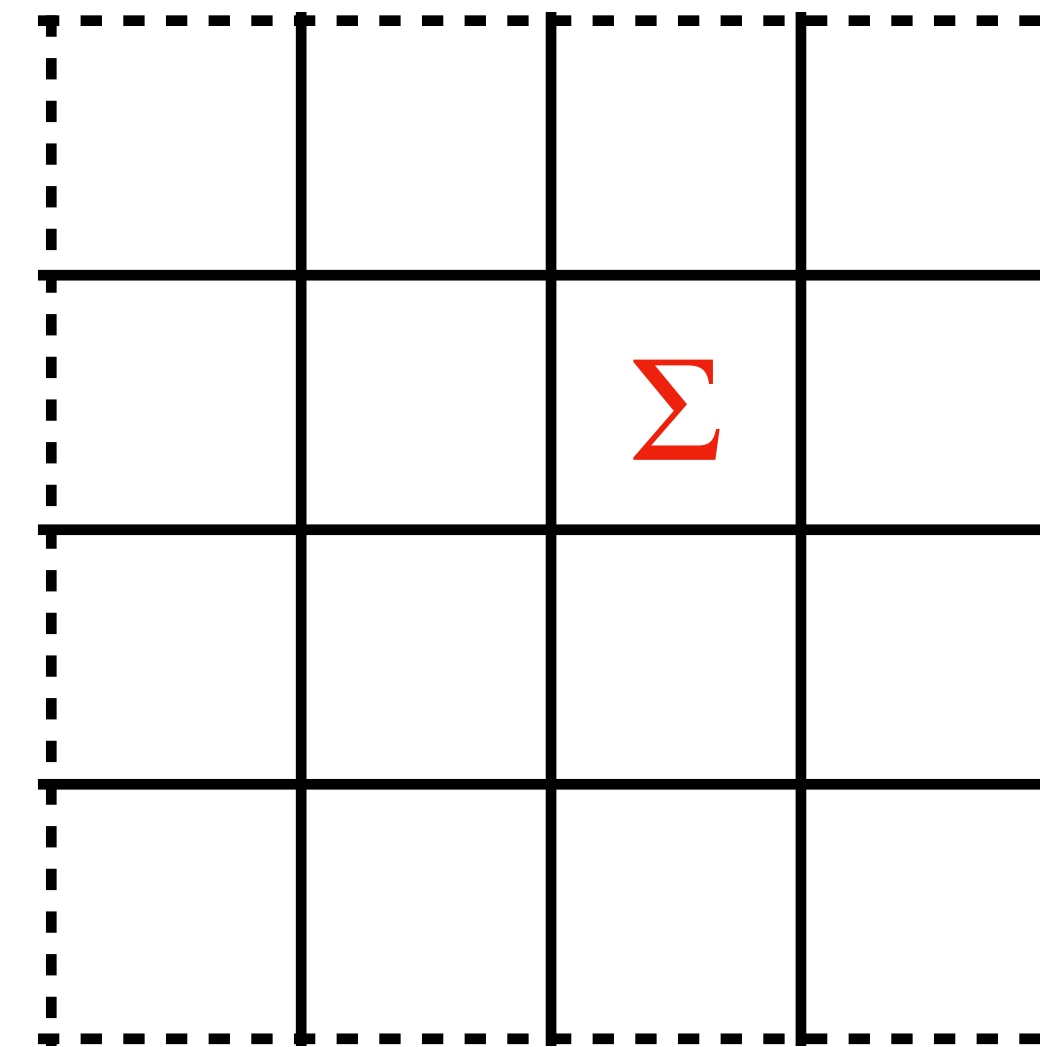
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$$\mathcal{A}_E \supsetneq$$



⇒ center: **corner Wilson loops**

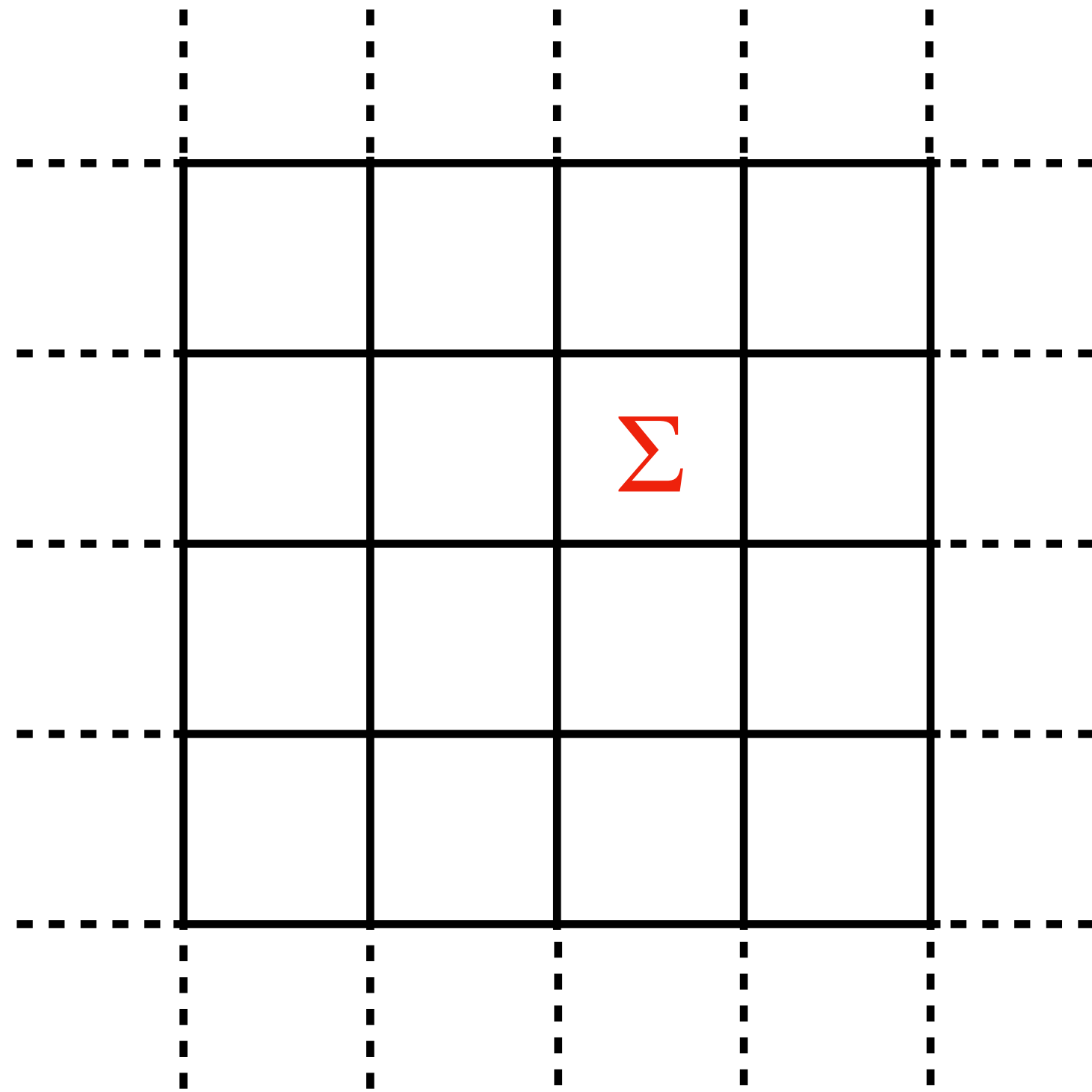
(+ certain loop dressed electric flux for nA theories)

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Standard regional center algebras

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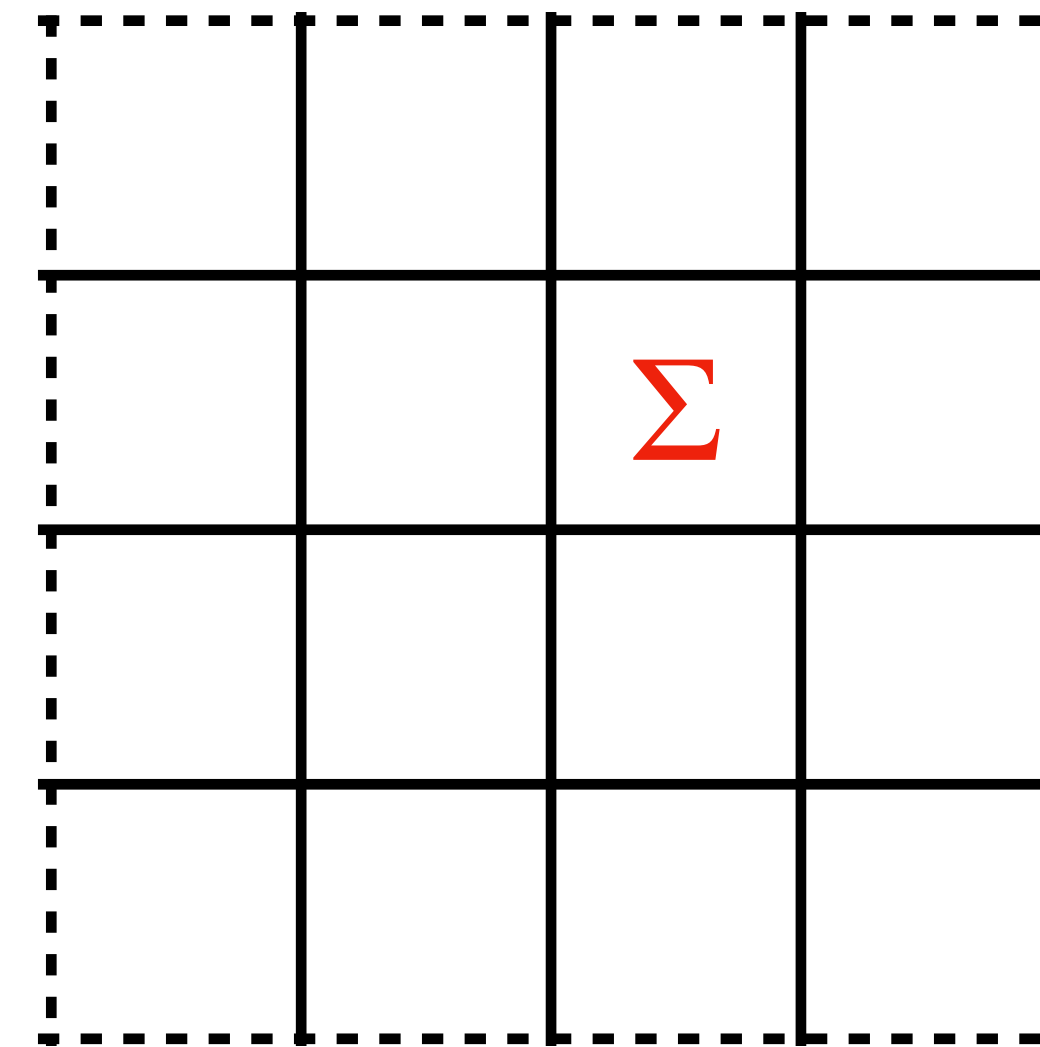
all regional loops
+ electric fields

⇒ by Gauss: get (Casimirs of) normal electric flux as center

[Casini, Huerta, Rosbal '13; Soni, Trivedi '15; Delcamp, Dittrich, Riello '16]

$$\mathcal{A}_E \not\supseteq$$

⇒ vN entropy non-distillable (doesn't measure entanglement)



all regional loops
+ **bulk** electric fields

⇒ center: **corner Wilson loops**

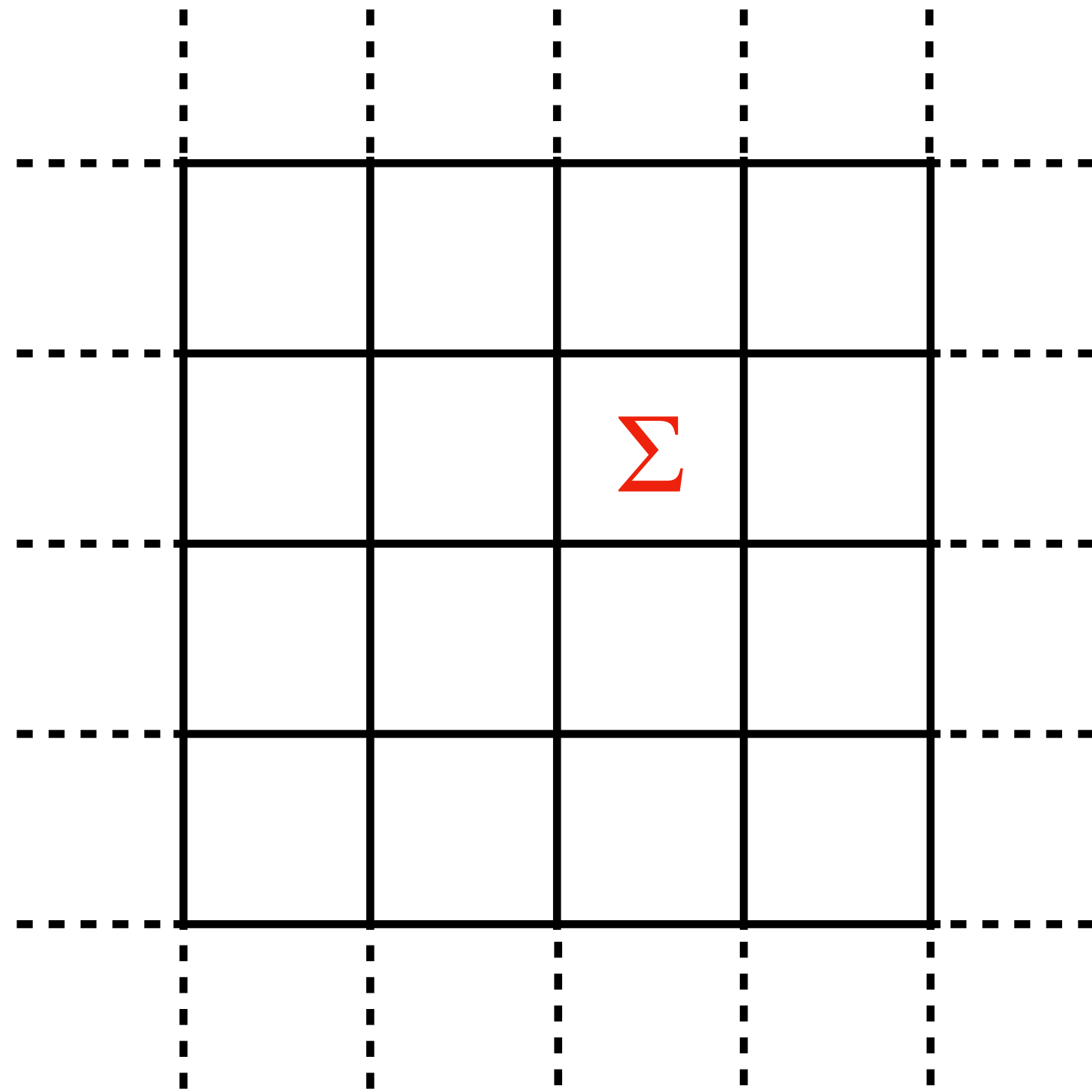
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$$\mathcal{A}_M$$

$$S_{\text{vN}}(\rho_{\text{int}}) = \sum_q p_q S_{\text{vN}}(\rho^q) + H(\{p_q\}) + \langle \log d_q \rangle$$

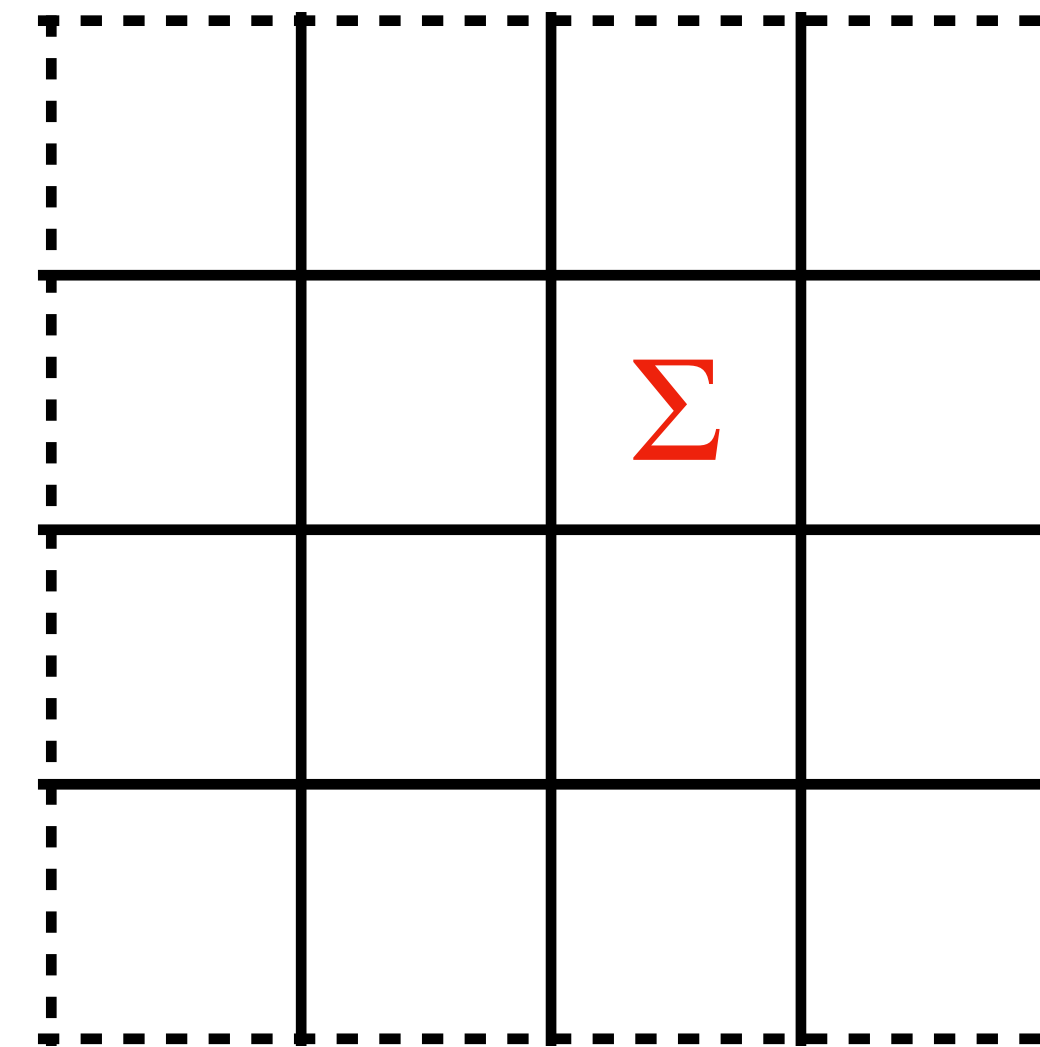
Including relational regional algebras



all regional loops
+ electric fields

⇒ by Gauss: get (Casimirs of) normal electric flux as center

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all regional loops
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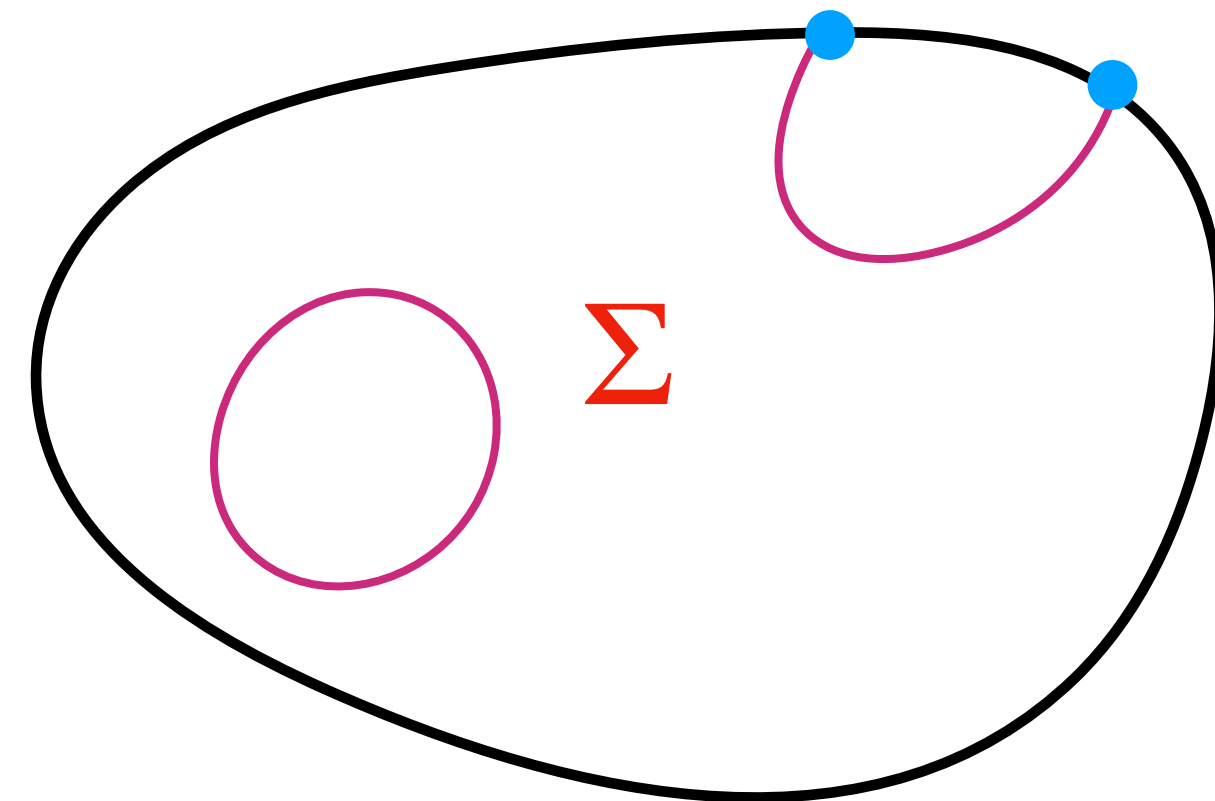
[Casini, Huerta, Rosbal '13; Delcamp, Dittrich, Riello '16]

$$\mathcal{A}_{\text{ext}} \supsetneq \mathcal{A}_E \supsetneq \mathcal{A}_{\text{int}} \supsetneq \mathcal{A}_M$$

