# New Results on Spin Density Matrix Elements for $\boldsymbol{\rho}^{0}$ at Hermes 

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## Outline

- Definition of Spin Density Matrix Elements (SDMEs)
- Angular Distribution of Vector Meson Decay
- The Hermes Data
- Method of SDMEs extraction
- 23 Extracted SDMEs
- Kinematical dependences of the SDMEs
- Summary


## Exclusive, Diffractive Electroproduction of $\boldsymbol{\rho}^{0}$

$$
\mathbf{e}+\mathbf{N} \rightarrow \mathbf{e}^{\prime}+\mathbf{N}+\boldsymbol{\rho}^{\mathbf{0}}
$$



$$
\begin{aligned}
& Q^{2}=-q^{2}=-\left(k-k^{\prime}\right)^{2} \\
& W^{2}=(q+p)^{2} \\
& t=(q+v)^{2}
\end{aligned}
$$

- Two gluon exchange mechanism at higher energies
- Quark exchange mechanism at intermediate energies
- Spin structure of the $\rho^{0}$ production
(How the helicity of the of the $\rho^{0}$ meson is related to the helicity of the virtual photon)
- Spin Density Matrix Elements SDMEs


## Spin-Density Matrix of the Vector Meson

- $\rho(\mathrm{V})=\frac{1}{2} \mathrm{~T} \rho(\gamma) \mathrm{T}^{+} \begin{aligned} & - \text { Spin-density matrix of the vector meson } \rho(\mathrm{V}) \text { in } \\ & \text { terms of the photon spin density matrix } \rho(\gamma) \text { and }\end{aligned}$ helicity amplitude T
- $\rho_{\lambda_{v} \lambda_{v^{\prime}}}^{\alpha}=\frac{1}{2 N_{\alpha}} \sum_{\lambda_{r} \lambda_{\gamma}} T_{\lambda_{v} \lambda_{\gamma}} \Sigma_{\lambda_{r} \lambda_{r}^{\prime} \lambda_{r}}^{\alpha} T_{\lambda_{v} \lambda_{r}^{\prime} \lambda_{r}^{\prime}}^{*}-\begin{aligned} & \text { spin-density matrix elements of } \\ & \text { the vector meson }\end{aligned}$
$T_{\lambda_{\mathrm{V}} \lambda_{N}^{\prime} \cdot \lambda_{\gamma} \lambda_{\mathrm{N}}}=\left\langle\lambda_{\mathrm{V}} \lambda_{N}^{\prime}\right| \mathrm{J}^{(\mathrm{em})} e^{\left(\lambda_{\gamma}\right)}\left|\lambda_{\mathrm{N}}\right\rangle \quad$ - helicity amplitudes
where $\lambda_{\mathrm{V}}, \lambda_{\gamma}, \lambda_{\mathrm{N}}-$ helicity of the vector meson, photon and proton
$\mathrm{J}^{(\mathrm{em})}$ - electromagnetic current, $\mathrm{e}^{(\lambda)}$ - photon polarization vector
$\lambda_{\gamma}=0$ - longitudinal polarization, $\lambda_{\gamma}= \pm 1-$ transverse polarization
$\sum_{\lambda_{\gamma} \lambda_{\gamma}^{\prime}}^{\alpha}-\begin{aligned} & (\alpha=0, \ldots 8) \text { nine hermitian matrices representing states of photon } \\ & \text { polarization }\end{aligned}$ $\alpha=0$ - unpolarized transverse photon $\alpha=1,2$ - linear polarization $\alpha=3$ - circularly polarized photon $\quad \alpha=4-$ longitudinal photon $\alpha=5,6,7,8-$ longitudinal- transverse interference terms


## Spin Density Matrix Elements (SDMEs)

- It is not possible to separate contributions from longitudinal and transverse photon at constant beam energy.
- We measure SDMEs - $\mathrm{r}_{\lambda \lambda_{\mathrm{v}}}^{\alpha}$

$$
\begin{aligned}
& \mathrm{r}_{\lambda_{V} \lambda_{V}^{\prime}}^{04}=\frac{\rho_{\lambda_{V} i_{V}}^{0}+\varepsilon \mathrm{R} \rho_{\lambda_{V} i_{V}^{\prime}}^{4}}{1+\varepsilon \mathrm{R}} \quad \mathrm{R}=\frac{\sigma_{L}}{\sigma_{T}} \quad \varepsilon-\text { polarization parameter } \\
& \mathrm{r}_{\lambda_{V} \lambda_{V}^{\prime}}^{\alpha}=\frac{\rho_{\lambda_{V} \lambda_{V}}^{a}}{1+\varepsilon \mathrm{R}} \quad \alpha=1,2,3 \\
& \mathrm{r}_{\lambda_{V} \lambda_{V}^{\prime}}^{\alpha}=\sqrt{R} \frac{\rho_{\lambda_{V} \lambda_{V}}^{a}}{1+\varepsilon \mathrm{R}} \quad \alpha=5,6,7,8
\end{aligned}
$$

