

**CONTROVERSIAL ISSUES IN THE  
POLARIZED PARTON DENSITIES:**

$\Delta G, \Delta_s$ , positivity and higher twist

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Two main problems compared with  
**UNPOLARIZED** case:

a) Small range of  $Q^2$   
 $\Rightarrow$  **poor determination of  $\Delta G(x)$**

b) No neutrino and antineutrino data  
 $\Rightarrow$  **poor flavour separation** and  
 $\Rightarrow$  **can only measure  $\Delta q + \Delta \bar{q}$**

### 'REMEDIES'

1) Include data at lower  $Q^2$  via **HIGHER TWIST** terms

$$\begin{aligned} g_1^{expt} &= F_1^{expt} A_1^{expt} \\ &= \frac{F_2^{expt}}{2x(1 + R^{expt})} A_1^{expt} \\ &= g_1^{QCD} + h/Q^2 \end{aligned}$$

HT essential. [ Leader, Sidorov, Stamenov  
(LSS05)include HT]

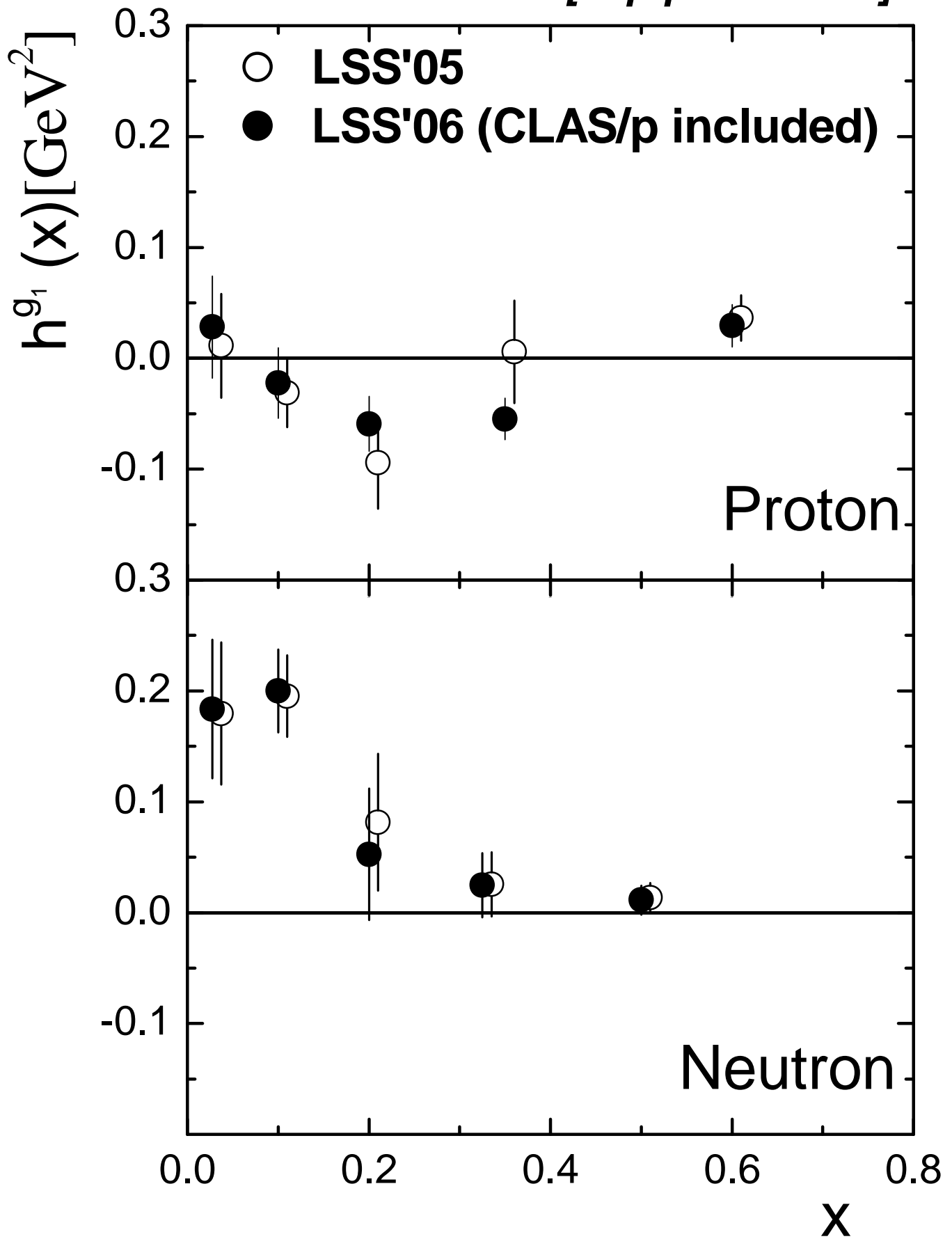
[Blümlein, Böttcher (BB02) claim HT not  
necessary; Hirai, Kumano, Saito (AAC06) do  
not include HT]