Edge modes and corner symmetries

electric corner symmetries

$$Q[\rho] = \int_{\partial\Sigma} \text{Tr}(\rho \star F_\Phi)$$



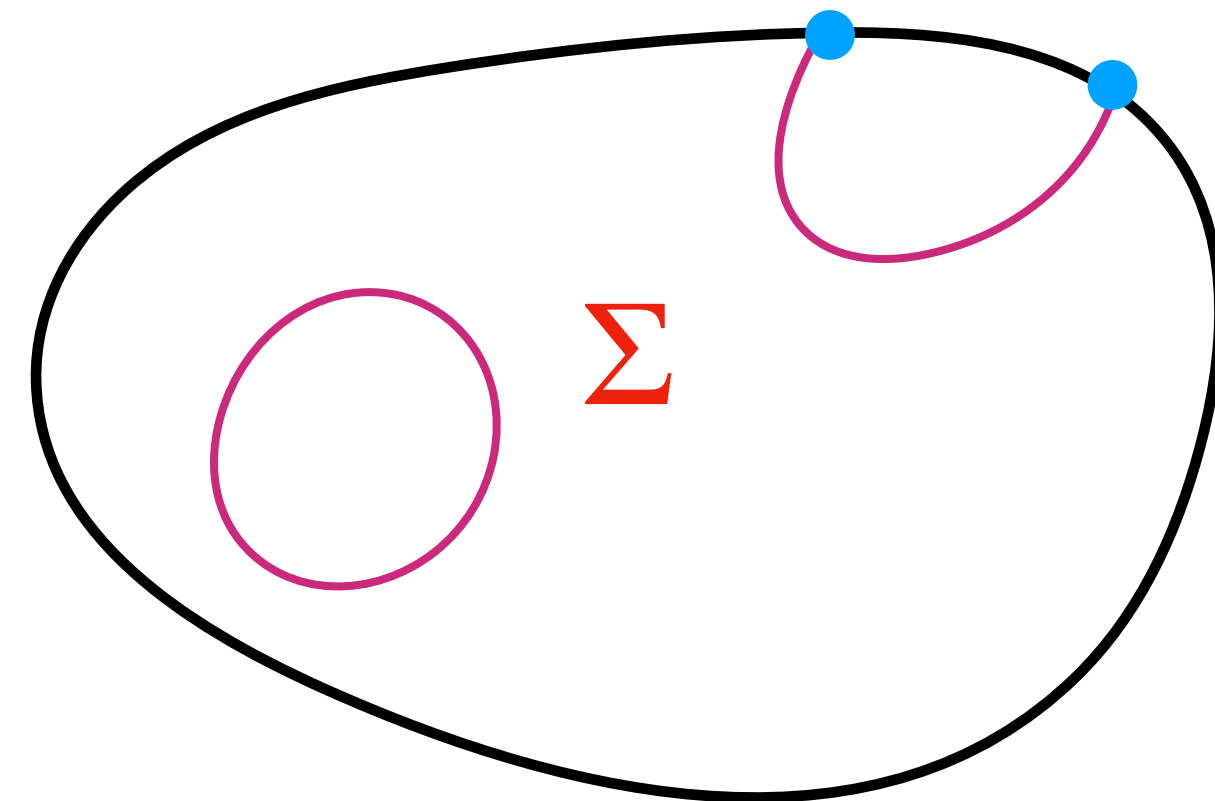
edge mode Φ

[Donnelly, Freidel '16; Geiller, Jai-akson '19; Ball, Law, Wong '24;...]

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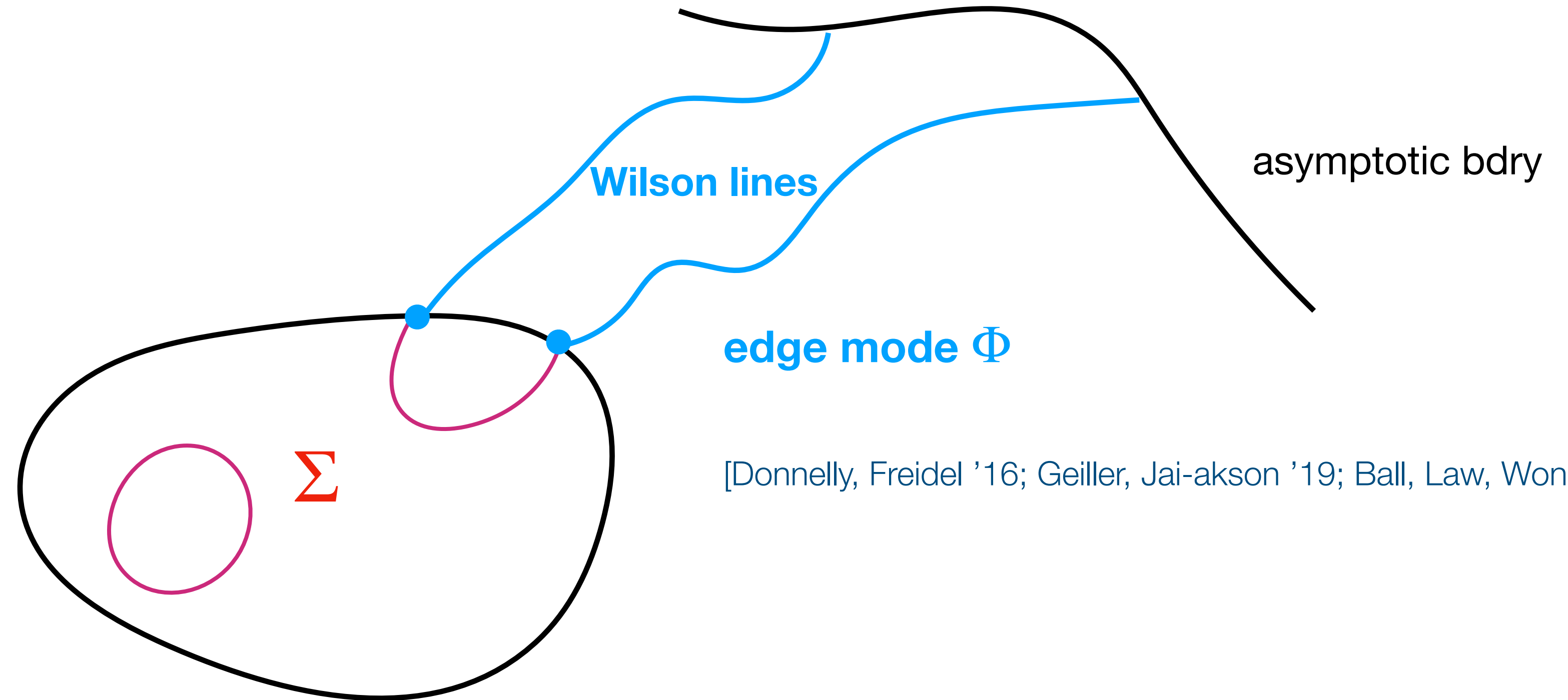
[Donnelly, Freidel '16; Geiller, Jai-akson '19; Ball, Law, Wong '24;...]

\Rightarrow how do you realize them?

Edge modes as QRFs

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[Donnelly, Freidel '16; Geiller, Jai-akson '19; Ball, Law, Wong '24;...]

⇒ how do you realize them? as QRFs

extrinsic edge mode QRFs:

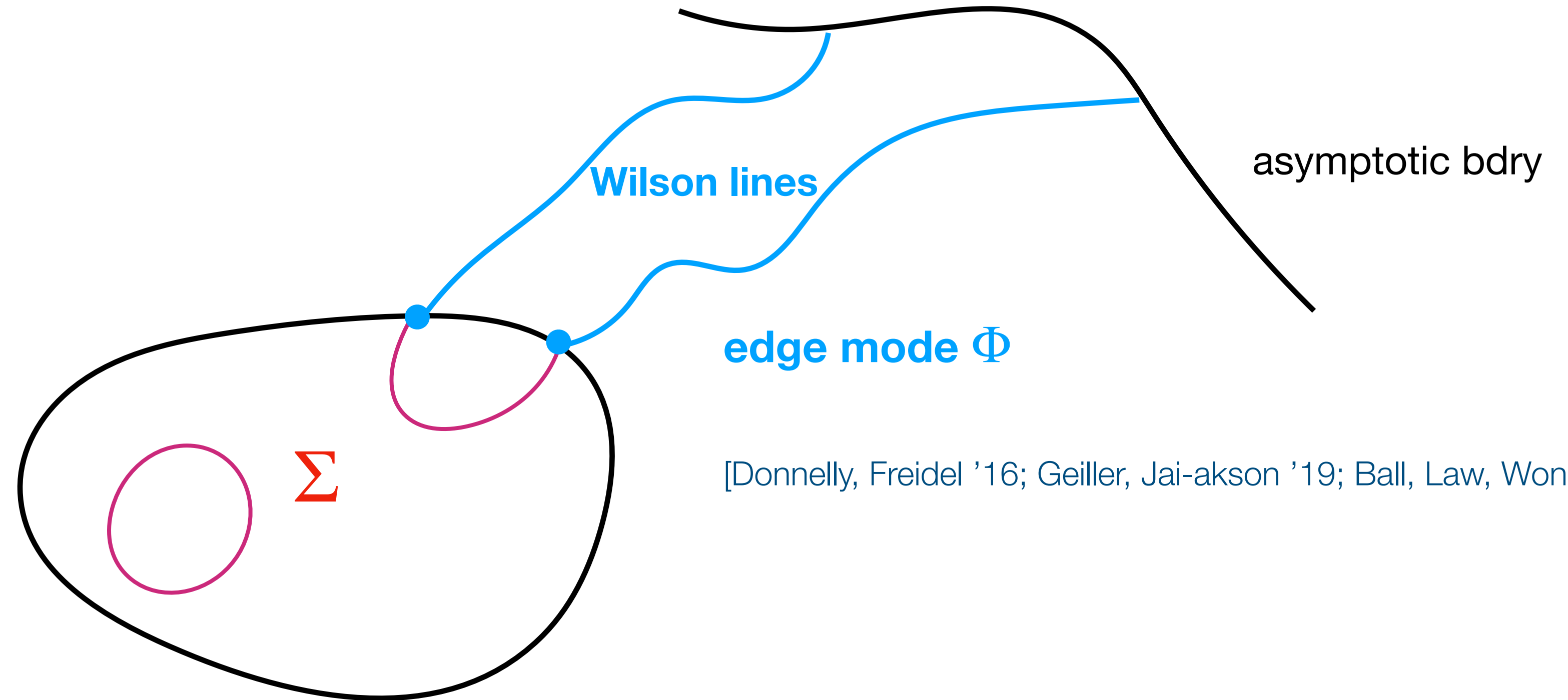
e.g. via extrinsic Wilson lines

[Carrozza, PH '21; Araújo-Regado, PH, Sartini, Tomova '24]

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e.g. via **Wilson lines anchored on the corner $\partial\Sigma$ or a Hodge decomp.**

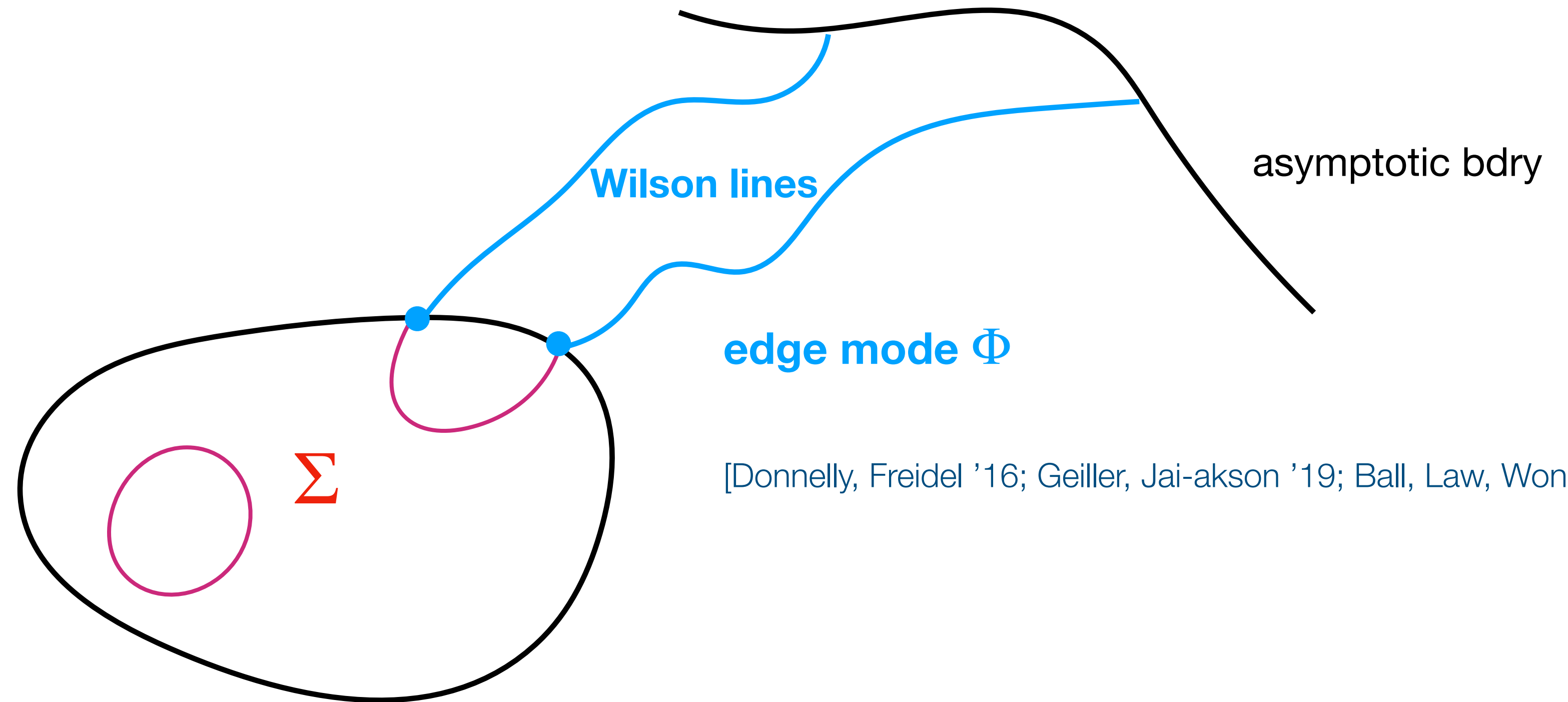
[Araújo-Regado, PH, Sartini, Tomova '24]

Edge modes as QRFs

electric corner symmetries

$$Q[\rho] = \int_{\partial\Sigma} \text{Tr}(\rho \star F_\Phi)$$

⇒ are reorientations of extrinsic QRFs!



[Donnelly, Freidel '16; Geiller, Jai-akson '19; Ball, Law, Wong '24;...]

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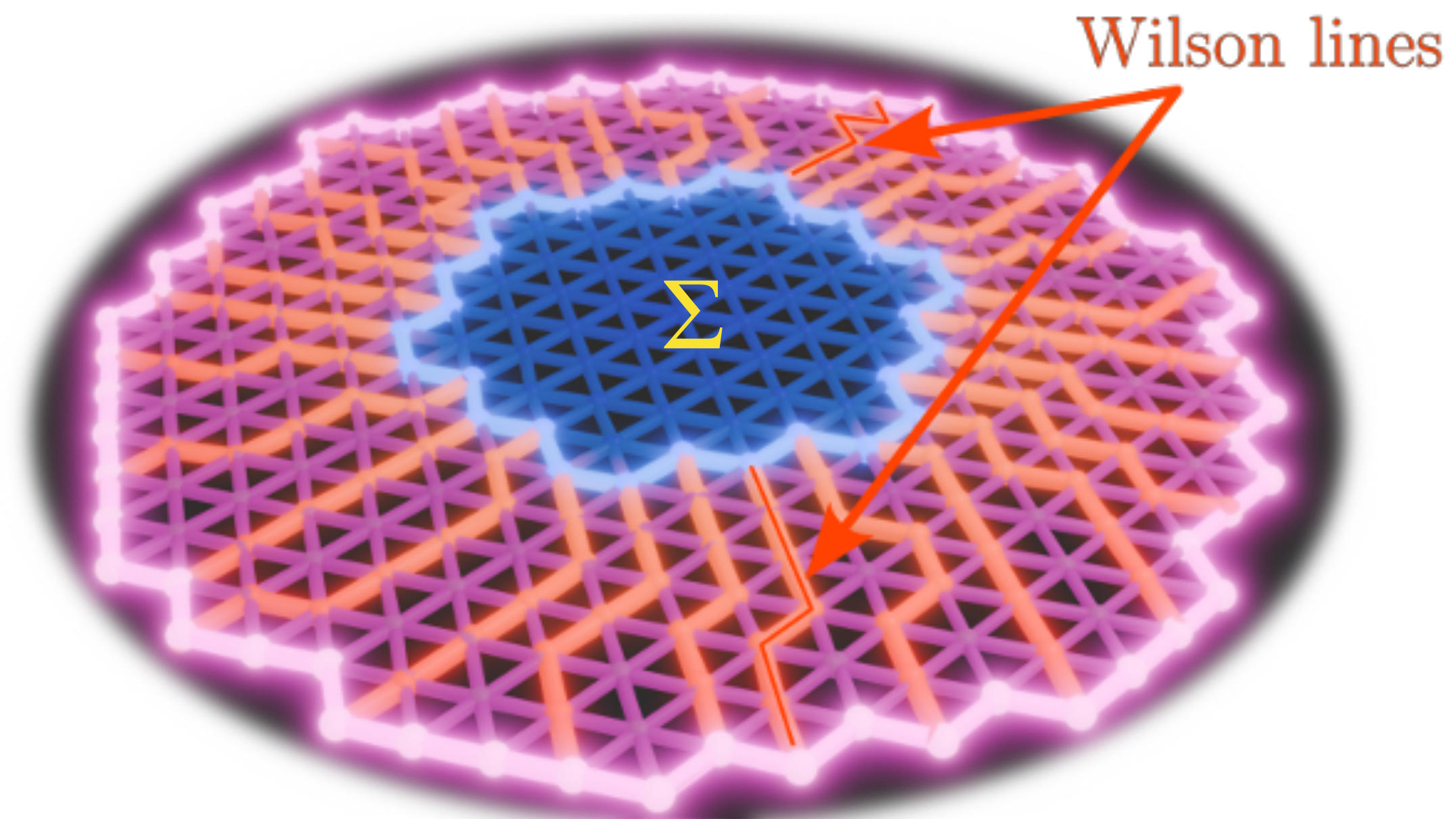
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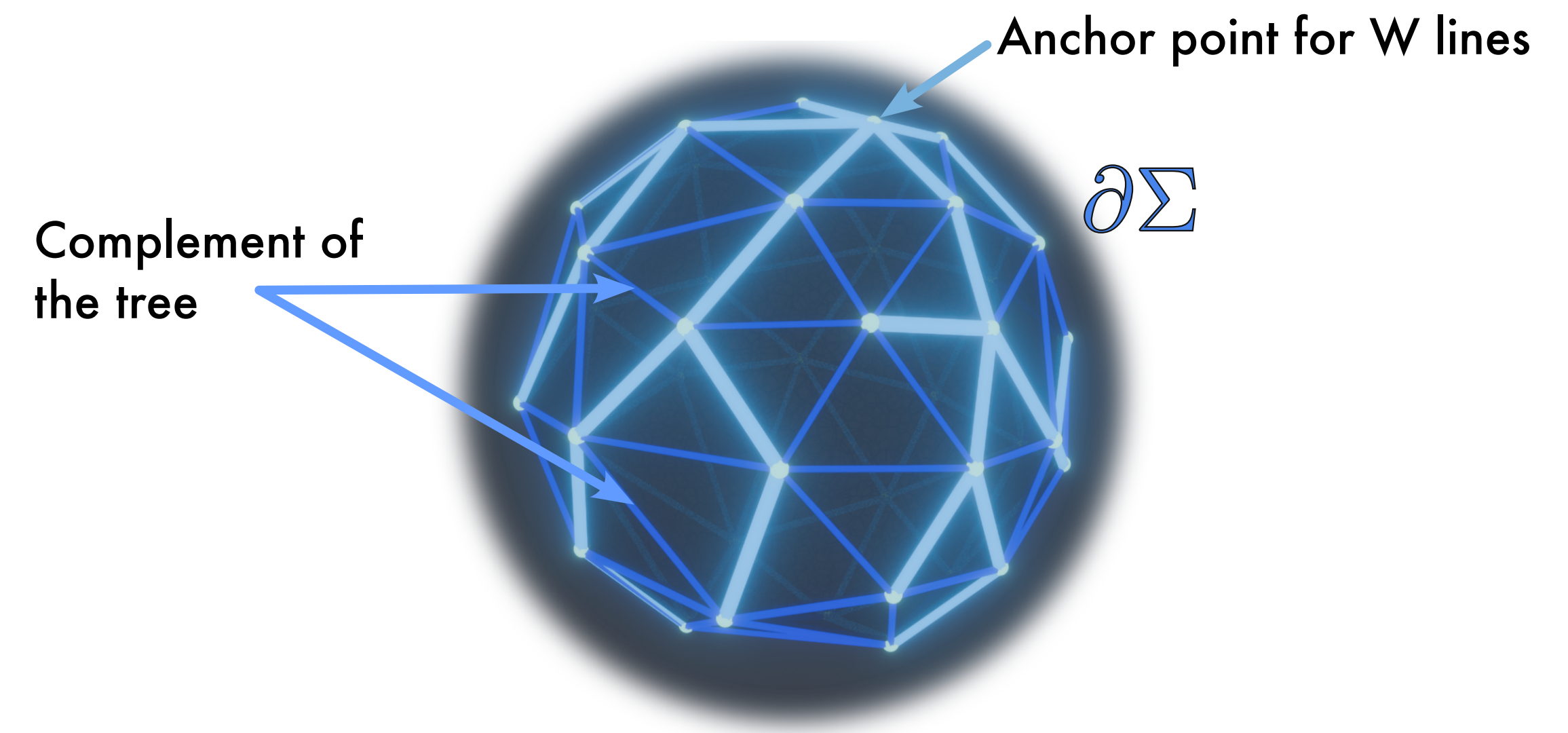
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Extrinsic and intrinsic QRFs on the lattice