- SCHC - s-channel helicity conservation
helicity of the virtual photon $=$ helicity of the vector meson

$$
\begin{aligned}
& \mathrm{T}_{01}=\mathrm{T}_{10}=\mathrm{T}_{-10}=\mathrm{T}_{0-1}=\mathrm{T}_{-11}=\mathrm{T}_{1-1}=0 \\
& \mathrm{~T}_{00} \neq 0, \mathrm{~T}_{11} \neq 0 \mathrm{~T}_{-1-1} \neq 0 \\
& r_{00}^{04}, \operatorname{Re}\left\{r_{1-1}^{1}\right\}, \operatorname{Im}\left\{r_{1-1}^{2}\right\}, \operatorname{Re}\left\{r_{10}^{5}\right\}, \operatorname{Im}\left\{r_{10}^{6}\right\}, \operatorname{Im}\left\{r_{10}^{7}\right\}, \operatorname{Re}\left\{r_{10}^{8}\right\} \neq 0
\end{aligned}
$$

- NPE - Natural Parity Exchange process dominance the exchange particle have quantum numbers $\left(\mathrm{J}^{\mathrm{P}}=0^{+}, 1^{-}, 2^{+} \ldots\right)$

$$
\mathrm{T}_{00}, \mathrm{~T}_{11}=\mathrm{T}_{-1-1}, \mathrm{~T}_{01}=-\mathrm{T}_{0-1}, \mathrm{~T}_{10}=-\mathrm{T}_{-10}, \mathrm{~T}_{1-1}=\mathrm{T}_{-11}
$$

## Decay Angles Definition


$\gamma^{*} \mathrm{p}$ - center - of momentum frame $\Phi$ - the azimuthal production angle of $\rho^{0}$ meson

$\rho^{0}$ - rest frame
$\theta, \varphi$ - polar and azimuthal decay angle of the meson $\pi^{+}$relative to the $\rho^{0}$ spin quantization axis, which is along the direction oposite the direction of the recoiling target -p '

## Decay Angular Distribution in terms of SDMEs

$-\epsilon \sin 2 \Phi\left(\sqrt{2} \operatorname{Im}\left(r_{10}^{2}\right) \sin ^{2} \Theta \sin \phi+\operatorname{Im}\left(r_{1-1}^{2}\right) \sin 2 \Theta \sin 2 \phi\right)$
$+\sqrt{2 \epsilon(1+\epsilon)} \cos \Phi\left(r_{11}^{5} \sin ^{2} \Theta+r_{00}^{5} \cos ^{2} \Theta-\sqrt{2} R e r_{10}^{5} \sin 2 \Theta \cos \phi-\right.$ $\left.r_{1-1}^{5} \sin ^{2} \Theta \cos 2 \phi\right)$
$\left.+\sqrt{2 \epsilon(1+\epsilon)} \sin \Phi\left(\sqrt{2} \operatorname{Im}\left(r_{10}^{6}\right) \sin 2 \Theta \sin \phi+\operatorname{Im}\left(r_{1-1}^{6}\right) \sin ^{2} \Theta \sin 2 \phi\right)\right]$
$W^{\text {long.pol. } .}(\cos \Theta, \phi, \Phi)=\frac{3}{4 \pi} P_{\text {beam }}[$
$\sqrt{1-\epsilon^{2}}\left(\sqrt{2} \operatorname{Im}\left(r_{10}^{3}\right) \sin 2 \Theta \sin \phi+I m\left(r_{1-1}^{3}\right) \sin ^{2} \Theta \sin 2 \phi\right)$
$+\sqrt{2 \epsilon(1-\epsilon)} \cos \Phi\left(\sqrt{2} \operatorname{Im}\left(r_{10}^{7}\right) \sin 2 \Theta \sin \phi+\operatorname{Im}\left(r_{1-1}^{7}\right) \sin ^{2} \Theta \sin 2 \phi\right)$
$+\sqrt{2 \epsilon(1-\epsilon)} \sin \Phi\left(r_{11}^{8} \sin ^{2} \Theta+r_{00}^{8} \cos ^{2} \Theta-\sqrt{2} R e\left(r_{10}^{8}\right) \sin 2 \Theta\right.$
$\left.\left.\cos \phi-r_{1-1}^{8} \sin ^{2} \Theta \cos 2 \phi\right)\right]()$