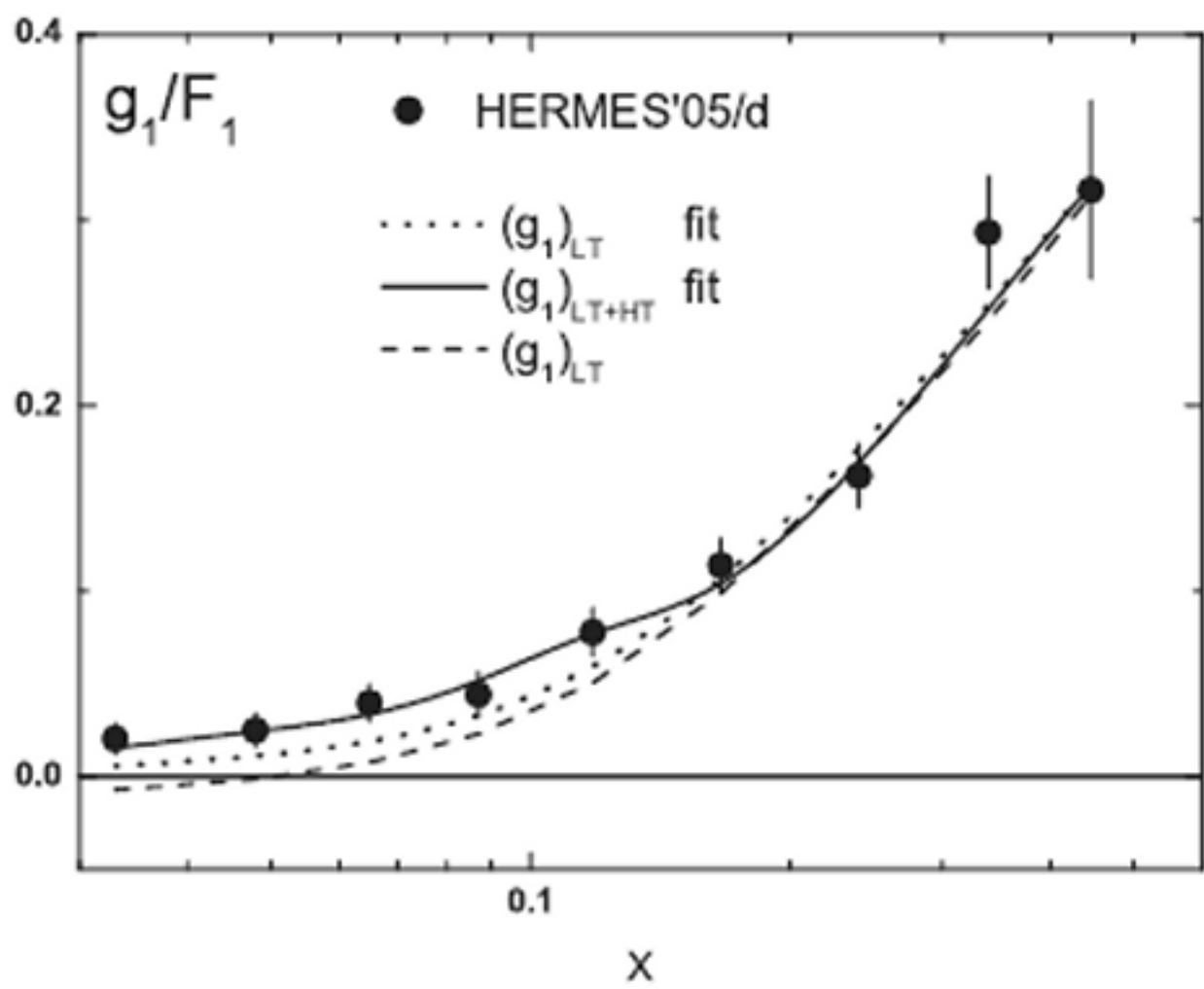
An example: HERMES  $g_1^d$  data at low  
 $Q^2 = 1\text{GeV}^2$  compared with COMPASS data  
at  $6\text{GeV}^2$