Extrinsic: congruence of Wilson lines in the complement

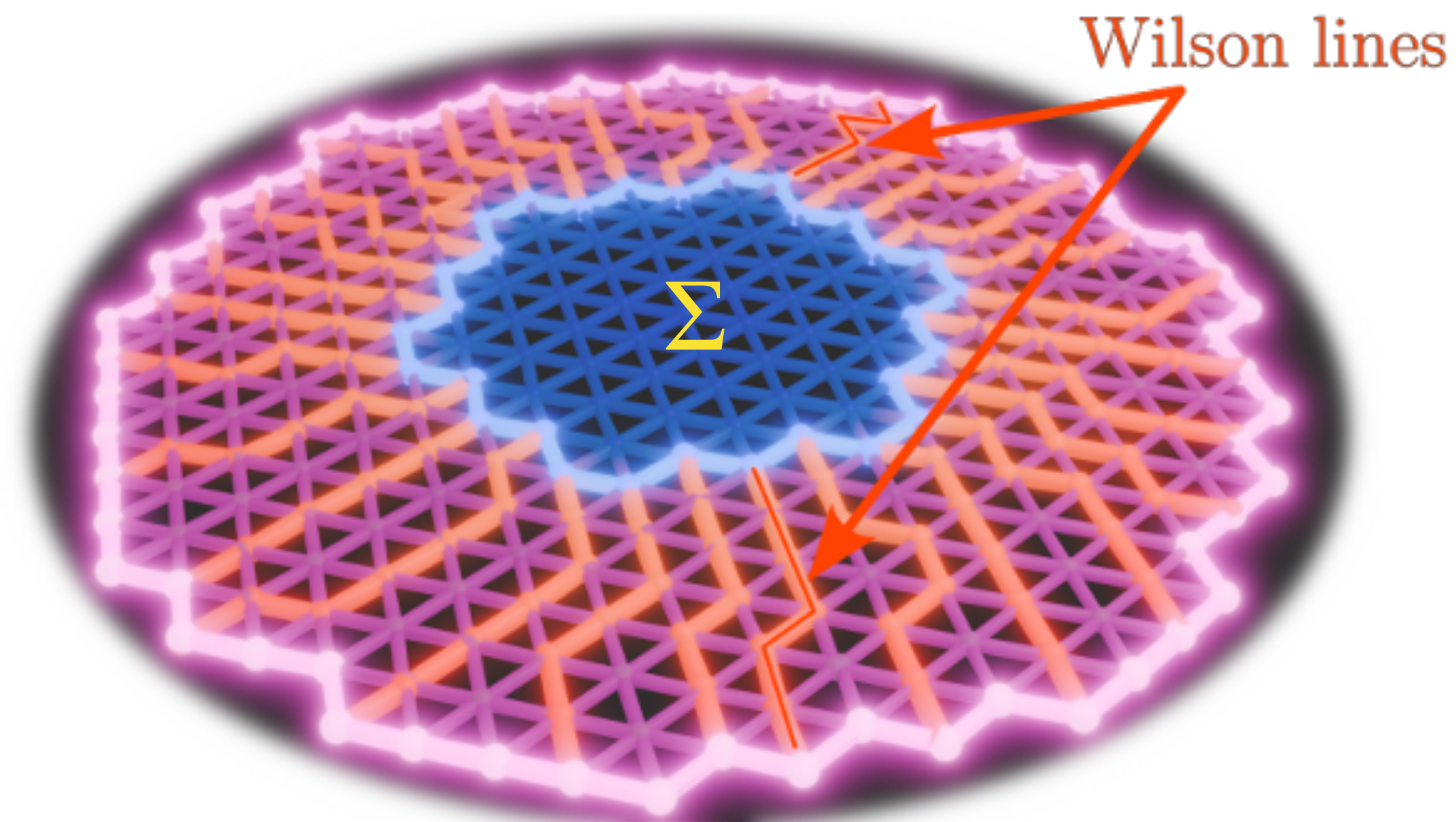


Spanning tree anchored on node N of the corner



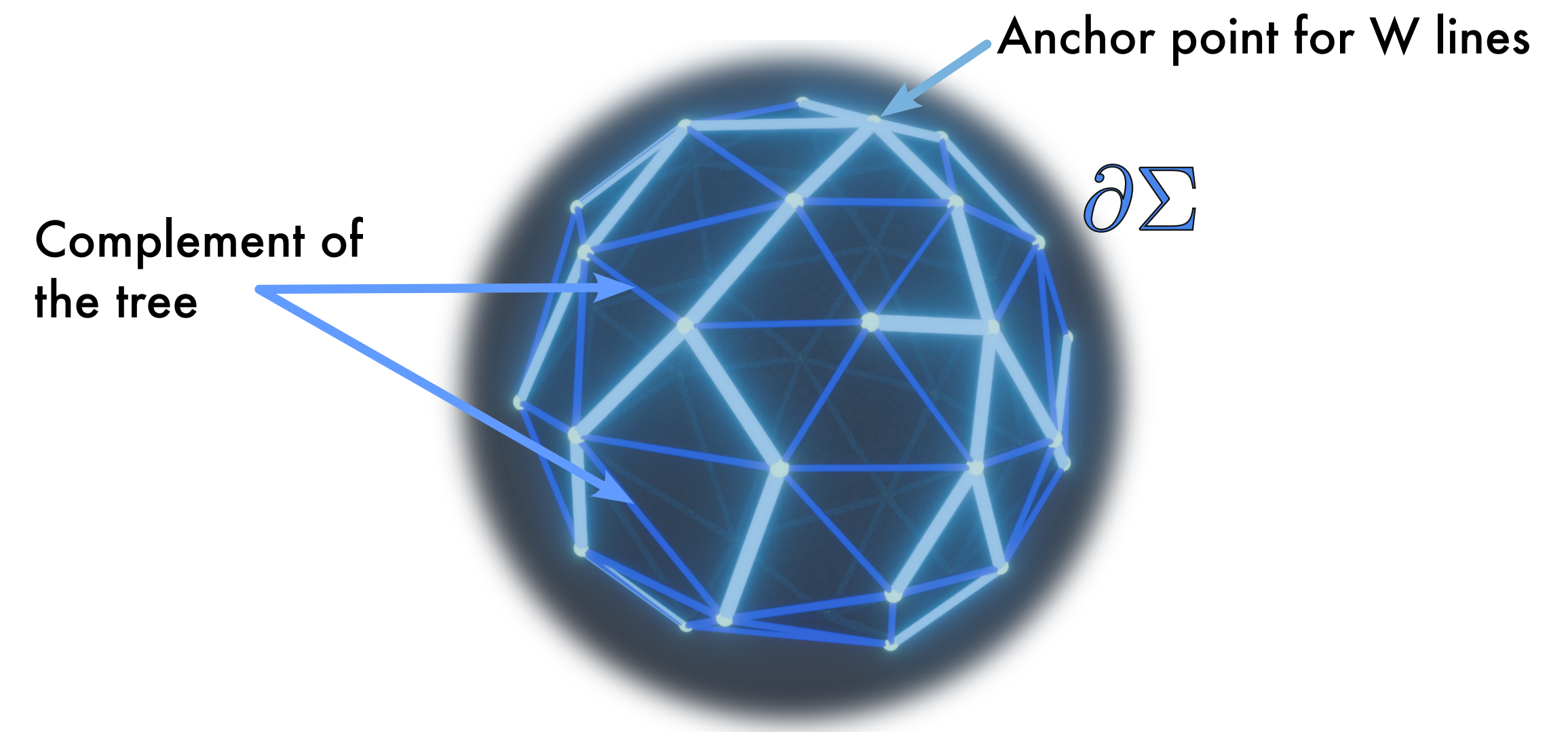
Extrinsic and intrinsic QRFs on the lattice

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complete ideal QRF: $\mathcal{H}_\Phi = L^2(G^{V_{\partial\Sigma}})$

Spanning tree anchored on node N of the corner

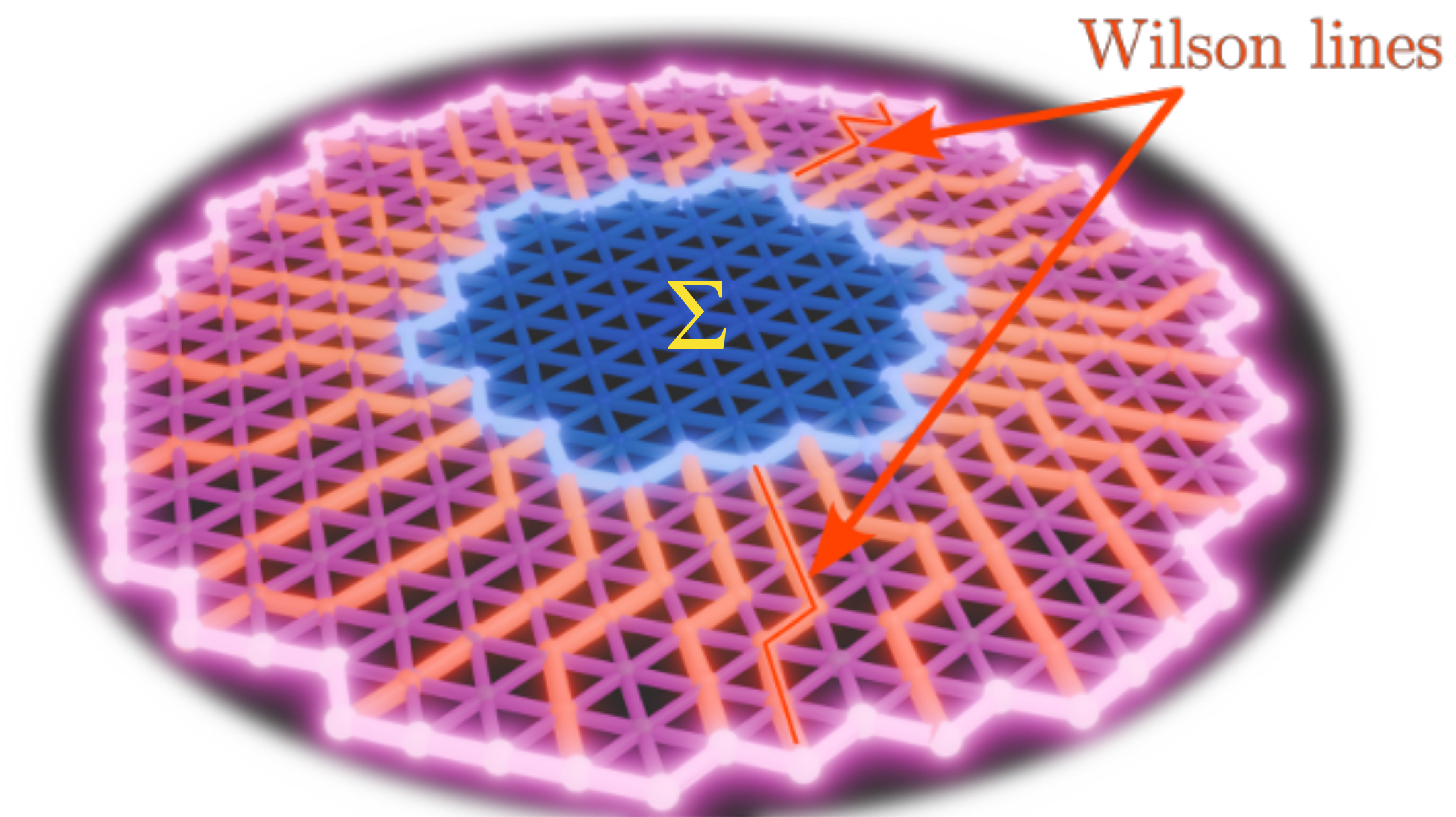


incomplete ideal QRF: $\mathcal{H}_{\tilde{\Phi}} = L^2(G^{V_{\partial\Sigma}^{-1}})$

⇒ can't deparametrize/dress at anchor point N

Extrinsic relational algebras

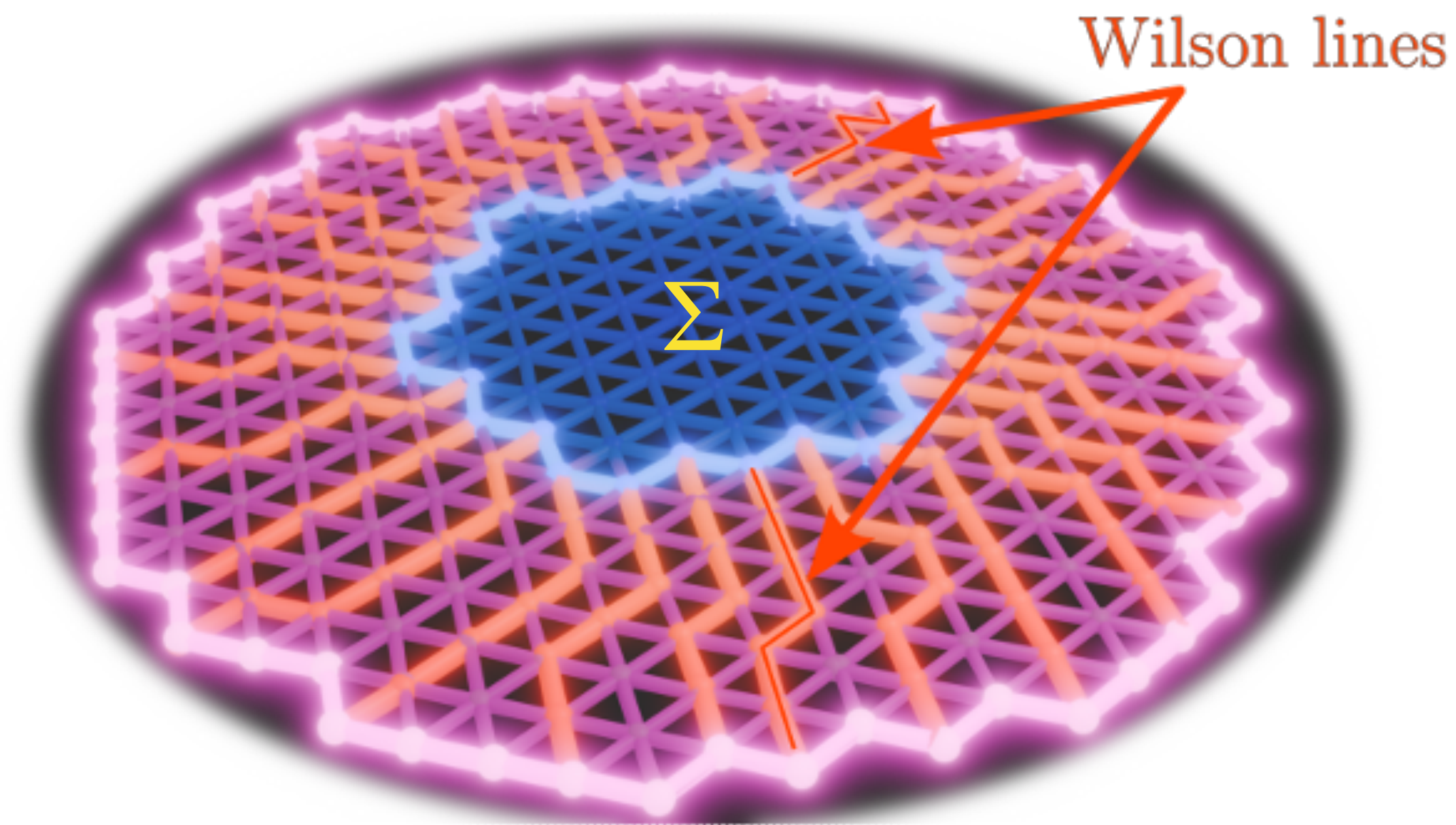
dress all the regional data in Σ (Wilson lines, electric fields) with extrinsic Wilson line QRF to asymp. bdry