15 unpolarized SDMEs

8 polarized SDMEs

## Information about Hermes Experimental Data

- Polarized positron (electron) beam of energy $\mathrm{E}=27.6 \mathrm{GeV}$
- The average lepton beam polarization was 0.53 for both positive and negative beam helicities
- Targets: Hydrogen, Deuterium
- Data collected in years 1996-2000


## Selection of Diffractive Exlusive $\rho^{0}$ Events

- Event has only 3 tracks, scattered lepton and two pions $\pi+\pi-$
- The $\rho^{0}$ meson is selected by mass constraints

$$
0.6<\mathrm{M}_{\pi+\pi-}<1.0 \mathrm{GeV}
$$

and veto constraints $\mathrm{M}_{\mathrm{K}+\mathrm{K}_{-}}>1.06 \mathrm{GeV}$

- Diffractive events were selected by requiring $-\mathrm{t}^{\prime}=\mathrm{t}-\mathrm{t}_{\text {min }}<0.6 \mathrm{GeV}$
- Exlusive events $-1<\delta E=\frac{M_{x}^{2}-M_{\text {targ }}^{2}}{2 \mathrm{M}_{\text {targ }}}<0.6 \mathrm{GeV}$
- 9600 - events H, 16000 - events D

$\delta$ E distributions for exlusive diffractive $\rho^{0}$ production for different kinematical bins (circles), compared to SIDIS background calculated by PYTHIA MC (histogram )


## Extraction of SDMEs

- SDMEs were determined by minimizing the difference between 3-dimesional matrix of data and a sample of MC events.
(1) 3-dimensional matrix of data in variables $(\cos (\theta), \varphi, \Phi)$ binned in $(8,8,8)$ bins
(2) 3-dimensional matrix of background events
(3) 3-dimensional matrix of MC events generated with uniform angular distribution, reweighted with angular distribution function $\mathrm{W}(\cos (\theta), \varphi, \Phi)$ which depends on the SDMEs
(1) - (2) was fitted by (3) with a binned Maximum Likelihood Method where SDMEs were treated as free parameters.


## Fitted Angular Distribution



$\psi=\varphi-\Phi$


- Closed circles represent measured data
- MC distribution fitted to the data
- isotropically genereted events used as an input for the fits


## 23 Unpolarized and Polarized SDME on Hydrogen and Deuterium



## The $\mathbf{Q}^{2}$ - dependence of the $\mathbf{1 5}$ upolarized SDMEs



## The -t'- dependence of the $\mathbf{1 5}$ unpolarized SDME



## Test of NPE dominance

$$
1-r_{00}^{04}+2 r_{1-1}^{04}-2 r_{11}^{1}-2 r_{1-1}^{1}=0 \quad \text { For NPE }
$$




## Longitudinal-to-Transverse Cross-Section Ratio



$$
\mathrm{R}_{\rho}^{\text {SCHC }}=\frac{1}{\varepsilon} \times \frac{\mathrm{r}_{00}^{04}}{1-\mathrm{r}_{00}^{04}}
$$

## Summary

$\square 23$ SDMEs were obtained with the Likelihood method for $\rho^{0}$ production on proton and deuteron targets.
$\square$ No significant deviation is seen between the SDMEs from proton and deuteron data and their kinematic dependences.
$\square$ Violation of SCHC was shown for non-zero values of several SDMEs on hydrogen and deuterium.

- 15 unpolarized SDMEs were extracted for four $\mathrm{Q}^{2}$ bins and four $-t$ ' bins for proton and deuteron. Several clean kinematic dependences of SDMEs on $\mathrm{Q}^{2}$ and -t' are observed.
$\square$ Test of Natural Parity Exchange was performed for different kinematic bins. An indication of unnatural parity exchange amplitudes is seen in the proton data.
$\square \mathrm{R}^{\text {SCHC }}$ was obtained for four $\mathrm{Q}^{2}$ bins under the assumption of SCHC