AAC06 explain difference as possibly due to  
gluon.

LSS05 claim due to HT. See Fig. 1 and Fig. 2

*LSS'05 [hep-ph/0512114]*





## 2) Impose $SU(3)$ sum rule

Notation:

$$\Delta q \equiv \Delta q(Q^2) = \int_0^1 dx \Delta q(x, Q^2)$$

$$\begin{aligned} a_8 &\equiv \Delta u + \Delta \bar{u} + \Delta d + \Delta \bar{d} \\ &\quad - 2(\Delta s + \Delta \bar{s}) \\ &= 3F - D \\ &= 0.585 \pm 0.025 \end{aligned}$$

Leader, Stamenov (2003) showed, depending on which data is used, E155 or E143, that

**IF  $(\Delta s + \Delta \bar{s}) \geq 0$ , then**

$$a_8 \leq 0.089 \pm 0.058 \quad (1)$$

or

$$a_8 \leq 0.197 \pm 0.068 \quad (2)$$

But analysis of hyperon decays [Ratcliffe (1999)] implies

$$a_8 = 0.585 \pm 10\%$$

New analysis of SIDIS [ de Florian, Navarro, Sassot (2006)(deFNS06)]implies

$$a_8 = 0.585 - 8\% \text{ or } -12\%$$

depending on choice of FRAGMENTATION functions.

These values significantly contradict the bounds in (1) and (2).

**Conclusion:** A positive value of the first moment  $\Delta_s + \Delta_{\bar{s}}$  is almost impossible

What are the experimental results on  $(\Delta_s + \Delta_{\bar{s}})$  ?

BB02:  $-0.148 \pm 0.034$

LLS05:  $-0.132 \pm 0.018$

AAC06:  $-0.12 \pm 0.04$

deFNS06:  $-0.116$

HERMES05:  $+0.056 \pm 0.066 \pm 0.018$

There appears to be incompatibility. Is this real or is it due to underestimating the errors in the analyses ? It will be very interesting to have more accurate data from SIDIS.

What about the  $x$  – *dependence* of  $\Delta s(x) + \Delta \bar{s}(x)$  ?

Surprisingly, it turns out that **POSITIVITY** plays a crucial role

3) **Impose POSITIVITY**

$$|\Delta q(x)| \leq |q(x)|$$

**NB** Impose at **LOWEST**  $Q^2$  involved in evolution.



The shape of  $\Delta s(x) + \Delta \bar{s}(x)$  at medium  $x$  is tightly controlled by the unpolarized  $s(x)$ .