\Rightarrow rel. obs. $O_{|\Phi}^g(a)$ with $a \in \mathcal{A}_\Sigma$

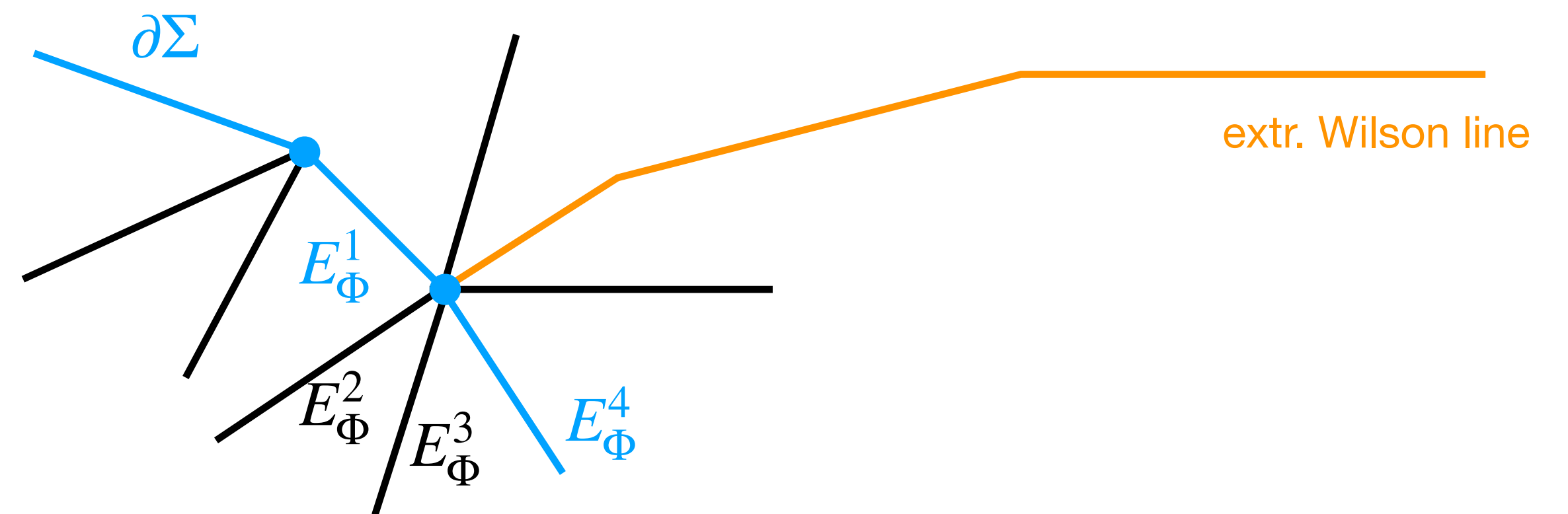
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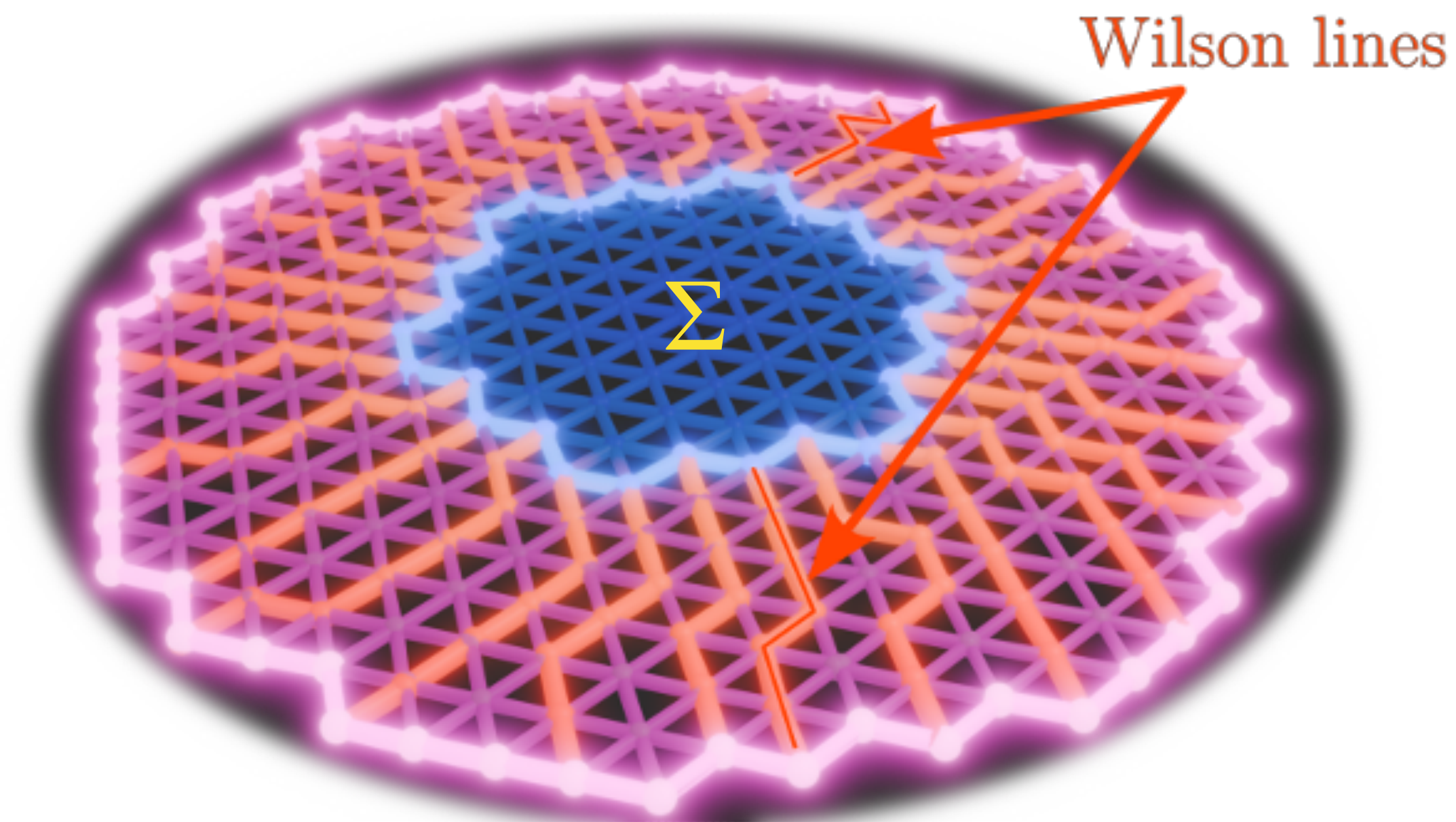
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But on $\mathcal{H}_{\text{phys}}$ Gauss's law gives us reorientations for free:



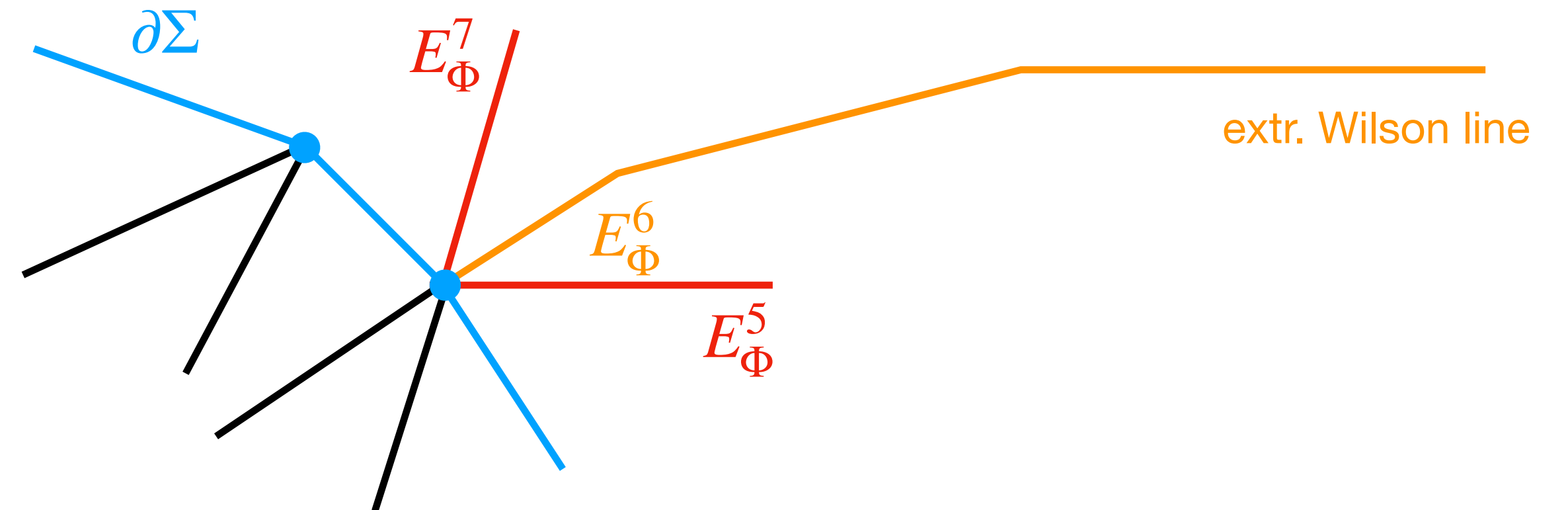
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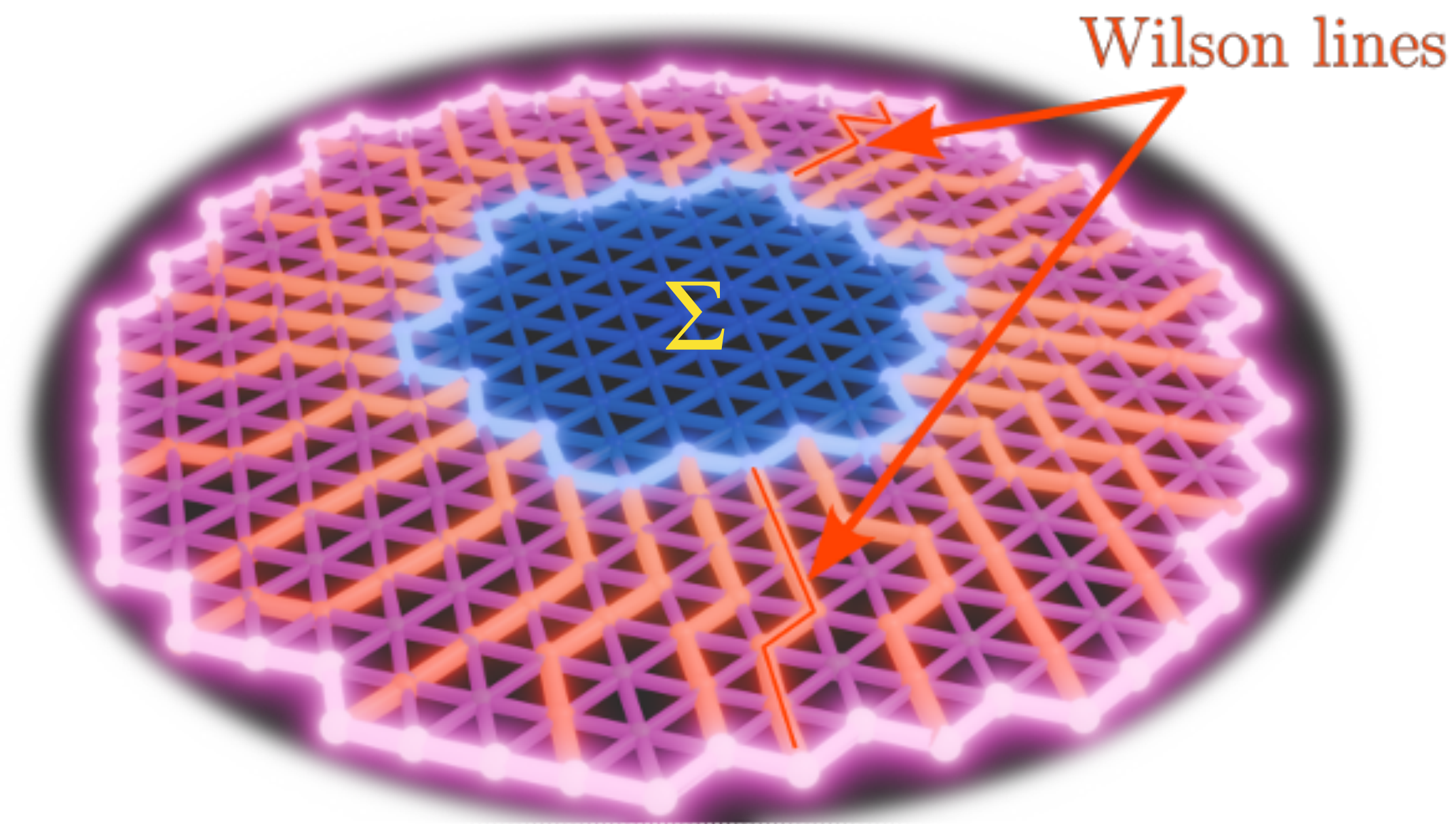
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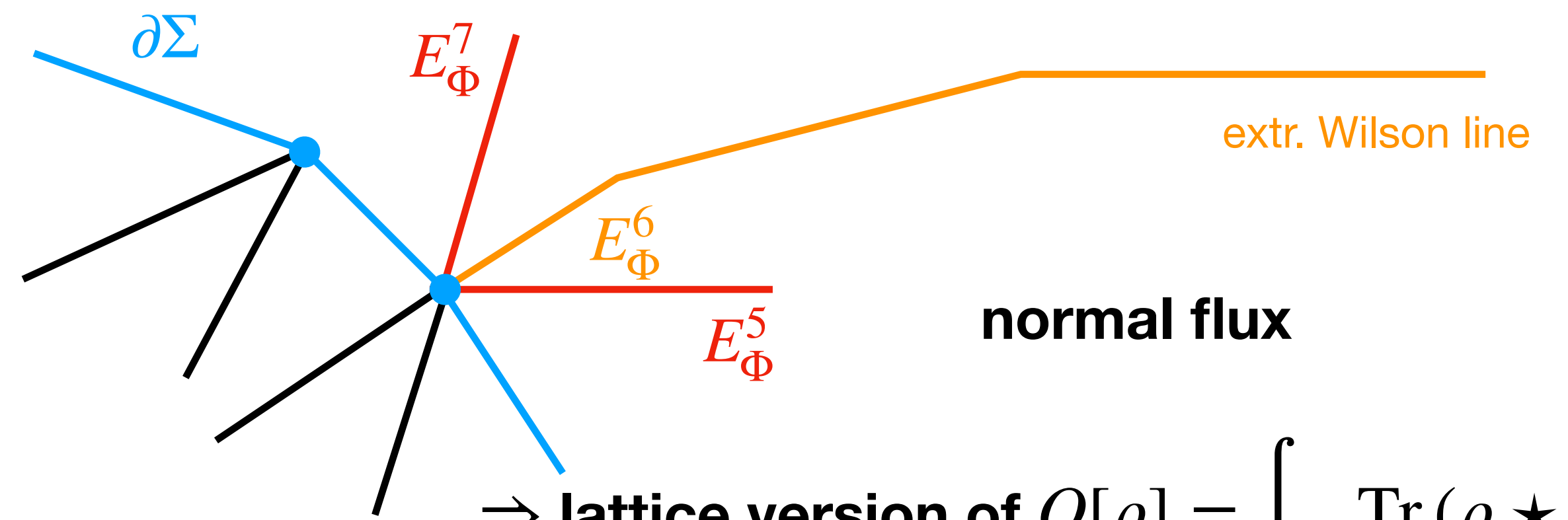
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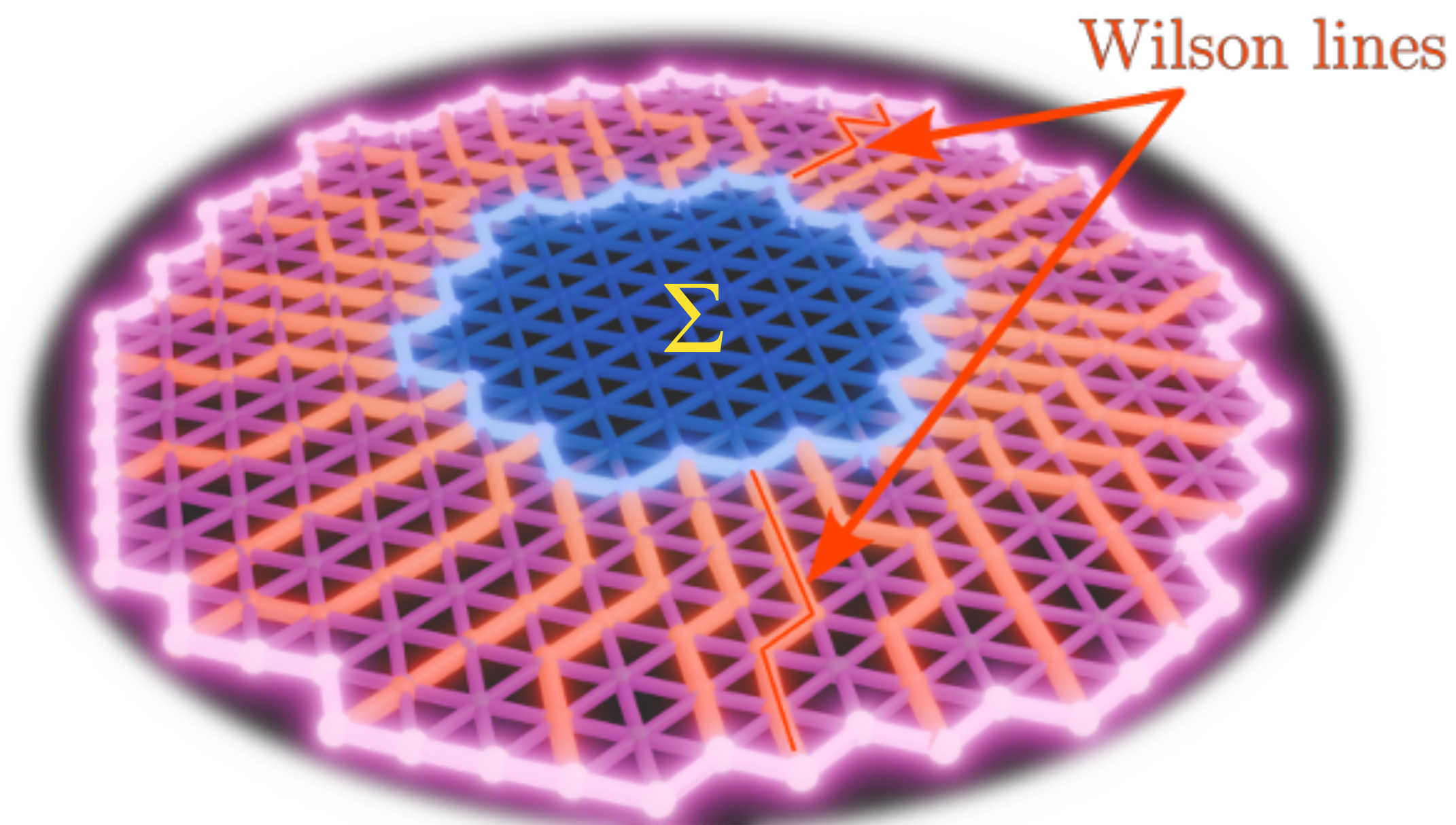
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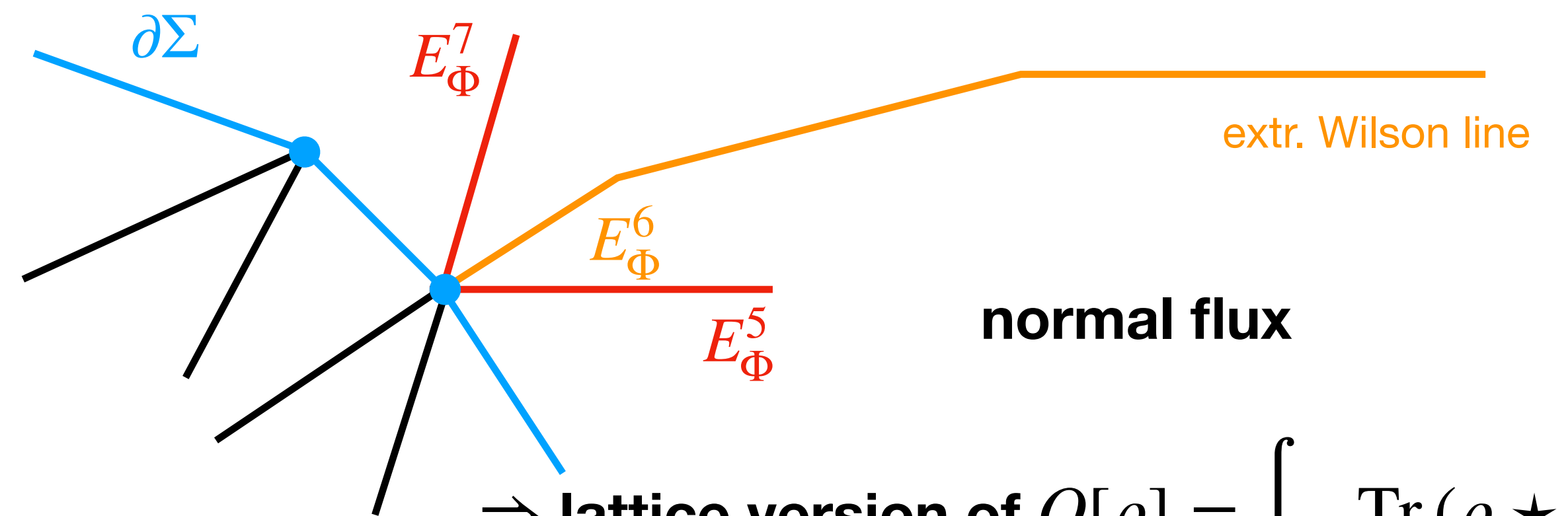
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But on $\mathcal{H}_{\text{phys}}$ Gauss's law gives us reorientations for free:



\Rightarrow lattice version of $Q[\rho] = \int_{\partial\Sigma} \text{Tr}(\rho \star F_\Phi)$

\Rightarrow electric corner symmetries are **reorientations** of Φ : $|g'\rangle_\Phi \mapsto V_\Phi(g) |g'\rangle_\Phi = |g'g^{-1}\rangle_\Phi$

