

Recent developments in perturbative QCD

Gavin Salam

LPTHE, Universities of Paris VI and VII and CNRS

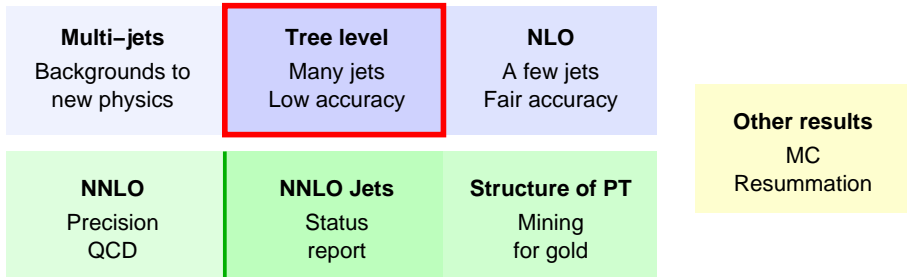
DIS 2006, Tsukuba, Japan,
20 April 2006

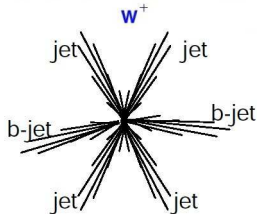
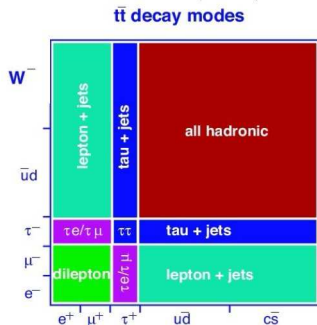
Concentrate on the effort to 'get QCD in shape for LHC era':

- ▶ Predicting multi-jet final-states:
[because new-particle signatures involve many jets]
 - ▶ New tree level techniques
 - ▶ NLO and progress in 1-loop calculations
- ▶ Aiming for accuracy
[because NLO theory is often far behind HERA/LEP precision]
 - ▶ NNLO jets: status & progress report
 - ▶ What NNLO is teaching us about QCD itself
- ▶ Other developments (mostly 'all-order' QCD)

Some recently very active fields, not covered:

- ▶ Small-x saturation ➡ talk by Iancu
- ▶ Generalised parton distributions – a field in its own right
➡ talk by Diehl, + hep-ph/0512201





All-hadronic

(BR~46%, huge bckg)

Juste – Lepton Photon '05

Heavy objects: multi-jet final-states

- ▶ Need to understand QCD multi-jet production (background)
- ▶ Max # jets: tree level ≤ 8 jets

MadEvent,AlpGen,Helac/Phegas

CompHEP,Grace,Amegic

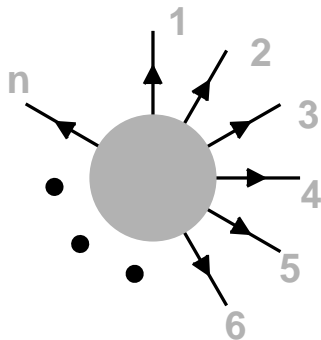
# jets	# events for 10 fb^{-1}
3	$9 \cdot 10^8$
4	$7 \cdot 10^7$
5	$6 \cdot 10^6$
6	$3 \cdot 10^5$
7	$2 \cdot 10^4$
8	$2 \cdot 10^3$

$p_t(\text{jet}) > 60 \text{ GeV}$, $\theta_{ij} > 30 \text{ deg}$, $|y_{ij}| < 3$

Draggiotis, Kleiss & Papadopoulos '02

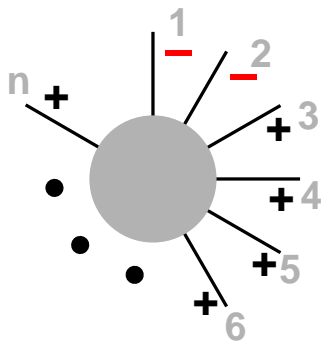
$$\mathcal{A}^{\text{tree}}(1, 2, \dots, n) = g^{n-2} \sum_{\text{perms}} \underbrace{\text{Tr}(T_1 T_2 \dots T_n)}_{\text{colour struct.}} \underbrace{\mathcal{A}^{\text{tree}}(1, 2, \dots, n)}_{\text{colour ordered amp.}}$$

$$\mathcal{A}^{\text{tree}}(1, 2, \dots, n) = g^{n-2} \sum_{\text{perms}} \underbrace{\text{Tr}(T_1 T_2 \dots T_n)}_{\text{colour struct.}} \underbrace{\mathcal{A}^{\text{tree}}(1, 2, \dots, n)}_{\text{colour ordered amp.}}$$



n	# diags	# col-ord diags
4	4	3
5	25	10
6	220	36
7	2485	133
8	34300	501
9	559405	1991
10	10525900	7335
	<i>non-planar</i>	<i>planar</i>

$$A^{\text{tree}}(1, 2, \dots, n) = g^{n-2} \sum_{\text{perms}} \underbrace{\text{Tr}(T_1 T_2 \dots T_n)}_{\text{colour struct.}} \underbrace{A^{\text{tree}}(1, 2, \dots, n)}_{\text{colour ordered amp.}}$$



Helicity amplitude: simplifies!