The function  $s(x)$  has changed in the many analyses of the unpolarized data, since the days of Glück, Reya, Vogelsang (GRV 1996).

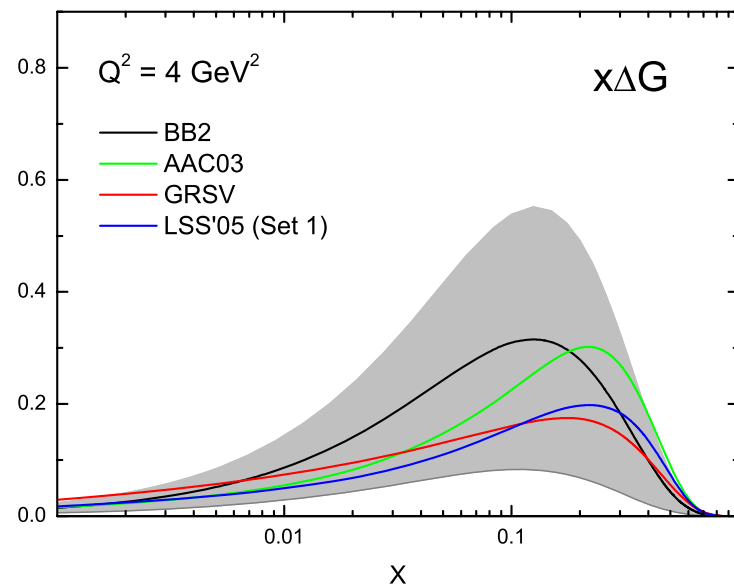
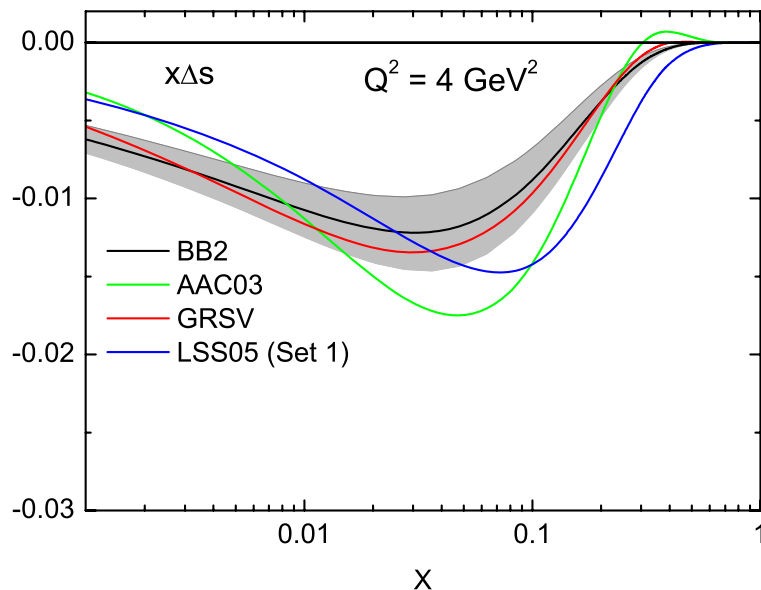
Fig.3 compares the differences between the  $\Delta s(x) + \Delta \bar{s}(x)$  from various analyses.

Fig.4 shows the role of positivity in creating these differences.

WHY DO PEOPLE CONTINUE TO USE THE GRV UNPOLARIZED DENSITIES, GIVEN THE MANY, MORE UP TO DATE ONES, AVAILABLE ??????????????????????