Extrinsic relational algebras

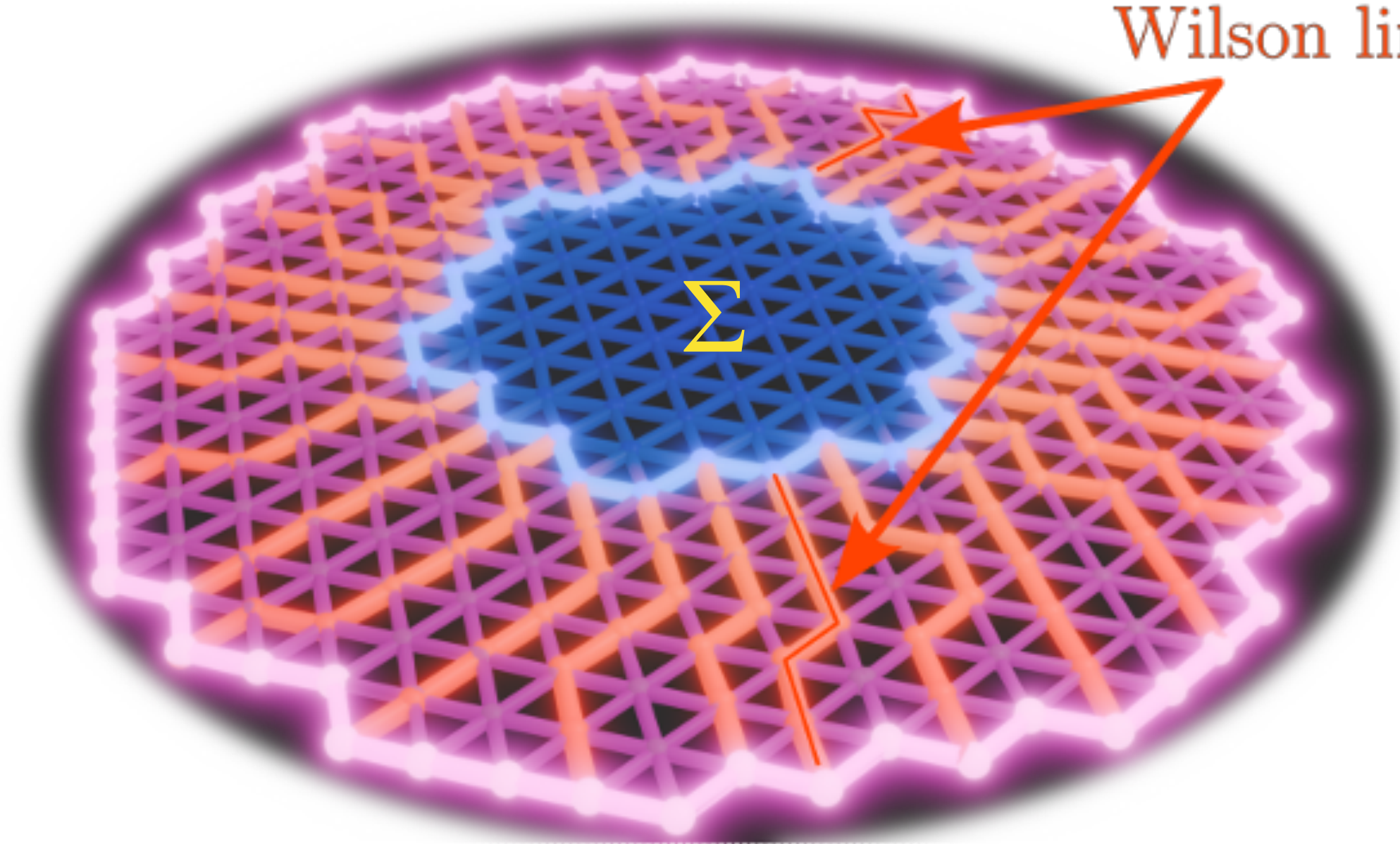
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Wilson lines

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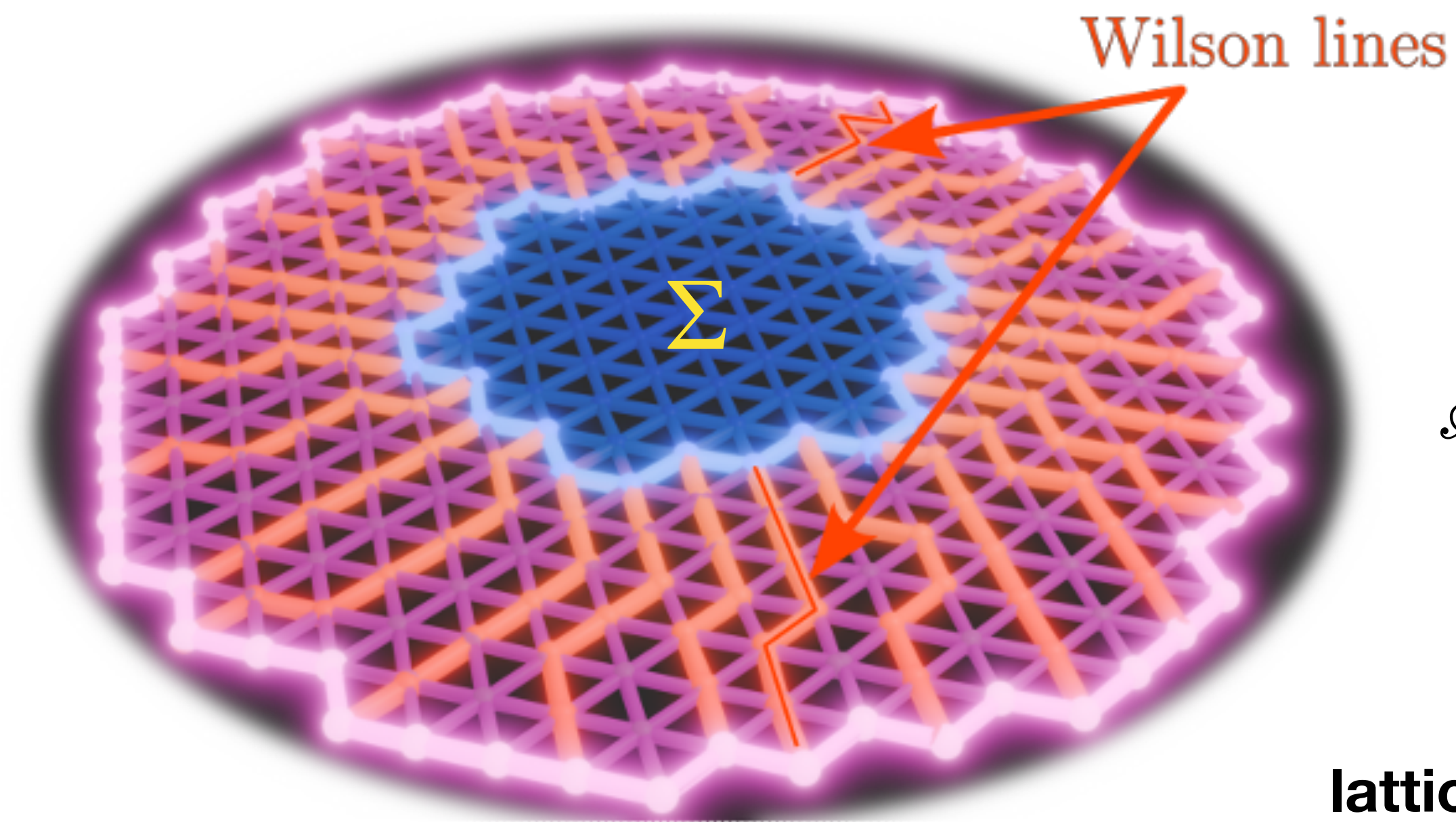
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$$\begin{aligned} \mathcal{A}_{\text{ext}} &= \mathcal{A}_\Sigma \rtimes_{\text{ext}} G^{V_{\partial\Sigma}} = \left(\mathcal{A}_\Sigma \otimes L^2(G^{V_{\partial\Sigma}}) \right)^{G^{V_{\partial\Sigma}}} \\ &= \left\langle O_{|\Phi}^e(a), V_\Phi(g) \mid a \in \mathcal{A}_\Sigma, g \in G^{V_{\partial\Sigma}} \right\rangle'' \end{aligned}$$



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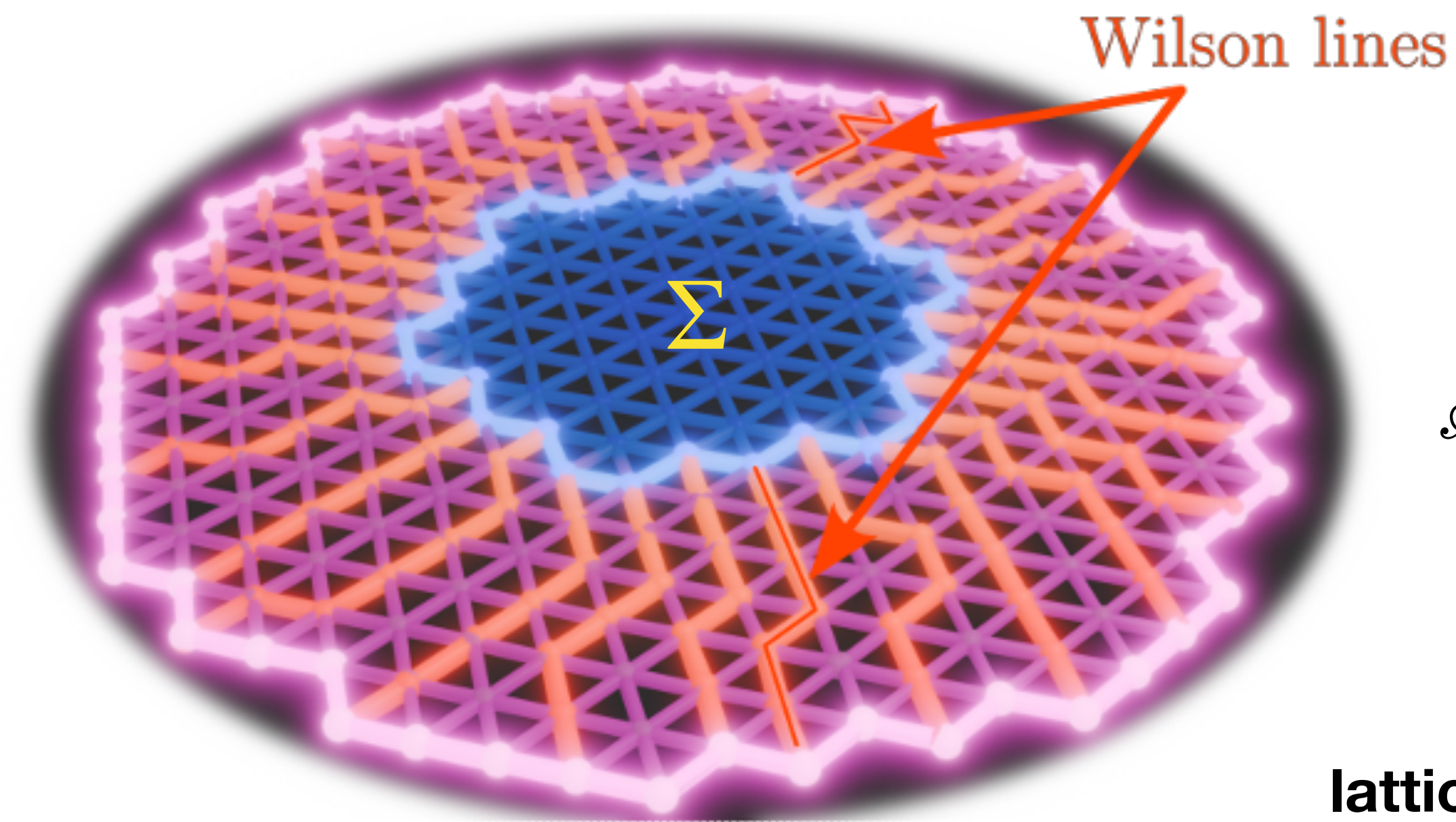
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lattice gauge theory analog of crossed product in gravity case

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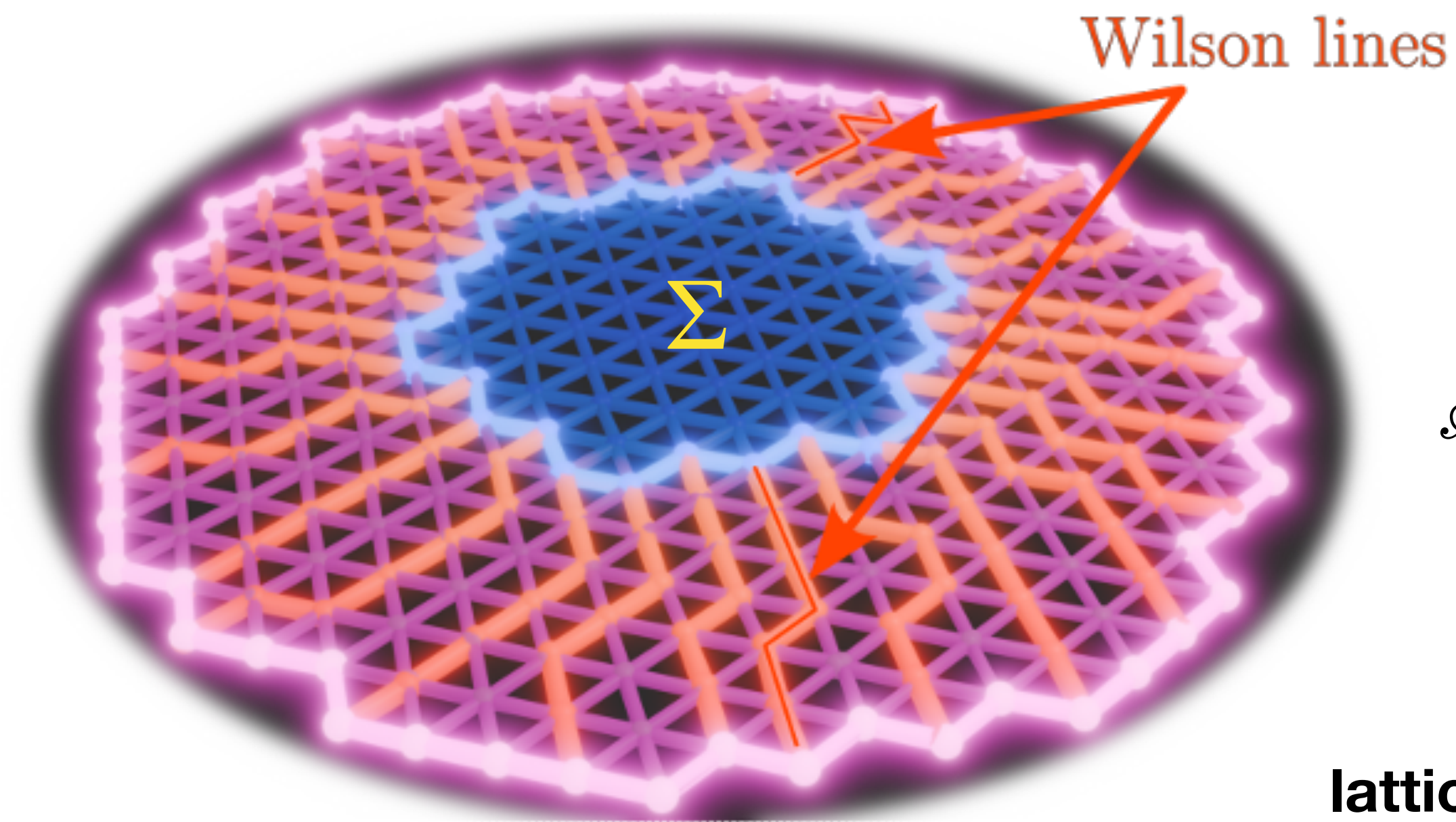
lattice gauge theory analog of crossed product in gravity case

\Rightarrow it is also a **factor**, can show using QRF techniques that indeed

$$\mathcal{H}_{\text{phys}} \simeq \mathcal{H}_{\text{ext}} \otimes \mathcal{H}_{\text{out}}$$

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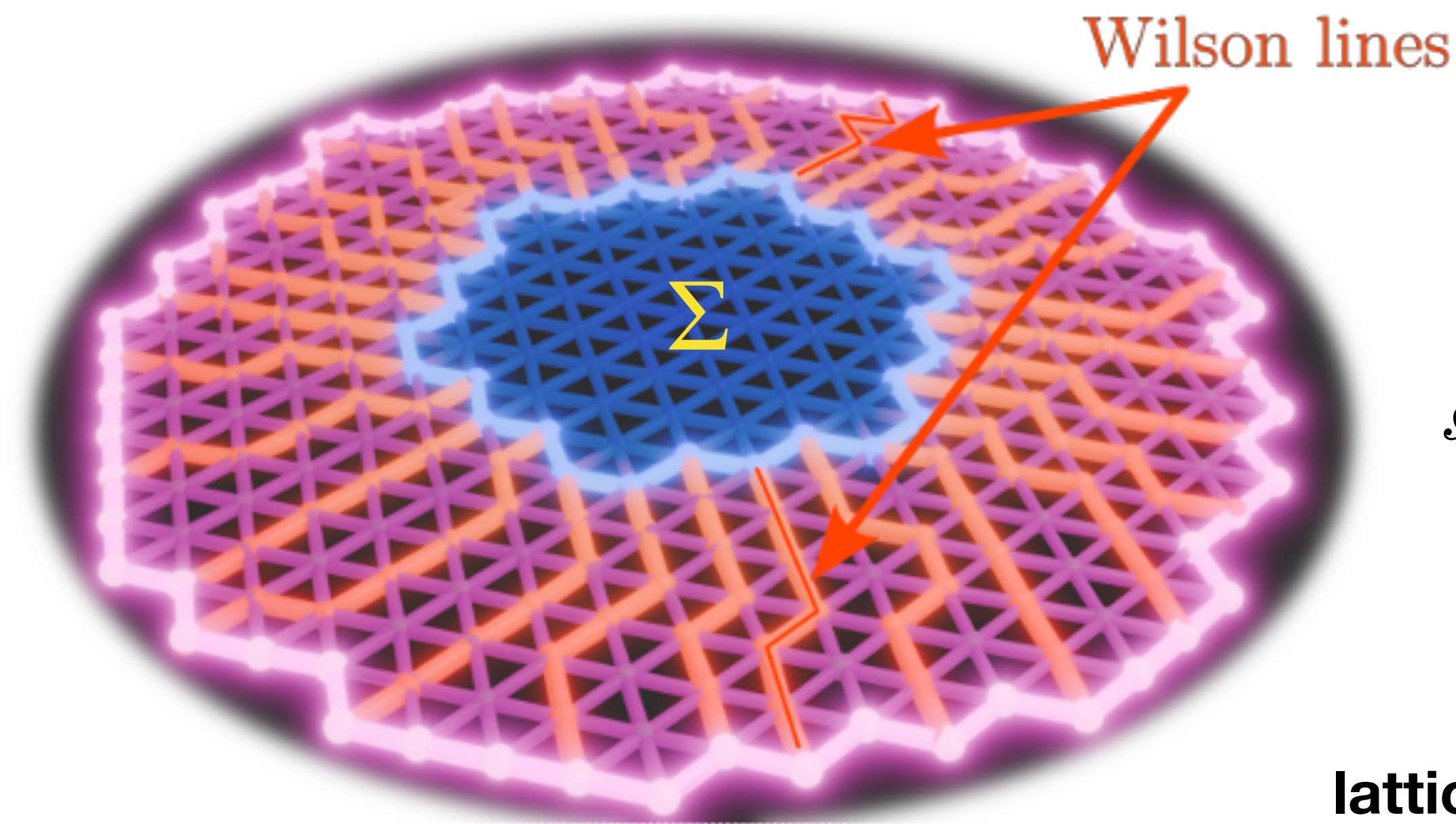
\Rightarrow it is also a **factor**, can show using QRF techniques that indeed

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\Rightarrow relative to extrinsic QRFs get distillable EE!

Extrinsic EE is relational

dress all the regional data in Σ (Wilson lines, electric fields) with extrinsic Wilson line QRF to asymp. bdry



⇒ rel. obs. $O_{|\Phi}^g(a)$ with $a \in \mathcal{A}_\Sigma$

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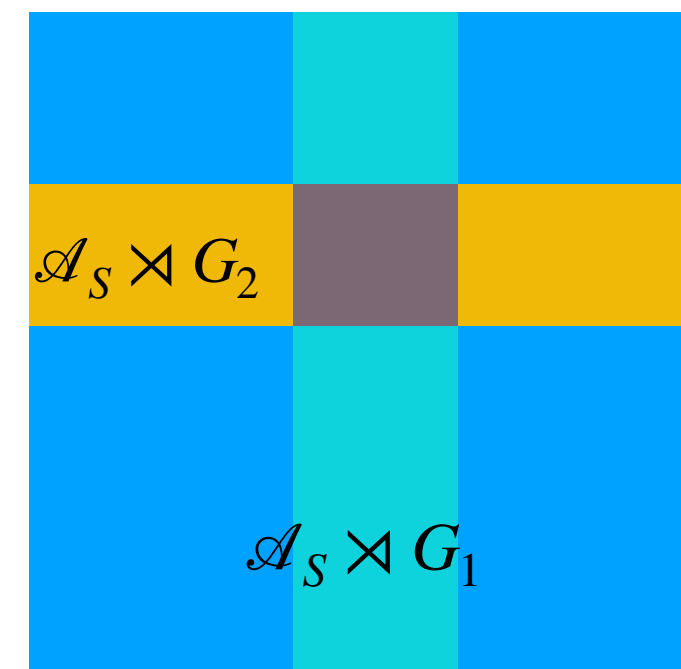
lattice gauge theory analog of crossed product in gravity case

⇒ it is also a **factor**, can show using QRF techniques that indeed

$$\mathcal{H}_{\text{phys}} \simeq \mathcal{H}_{\text{ext}} \otimes \mathcal{H}_{\text{out}}$$

⇒ relative to extrinsic QRFs get distillable EE!

subsystem relativity for diff. ext. QRFs:
distillable EE is QRF-dependent!