# NLO QCD PPD ( $\overline{\text{MS}}$ ) obtained by different groups

$x\Delta_S$  and  $x\Delta_G$  are **weakly** constrained from the present data on inclusive DIS



GRSV: Glück et al., hep-ph/0011215

BB: Blümlein, Böttcher, hep-ph/0203155

AAC03: Goto et al., hep-ph/0312112

LSS05: Leader et al., hep-ph/0512114

$x\Delta_{u_v}$  and  $x\Delta_{d_v}$  well consistent

# Impact of positivity constraints on $x\Delta s(x, Q^2)$

GRSV: Glück et al., hep-ph/0011215

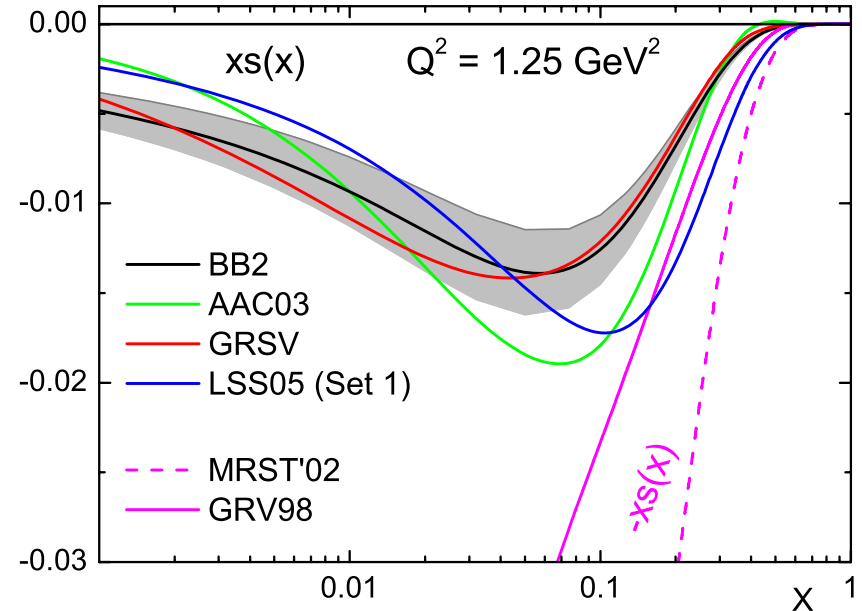
BB: Blümlein, Böttcher, hep-ph/0203155

AAC03: Goto et al., hep-ph/0312112

LSS05: Leader et al., hep-ph/0512114

$$|x\Delta f(x, Q_0^2)| \leq xf(x, Q_0^2)_{\text{GRV}}$$

$$|x\Delta f(x, Q_0^2)|_{\text{LSS}} \leq xf(x, Q_0^2)_{\text{MRST02}}$$



GRSV, BB and AAC have used the **GRV unpolarized** PD for constraining their PPD, while LSS have used those of **MRST'02**.

As a result,  $x|\Delta s(x)|$  (LSS) for  $x > 0.1$  is **larger** than the magnitude of the polarized strange sea densities obtained by the other groups.

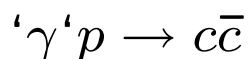
#### 4) Study $\Delta G$ via specific reactions

The range of  $Q^2$  in polarized DIS is too small to give a precise determination of  $\Delta G$ . Nonetheless essentially all analyses give positive  $\Delta G(x)$  with large error bands.

See Fig.3.

For a more precise determination we have to look at other possibilities.

The **gold plated** reaction is the photon fusion reaction

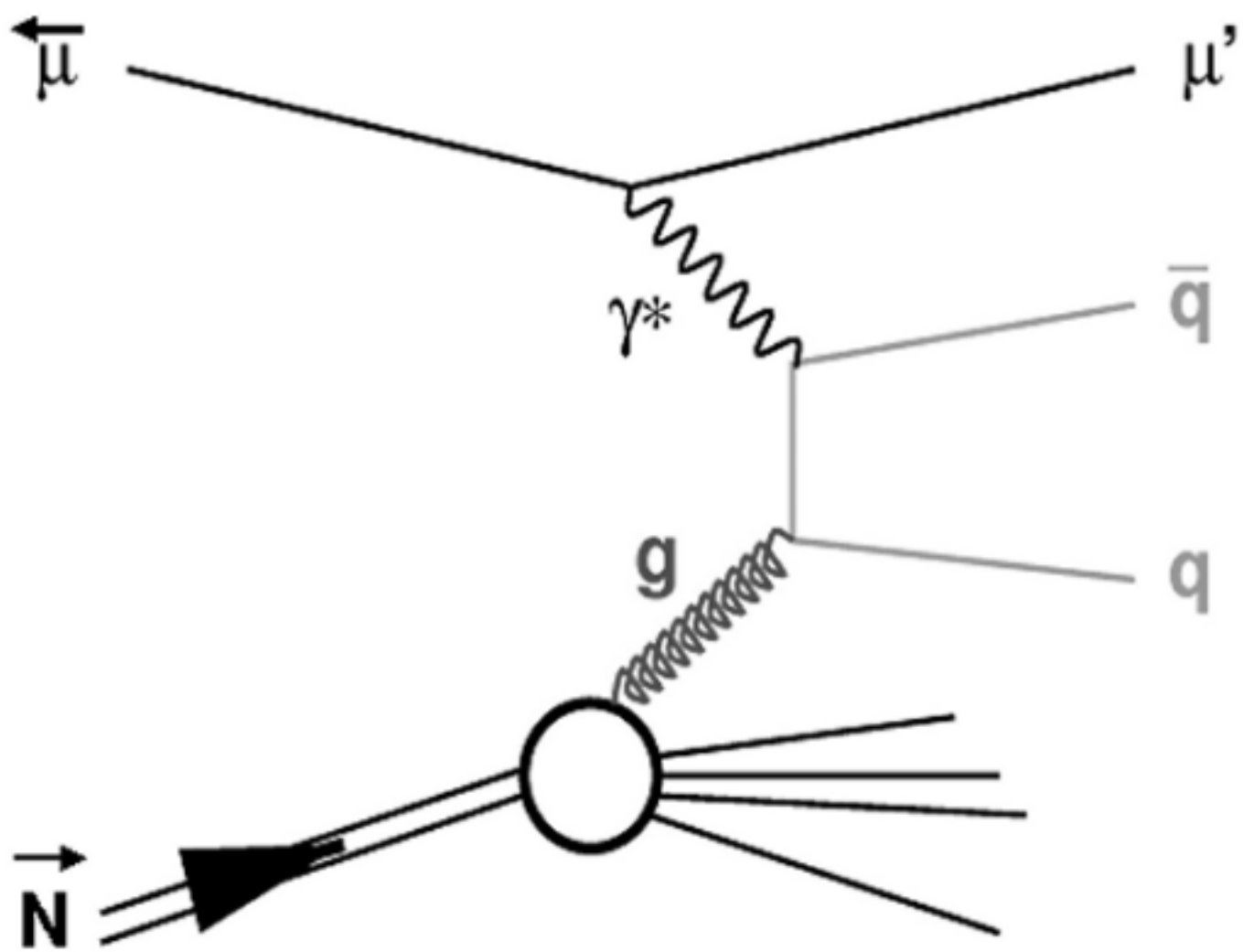


with identification of both charmed particles (open charm).

See Fig.5.

The next best (silver plated) is picking up one of the charmed particles.

Less clean is picking up two high  $P_T$  jets.



Old and new results are shown in Fig.6  
Errors are still large, but the situation looks  
intriguing.

Will there be a contradiction between  
HERMES and COMPASS?