3 ways to understand the electric center

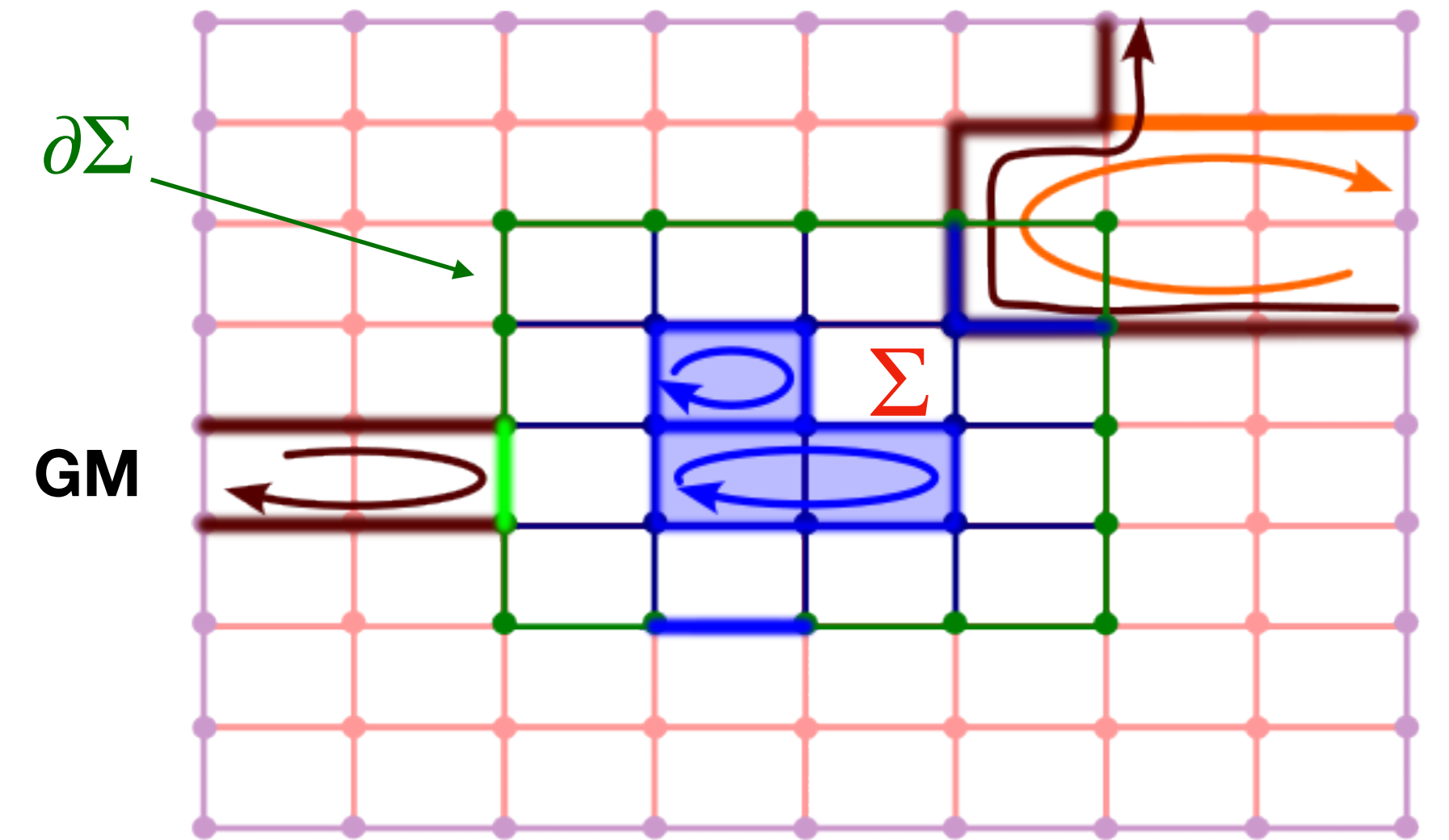
1. \mathcal{A}_{ext} contains a Goldstone mode:

rel. obs. describing int. rel. to ext. QRF $O_{|\Phi}^g(\tilde{\Phi})$

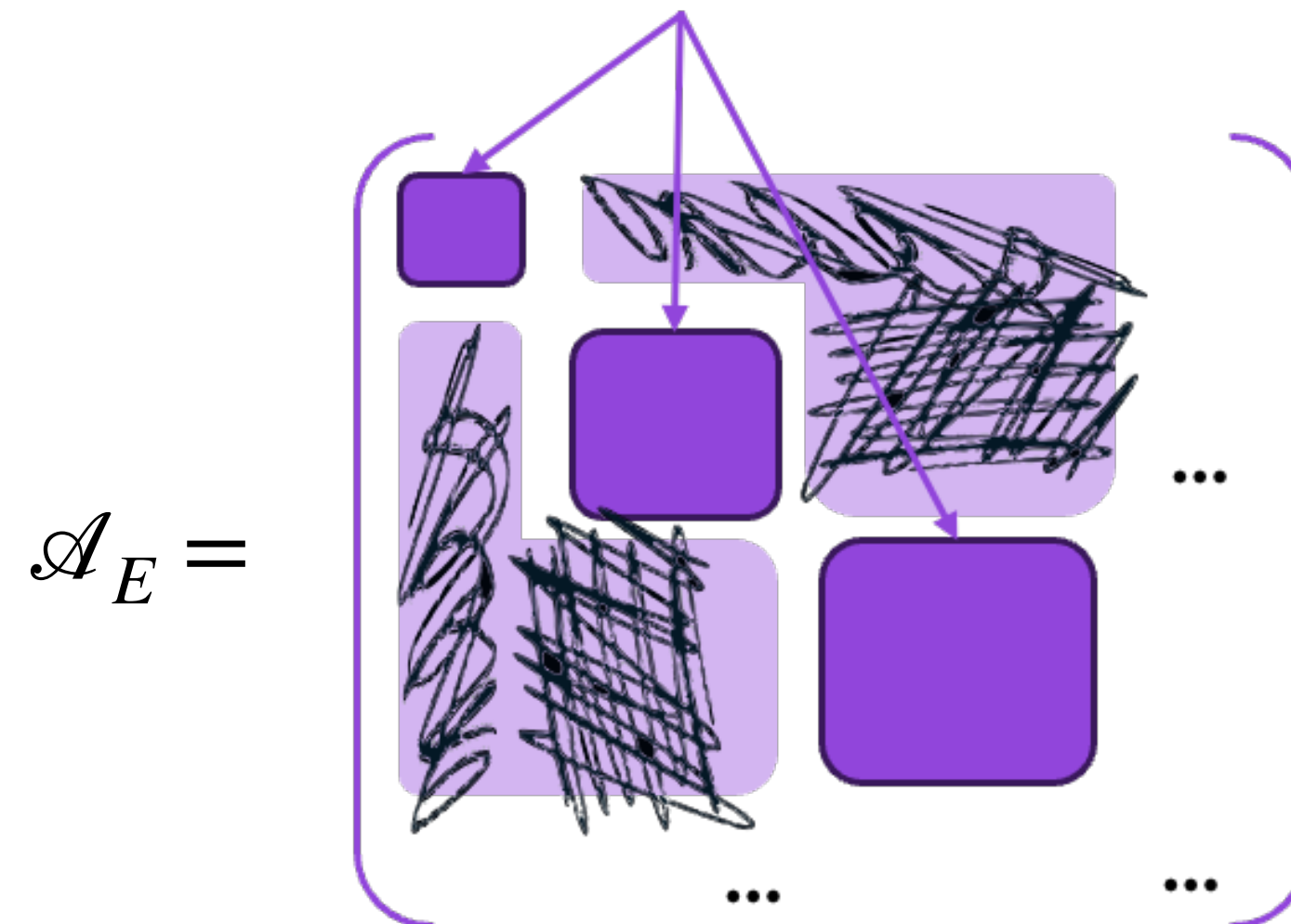
\Rightarrow nonlocal **gauge-inv.** QRF for the electric corner group

if no access to it \Rightarrow average over its orientations/corner group

$$\mathcal{A}_E = \mathbb{G}_{\partial\Sigma}(\mathcal{A}_{\text{ext}}) = \int_{G^{V_{\partial\Sigma}}} dg V_{\Phi}(g) \mathcal{A}_{\text{ext}} V_{\Phi}^{\dagger}(g)$$



Corner group irreps sectors



electric corner charge superselection:

reproduces “electric center” of

[Casini, Huerta, Rosbal '13; Soni, Trivedi '15; Delcamp, Dittrich, Riello '16]

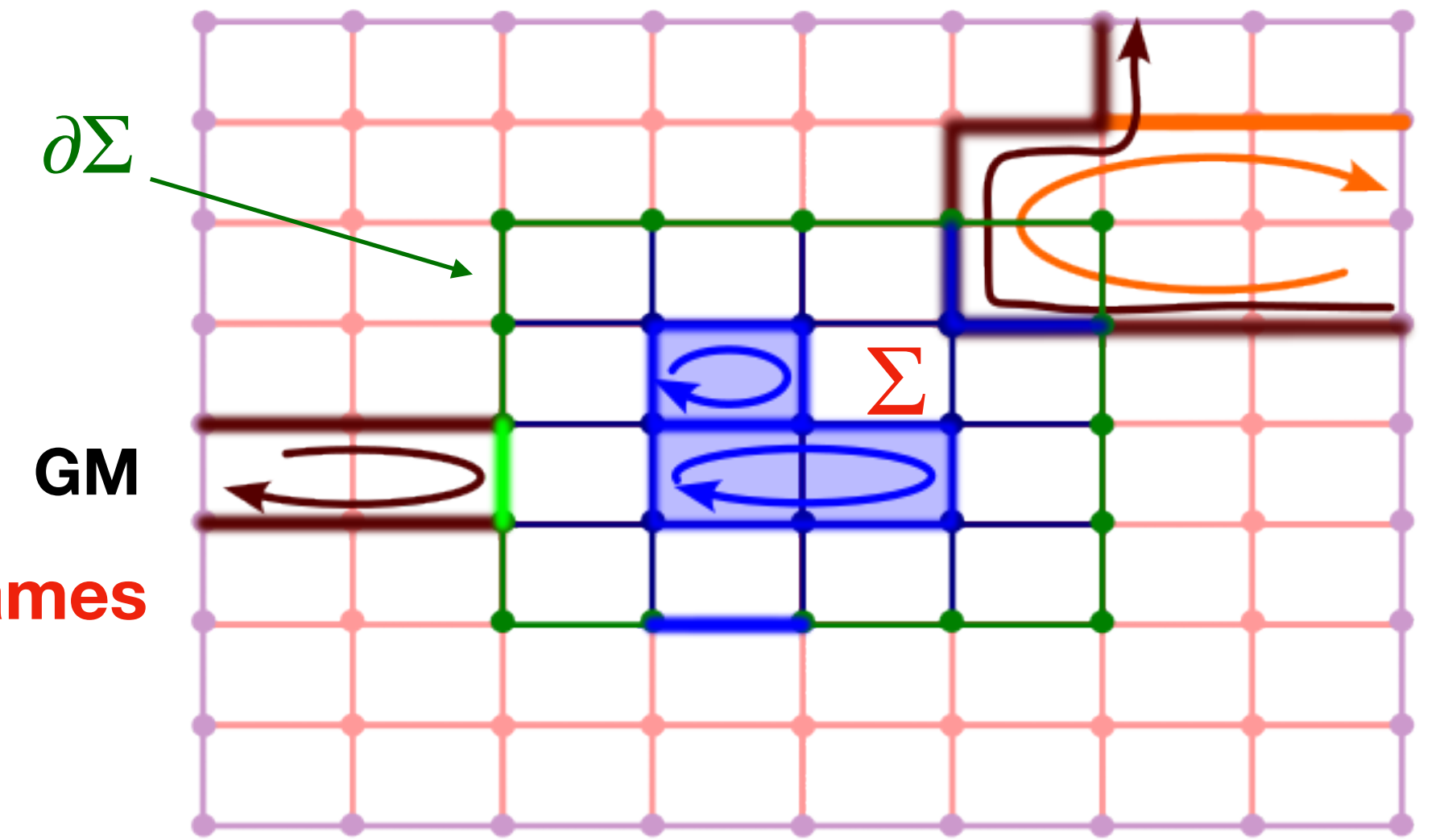
(center: Casimirs of normal flux)

3 ways to understand the electric center

2. what all extr. QRFs agree on:

$$\mathcal{A}_E = \bigcap_{\Phi} \mathcal{A}_{\text{ext}}^{\Phi} \quad (\text{subsystem relativity: agree on what's ext. QRF-indep.})$$

⇒ electric center is to ext. QRFs what Lorentz scalars are to Lorentz frames

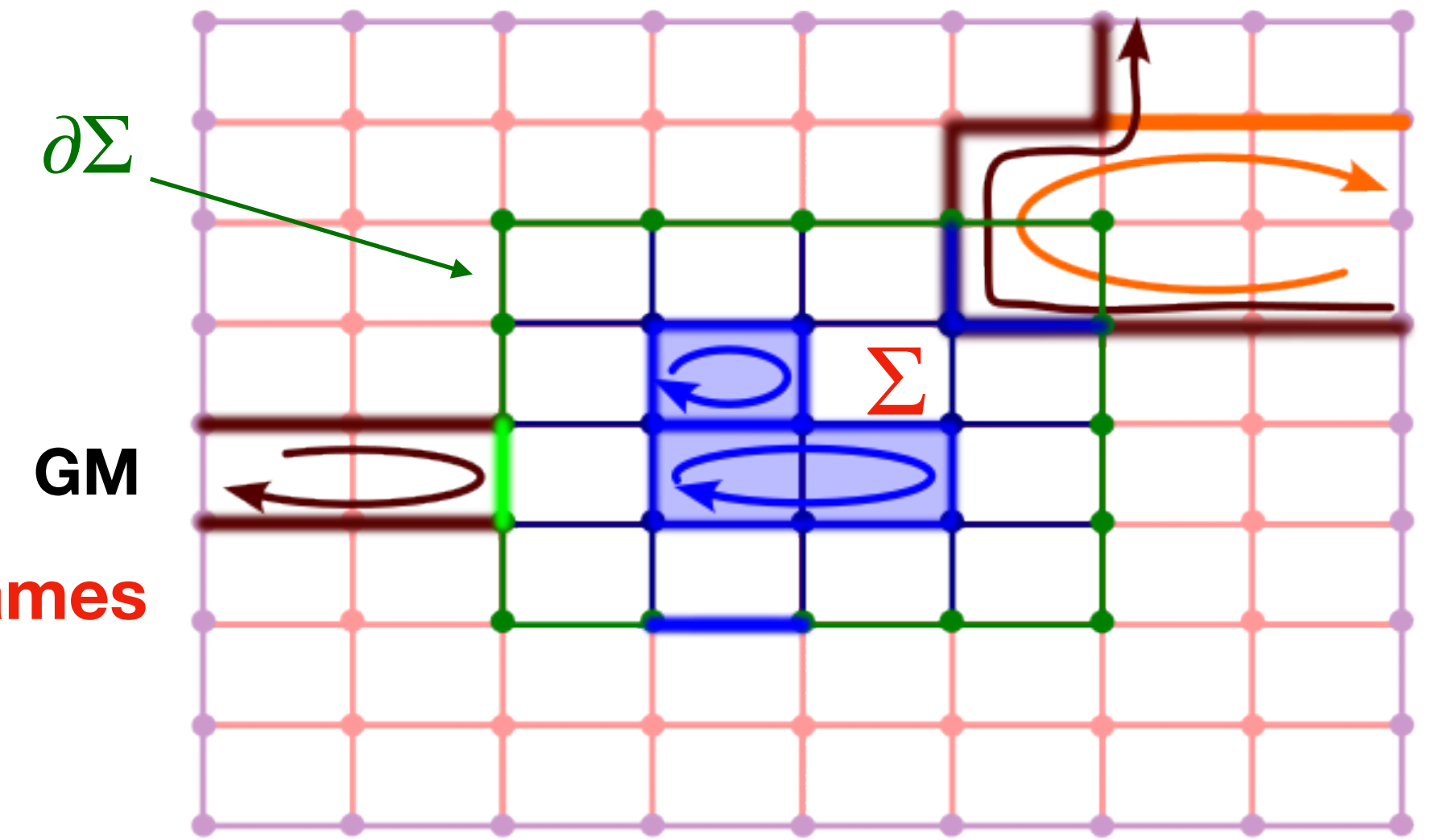


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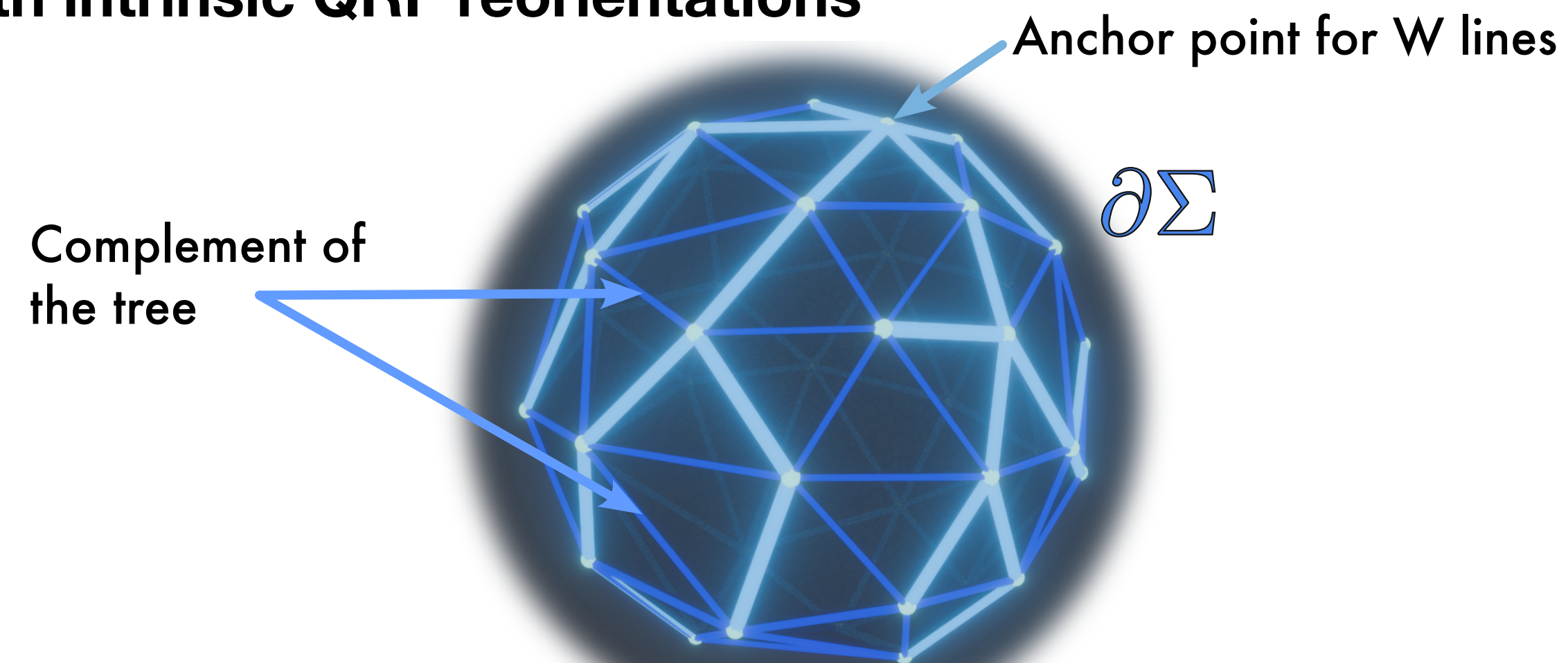
⇒ electric center is to ext. QRFs what Lorentz scalars are to Lorentz frames



3. as a (projected) crossed product of regional complement with intrinsic QRF reorientations

$$\begin{aligned} \mathcal{A}_E &= \mathcal{A}_{\Sigma} \rtimes_{\text{int}} G^{V_{\partial\Sigma}^{-1}} = \left(\mathcal{A}_{\Sigma} \otimes L^2(G^{V_{\partial\Sigma}^{-1}}) \right)^{G^{V_{\partial\Sigma}}} \\ &= \left\langle O_{|\tilde{\Phi}}^e(a), V_{\tilde{\Phi}}(g) \mid a \in \mathcal{A}_{\Sigma \setminus \tilde{\Phi}}, g \in G^{V_{\partial\Sigma}^{-1}} \right\rangle'' \end{aligned}$$

⇒ for **Abelian theories**: electric center IS a crossed product

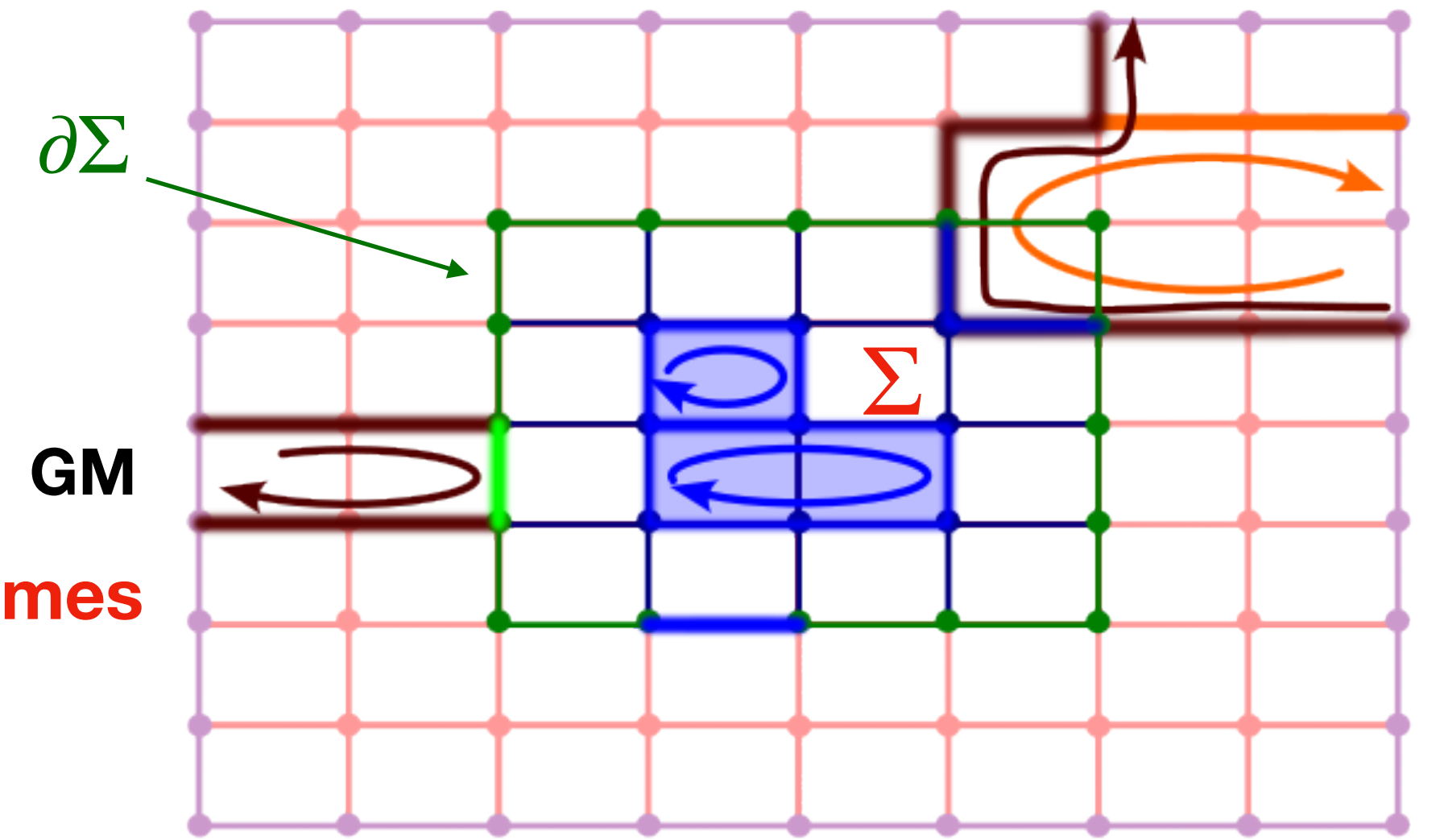


3 ways to understand the electric center

2. what all extr. QRFs agree on:

$$\mathcal{A}_E = \bigcap_{\Phi} \mathcal{A}_{\text{ext}}^{\Phi} \quad (\text{subsystem relativity: agree on what's ext. QRF-indep.})$$

⇒ electric center is to ext. QRFs what Lorentz scalars are to Lorentz frames



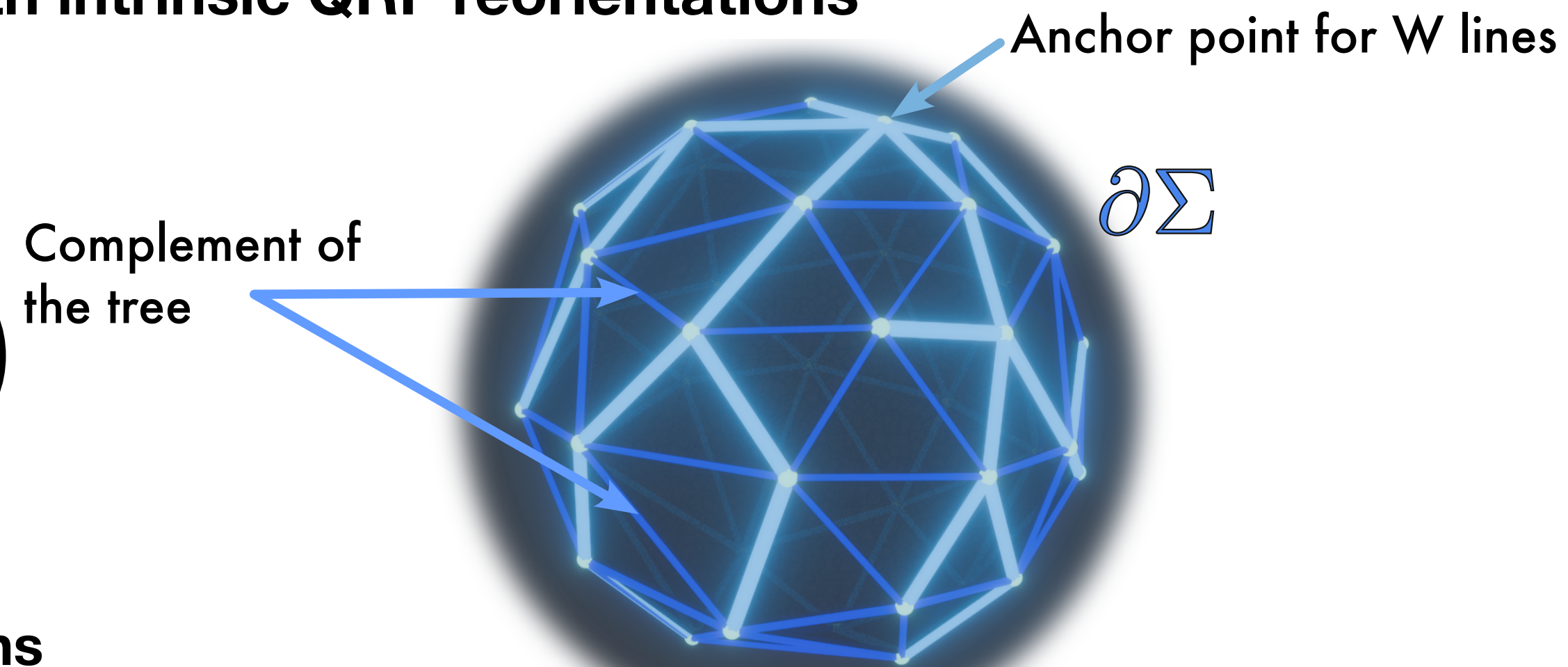
3. as a (projected) crossed product of regional complement with intrinsic QRF reorientations

$$\begin{aligned} \mathcal{A}_E &= \mathcal{G}_N \left(\mathcal{A}_{\Sigma} \rtimes_{\text{int}} G^{V_{\partial\Sigma^{-1}}} \right) = \left(\mathcal{A}_{\Sigma} \otimes L^2(G^{V_{\partial\Sigma^{-1}}}) \right)^{G^{V_{\partial\Sigma}}} \\ &= \mathcal{G}_N \left(\left\langle O_{|\tilde{\Phi}}^e(a), V_{\tilde{\Phi}}(g) \mid a \in \mathcal{A}_{\Sigma \setminus \tilde{\Phi}}, g \in G^{V_{\partial\Sigma^{-1}}} \right\rangle'' \right) \end{aligned}$$

incoherent avg. over gauge transf. at N

$$\mathcal{G}_N(\cdot) = \int_G dg U_N(g) \cdot U_N^\dagger(g)$$

for nA theories:
kills reorientations



Intrinsic relational algebras

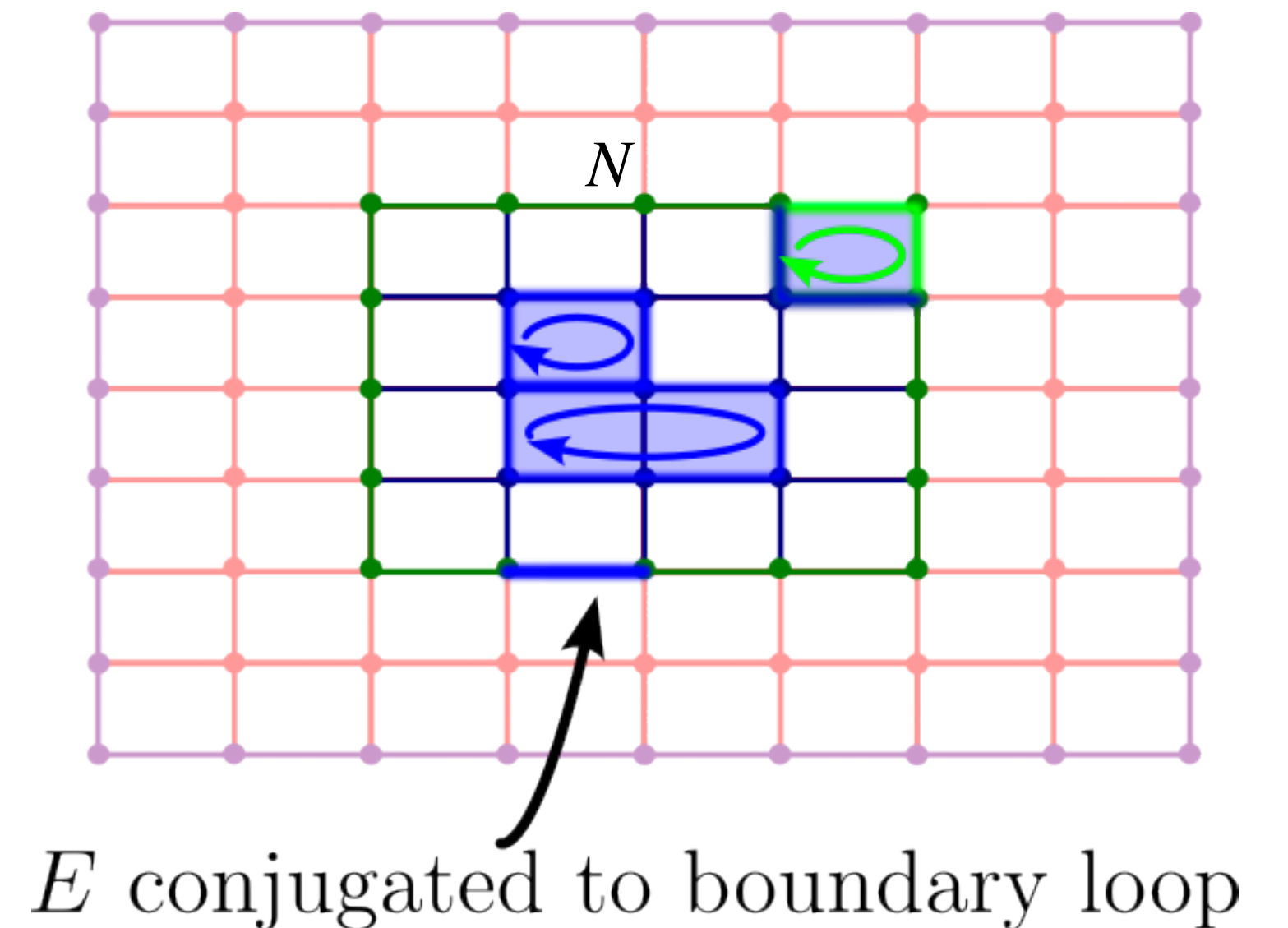
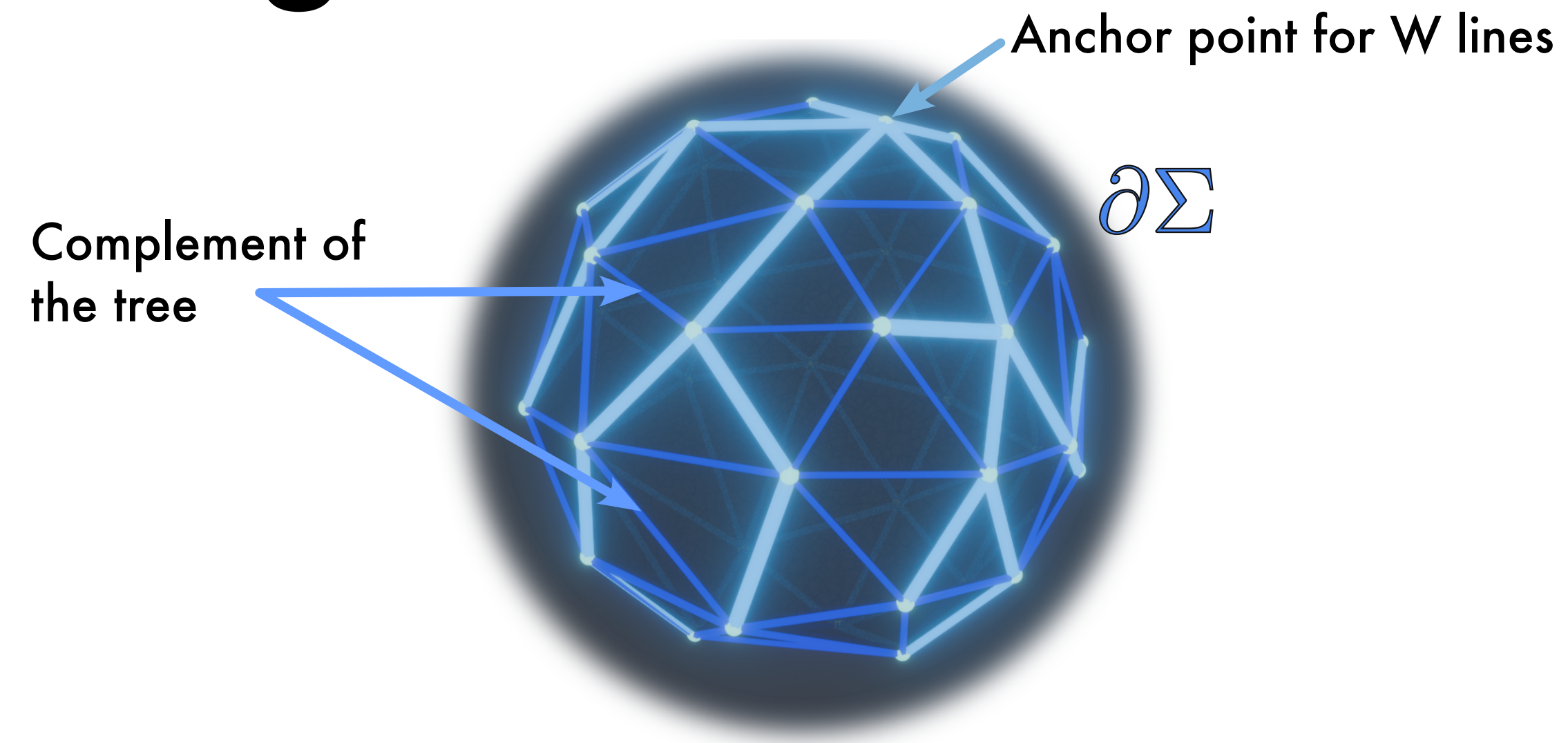
only relational observables of regional complement relative to intrinsic QRF $\tilde{\Phi}$

$$\mathcal{A}_{\text{int}} = \langle O_{|\tilde{\Phi}}^g(a) \mid a \in \mathcal{A}_{\Sigma \setminus \tilde{\Phi}}, g \in G^{V_{\partial\Sigma}} \rangle''$$

(shedding all electric fields tangential to frame tree)

\Rightarrow **factor** in Abelian case \Rightarrow **distillable EE**

\Rightarrow **central** in non-Abelian case \Rightarrow **vN entropy not fully distillable**



2 ways to understand the magnetic center

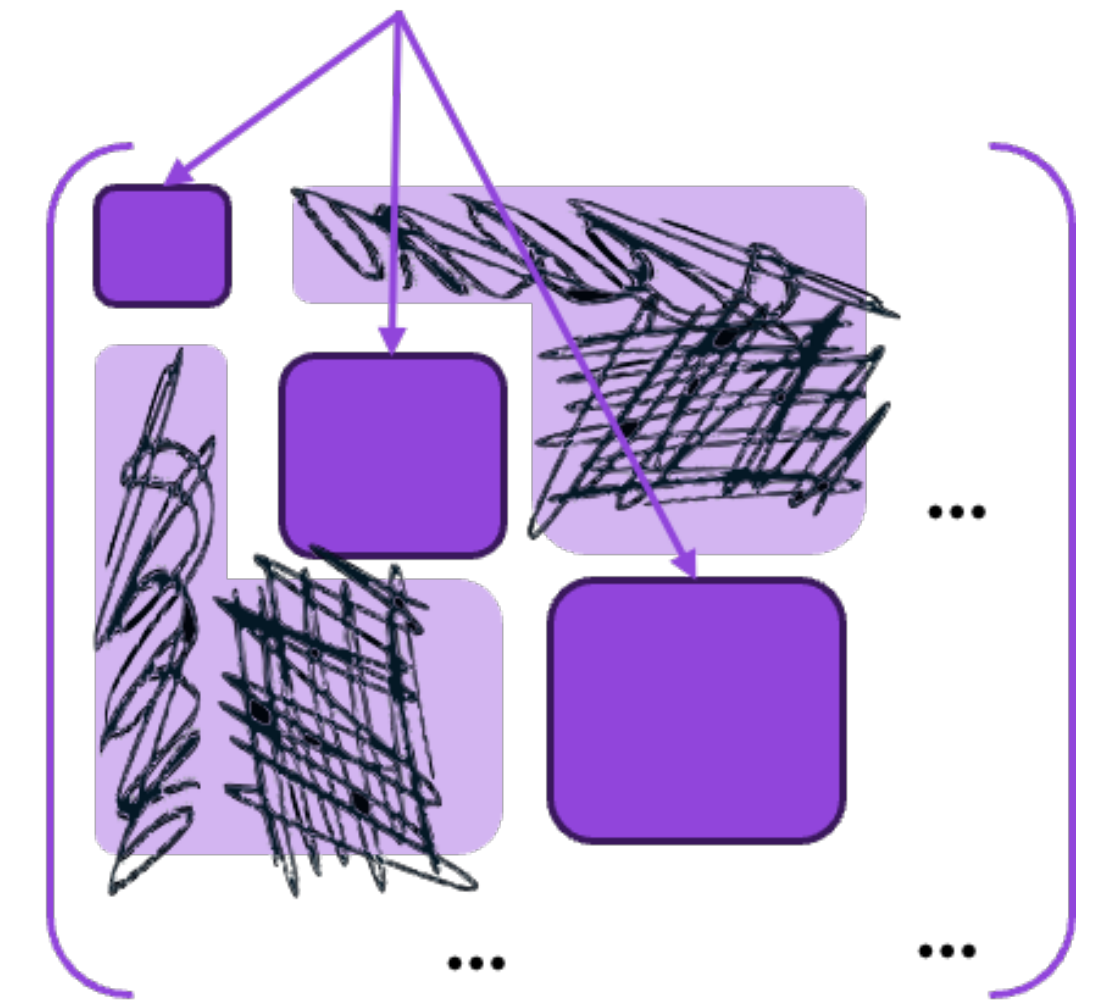
1. in the Abelian case as a twirl over a Pontryagin dual **magnetic** corner symmetry group (generated by corner Wilson loops)