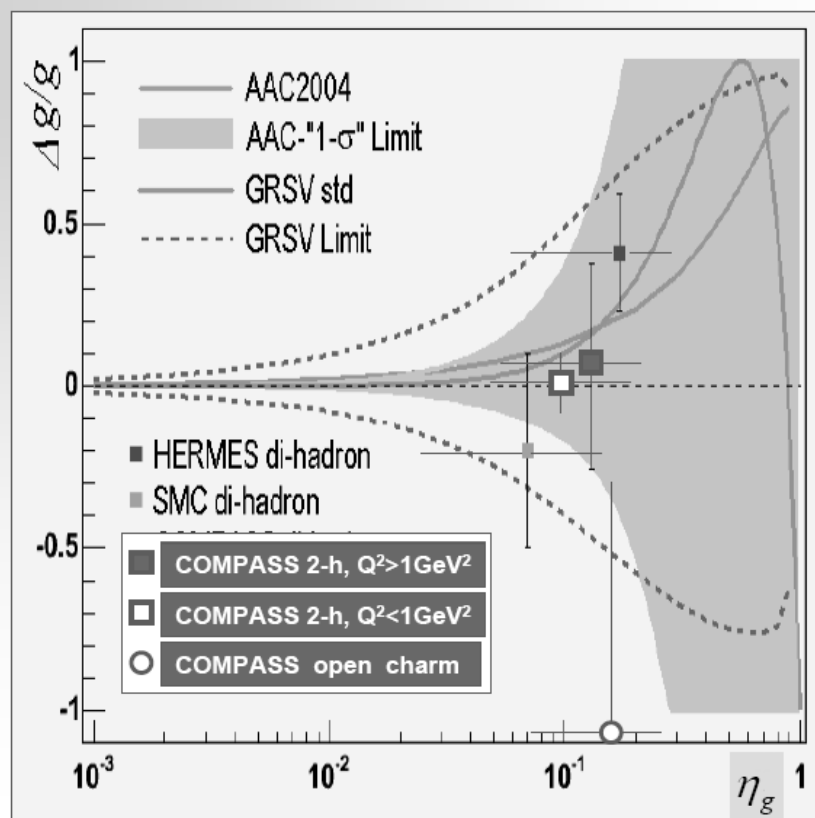
Is there a hint that  $\Delta G(x)$  changes sign?

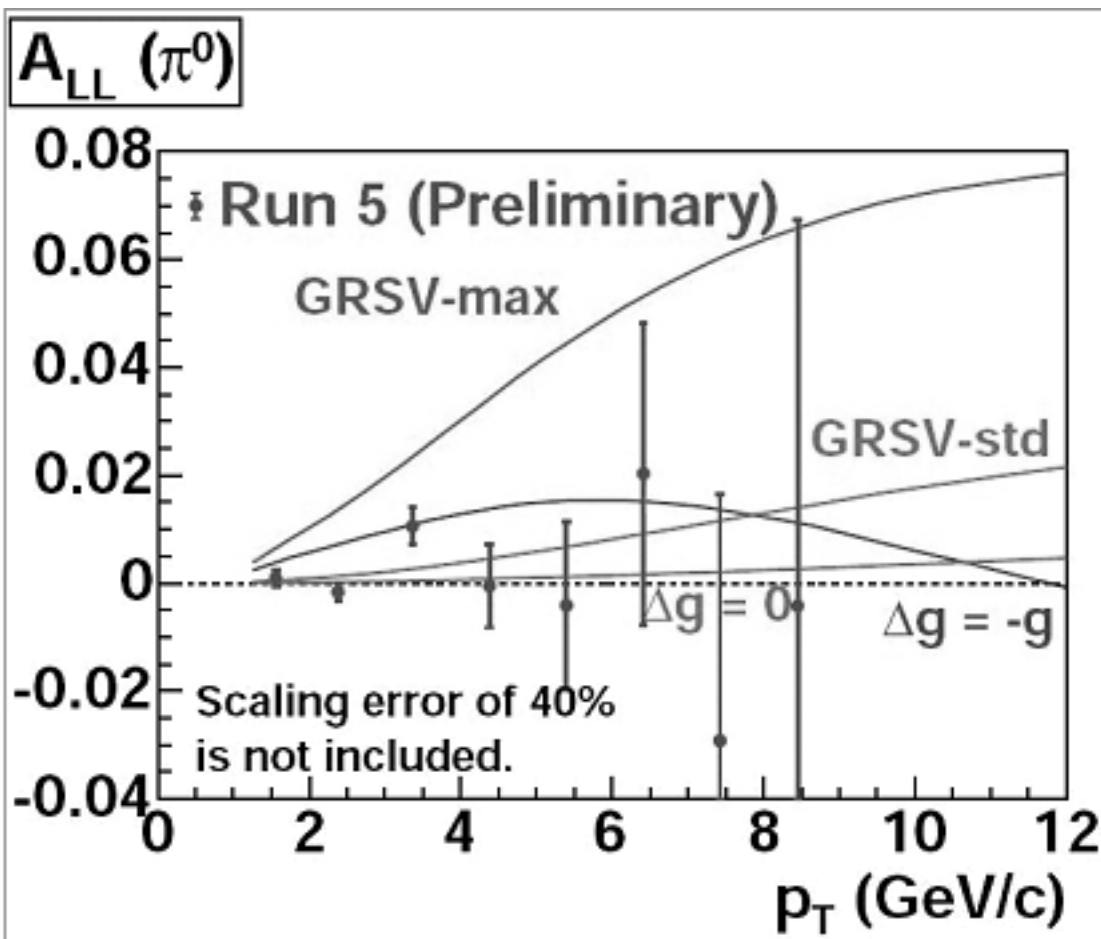
Another intriguing source of information:  $A_{LL}$   
in  $pp \rightarrow \pi X$  which is quadratic in  $\Delta G(x)$