⇒ magnetic corner charge superselection

lattice version of $Q_M[\kappa] = \int_{\partial\Sigma} \kappa F$

$$\mathcal{A}_M = \hat{G}_{\partial\Sigma}(\mathcal{A}_{\text{int}}) =$$

Corner group irreps sectors



2 ways to understand the magnetic center

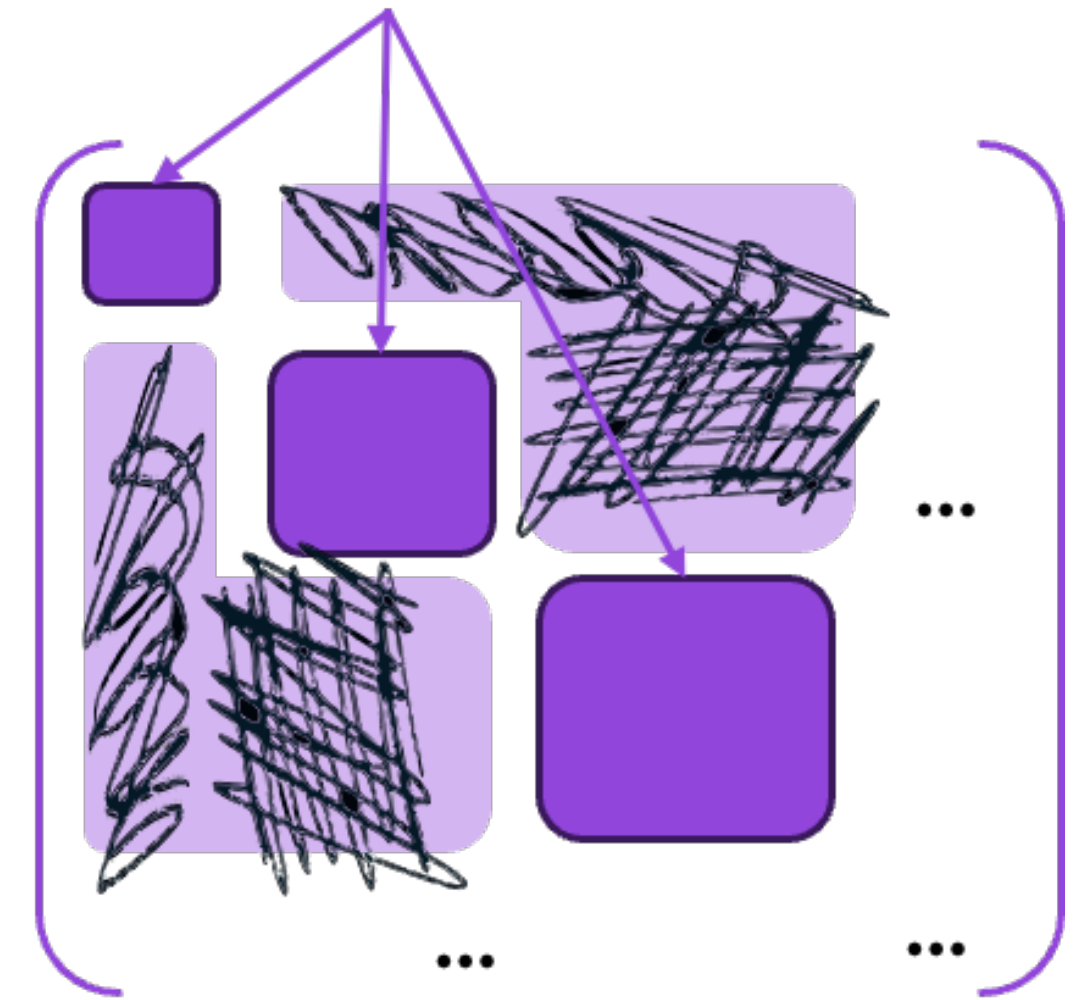
1. in the Abelian case as a twirl over a Pontryagin dual **magnetic** corner symmetry group (generated by corner Wilson loops)

⇒ magnetic corner charge superselection

lattice version of $Q_M[\kappa] = \int_{\partial\Sigma} \kappa F$

$$\mathcal{A}_M = \hat{G}_{\partial\Sigma}(\mathcal{A}_{\text{int}}) =$$

Corner group irreps sectors



2. **in general:** what all intr. QRFs agree on:

$$\mathcal{A}_M = \bigcap_{\tilde{\Phi}} \mathcal{A}_{\text{int}}^{\tilde{\Phi}} \quad (\text{subsystem relativity: agree on what's int. QRF-indep.})$$

⇒ magnetic center is to int. QRFs what Lorentz scalars are to Lorentz frames

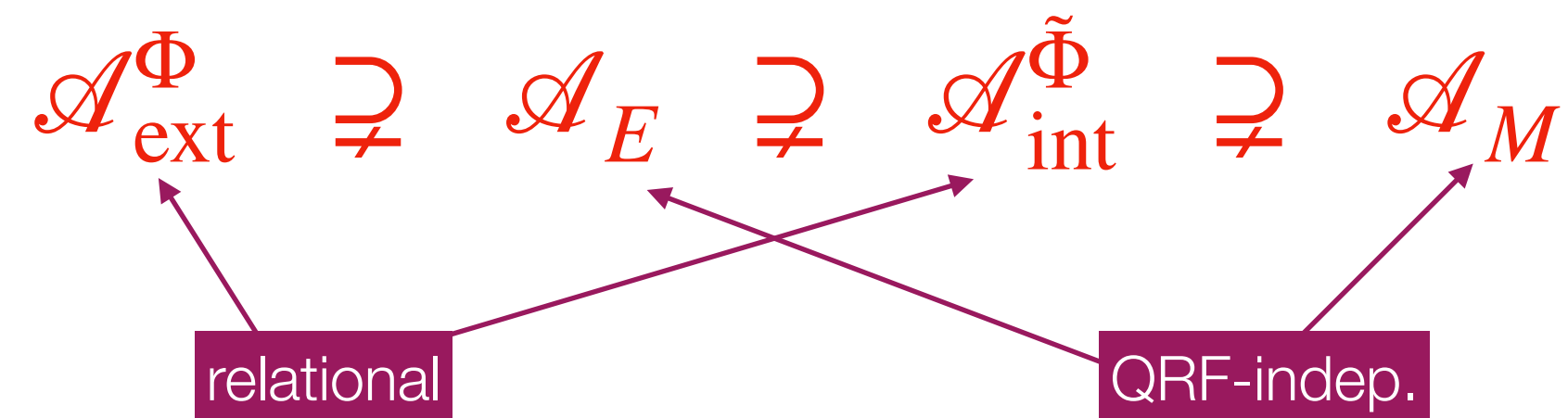
(generalizes magnetic center beyond finite groups, recovers [Casini, Huerta, Rosabal '13; Delcamp, Dittrich, Riello '16] as special cases)

finite Abelian

finite non-Abelian, 2+1d

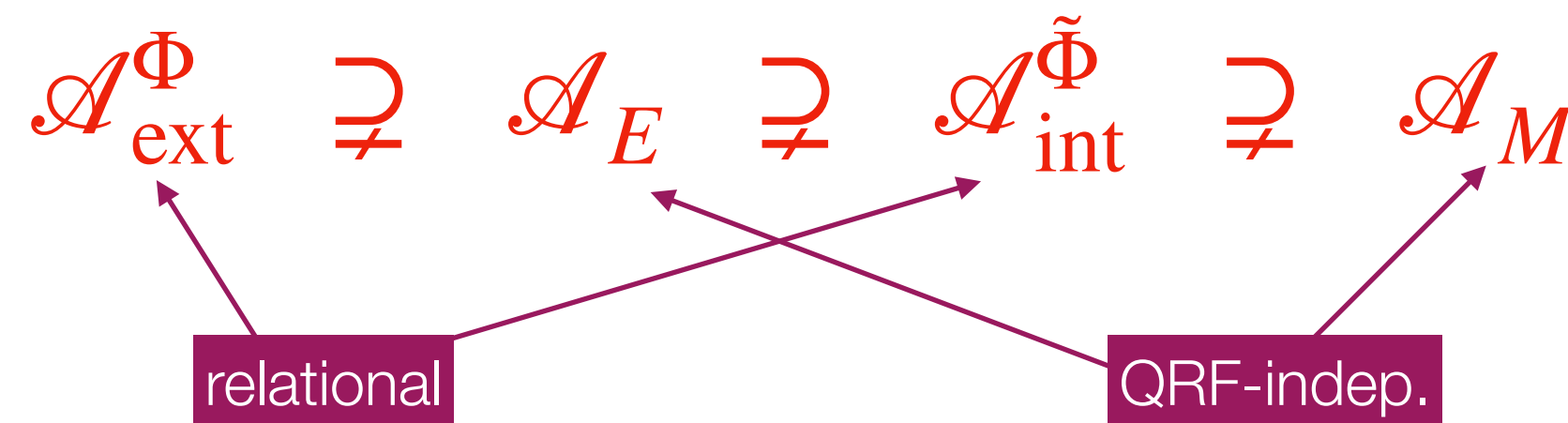
Algebra and entropy hierarchy

algebra	entanglement entropy	distillable
\mathcal{A}_{ext}	$S_{\text{vN}}(\rho_{\Sigma \Phi}^{\text{phys}})$	✓
\mathcal{A}_E	$H(\{p_\epsilon\}) + \sum_\epsilon p_\epsilon (\log d_\epsilon + S_{\text{vN}}(\rho_\epsilon))$	✗
\mathcal{A}_{int}	A: $S_{\text{vN}}(\rho_{\Sigma \setminus \tilde{\Phi} \tilde{\Phi}}^{\text{phys}})$	✓
	nA: $H(\{p_{\hat{\epsilon}}\}) + \sum_{\hat{\epsilon}} p_{\hat{\epsilon}} (\log d_{\hat{\epsilon}} + S_{\text{vN}}(\rho_{\hat{\epsilon}}))$	✗
\mathcal{A}_M	$H(\{p_{\bar{\epsilon}}^m\}) + \sum_{m,\bar{\epsilon}} p_{\bar{\epsilon}}^m (\log d_{\bar{\epsilon}}^m + S_{\text{vN}}(\rho_{\bar{\epsilon}}^m))$	✗



Algebra and entropy hierarchy

algebra	entanglement entropy	distillable
\mathcal{A}_{ext}	$S_{\text{vN}}(\rho_{\Sigma \Phi}^{\text{phys}})$	✓
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\mathcal{A}_{int}	A: $S_{\text{vN}}(\rho_{\Sigma \setminus \tilde{\Phi} \tilde{\Phi}}^{\text{phys}})$	✓
	nA: $H(\{p_{\hat{\epsilon}}\}) + \sum_{\hat{\epsilon}} p_{\hat{\epsilon}} (\log d_{\hat{\epsilon}} + S_{\text{vN}}(\rho_{\hat{\epsilon}}))$	✗
\mathcal{A}_M	$H(\{p_{\bar{\epsilon}}^m\}) + \sum_{m,\bar{\epsilon}} p_{\bar{\epsilon}}^m (\log d_{\bar{\epsilon}}^m + S_{\text{vN}}(\rho_{\bar{\epsilon}}^m))$	✗



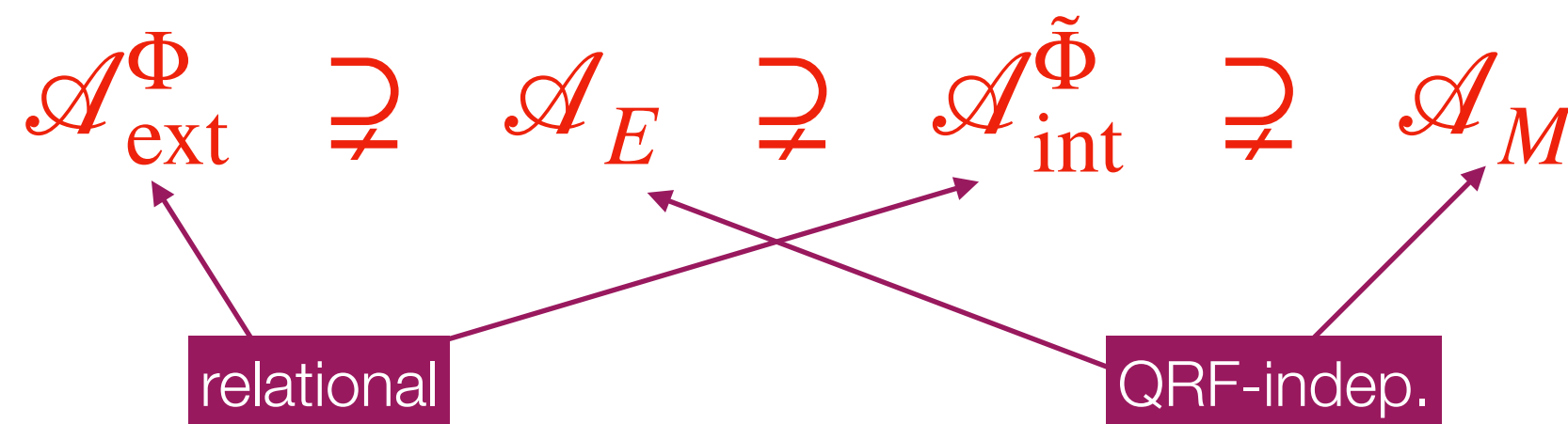
rel. entropy hierarchy:

$$S(\rho_{|\Phi} || \sigma_{|\Phi}) \geq S(\rho_E || \sigma_E) \geq S(\rho_{|\tilde{\Phi}} || \sigma_{|\tilde{\Phi}}) \geq S(\rho_M || \sigma_M)$$

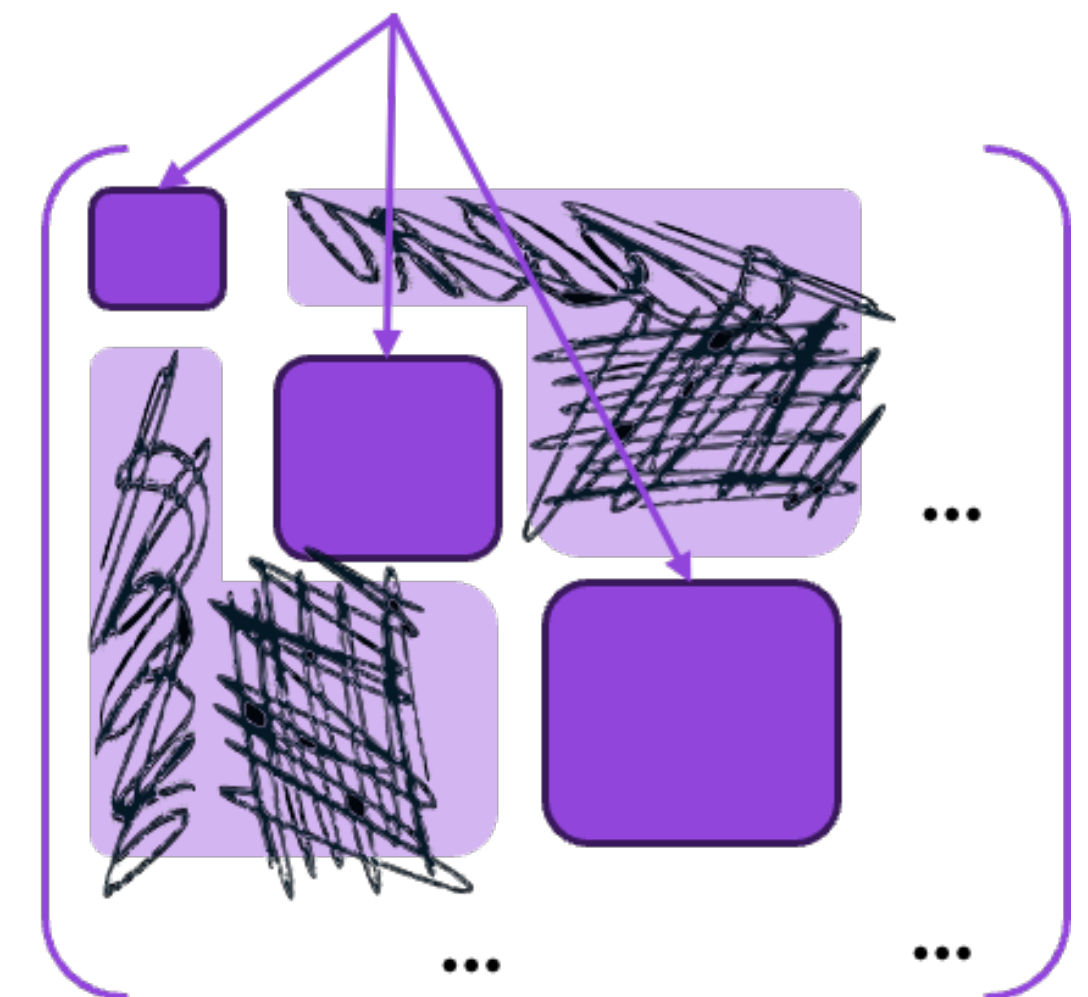
(same global states)

Algebra and entropy hierarchy

algebra	entanglement entropy	distillable
\mathcal{A}_{ext}	$S_{\text{vN}}(\rho_{\Sigma \Phi}^{\text{phys}})$	✓
\mathcal{A}_E	$H(\{p_\epsilon\}) + \sum_\epsilon p_\epsilon (\log d_\epsilon + S_{\text{vN}}(\rho_\epsilon))$	✗
\mathcal{A}_{int}	A: $S_{\text{vN}}(\rho_{\Sigma \setminus \tilde{\Phi} \tilde{\Phi}}^{\text{phys}})$	✓
	nA: $H(\{p_{\hat{\epsilon}}\}) + \sum_{\hat{\epsilon}} p_{\hat{\epsilon}} (\log d_{\hat{\epsilon}} + S_{\text{vN}}(\rho_{\hat{\epsilon}}))$	✗
\mathcal{A}_M	$H(\{p_{\bar{\epsilon}}^m\}) + \sum_{m,\bar{\epsilon}} p_{\bar{\epsilon}}^m (\log d_{\bar{\epsilon}}^m + S_{\text{vN}}(\rho_{\bar{\epsilon}}^m))$	✗



Corner group irreps sectors



rel. entropy hierarchy:

$$S(\rho_{|\Phi} || \sigma_{|\Phi}) \geq S(\rho_E || \sigma_E) \geq S(\rho_{|\tilde{\Phi}} || \sigma_{|\tilde{\Phi}}) \geq S(\rho_M || \sigma_M)$$

EE inequalities from incoherent corner twirls

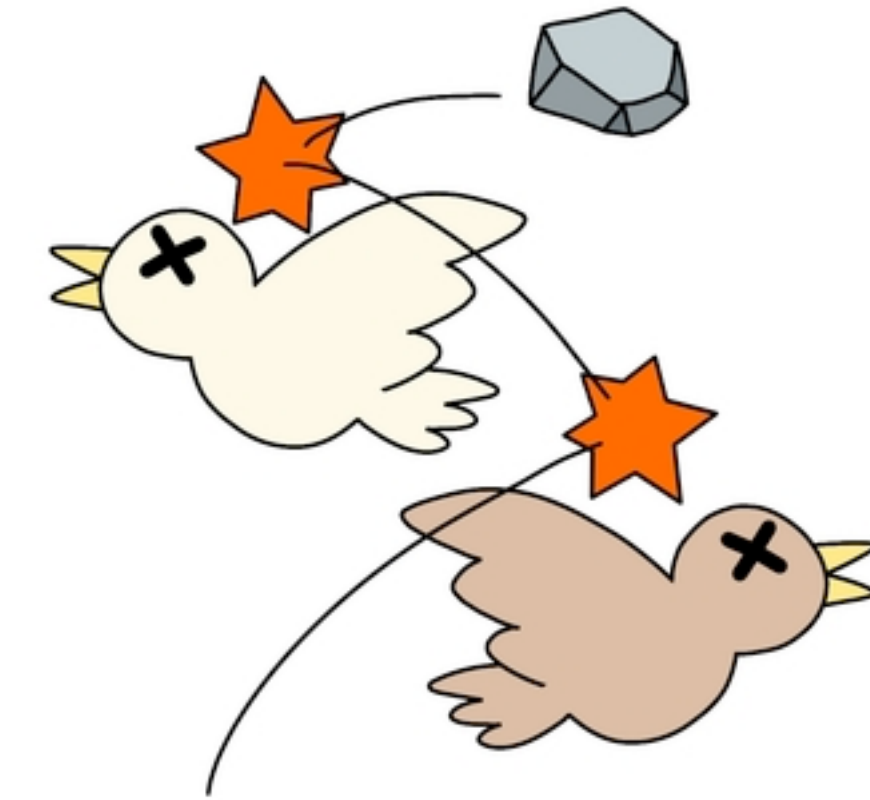
(charge superselection)

$$S_{\text{vN}}(\rho_{|\Phi}) \leq S_{\text{vN}}(\rho_E)$$

$$S_{\text{vN}}(\rho_{|\tilde{\Phi}}) \leq S_{\text{vN}}(\rho_M)$$

finite Abelian groups

Summary



- **QRFs help with intrinsically regulating UV divergences in EEs**
(von Neumann algebra type transition)
- **multitude of possible gauge-inv. subsystems associated with a subregion in gauge theories**
(algebra and entropy hierarchy)
- **QRFs perspectives as gauge-inv. subsystem partitions**
(may be factor algebras \Rightarrow distillable entanglement)
- **relational entanglement entropies**
(relational entropies are QRF/observer-dependent)
- **electric/magnetic center algebra is what all ext./int. QRFs agree on**