Preliminary PHENIX data are shown in Fig.7.

Again errors are still large, but if  $A_{LL}$  is  
almost zero, it will require either that gluons  
are unpolarized **in contradiction with almost  
all the DIS results**, or that  $\Delta G(x)$  changes  
sign as a function of  $x$ .

## Results on $\Delta G$







An interesting development: AAC06 have fitted the DIS data with a  $\Delta G(x)$  which changes sign, and which **might** fit the PHENIX  $A_{LL}$  data.

If  $\Delta G(x) \neq 0$ , a good way to study its sign is via  $A_{LL}$  in  $pp \rightarrow \gamma X$ , which is linear in  $\Delta G(x)$ .

THE 'GOOD' NEWS: **IF**  $\Delta G(x)$  is very small we will be facing the RESURRECTION of **A CRISIS IN THE PARTON MODEL - WHERE, OH WHERE, IS THE PROTON'S SPIN?**

Recall that the small value of the proton's  $a_0$  was explained as a cancellation:-

$$a_0 = \Delta\Sigma - N_f(\alpha_s/2\pi) \Delta G$$

thereby allowing  $\Delta\Sigma$ , the spin carried by the quarks, to be reasonably large (say  $\approx 0.6$ ).

Higher Twist: necessity depends on WHAT  
you fit with QCD

$$A_1^{expt} = \frac{g_1^{QCD} + h/Q^2}{F_1^{QCD} + H/Q^2}$$

$$\approx (\text{empirically}) \frac{g_1^{QCD}}{F_1^{QCD}}$$

No HT needed. [Leader, Sidorov, Stamenov  
(LSS)]

$$g_1^{expt} = F_1^{expt} A_1^{expt}$$

$$= \frac{F_2^{expt}}{2x(1 + R^{expt})} A_1^{expt}$$

$$= g_1^{QCD} + h/Q^2$$

HT essential. [ Leader, Sidorov, Stamenov  
(LSS05)include HT]

[Blümlein, Böttcher (BB02) do not include  
HT]

$$\begin{aligned}g_1^{hybrid} &= \frac{F_2^{QCD}}{2x(1 + R^{expt})} A_1^{expt} \\ &= g_1^{QCD} + h'/Q^2\end{aligned}$$

HT necessary. [Hirai, Kumano, Saito (AAC 06) do not include HT]