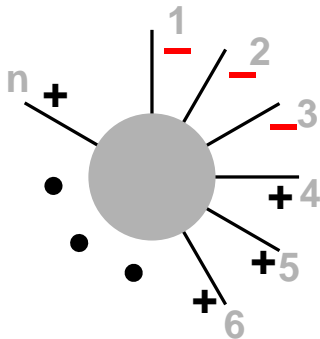
$$A^{\text{tree}}(- - + + \dots) = \frac{i \langle 12 \rangle^4}{\langle 12 \rangle \langle 23 \rangle \dots \langle n1 \rangle}$$

Parke & Taylor, Kunszt '85

Berends & Giele '88

Maximal Helicity Violating
(MHV)

$$A^{\text{tree}}(1, 2, \dots, n) = g^{n-2} \sum_{\text{perms}} \underbrace{\text{Tr}(T_1 T_2 \dots T_n)}_{\text{colour struct.}} \underbrace{A^{\text{tree}}(1, 2, \dots, n)}_{\text{colour ordered amp.}}$$



NEXT to Maximal Helicity
Violating (NMHV)

$$A_n(- + \dots + - -) =$$

$$= \frac{i}{\langle 12 \rangle \langle 23 \rangle \dots \langle (n-2)(n-1) \rangle \langle (n-2)(n-1) \rangle \langle (n-1)n \rangle \langle n1 \rangle [12]}$$

$$\times \left(\frac{\langle (n-1)n \rangle \langle 12 \rangle \langle (n-1)(n-2) \rangle \langle (n-1)^- | K_- | 2^- \rangle^2}{S_{3,n-1}} + \frac{\langle 1n \rangle \langle (n-1)(n-2) \rangle [12] \langle 1^- | K_- | (n-2)^- \rangle^2}{S_{1,n-3}} \right)$$

$$+ \frac{\langle n(n-1) \rangle \langle 1(n-1) \rangle \langle 1(n-2) \rangle [1n] [12] \langle (n-1)(n-2) \rangle \langle 1^- | K_- | (n-2)^- \rangle}{S_{1,n-3}}$$

$$+ \frac{\langle n1 \rangle \langle (n-1)1 \rangle \langle (n-1)2 \rangle [(n-1)n] [12] \langle (n-1)(n-2) \rangle \langle (n-1)^- | 2^- \rangle}{S_{3,n-1}}$$

$$- \langle 1(n-1) \rangle^2 S_{3,n-3} [12] \langle (n-1)(n-2) \rangle - \frac{\langle (n-1)n \rangle \langle 1n \rangle \langle 1(n-1) \rangle \langle 1^- | K_- | (n-2)^- \rangle [12]}{S_{1,n-3}}$$

$$- [n1] [n(n-1)] [12] \langle (n-1)(n-2) \rangle$$

$$\times \sum_{i=3}^{n-3} \left[\frac{\langle n(n-1) \rangle^2 \langle (n-1)1 \rangle \langle 1^- | K_{i,i-1} \tilde{\epsilon}_i | 1^+ \rangle}{S_{1,i-1} S_{1,i}} \right.$$

$$+ \frac{\langle n1 \rangle^2 \langle 1(n-1) \rangle \langle (n-1)^- | K_{i+1,n-1} \tilde{\epsilon}_i | (n-1)^+ \rangle}{S_{i+1,n-1} S_{i,n-1}}$$

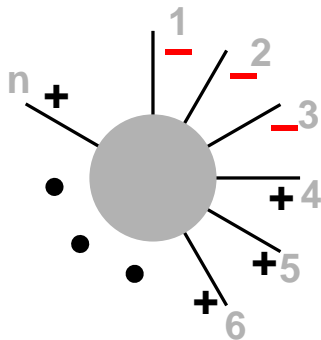
$$- \frac{\langle n1 \rangle \langle n(n-1) \rangle \langle (n-1)1 \rangle \langle (n-1)^- | K_{i+1,n} \tilde{\epsilon}_i | 1^+ \rangle}{S_{1,i} S_{i,n-1}}$$

$$- \frac{\langle n1 \rangle \langle n(n-1) \rangle^2 \langle (n-1)^- | K_{i+1,n} \tilde{\epsilon}_i | 1^+ \rangle \langle 1^- | K_{i,n} | n^- \rangle}{S_{1,i-1} S_{1,i} S_{i,n-1}}$$

$$\left. - \frac{\langle n1 \rangle^2 \langle n(n-1) \rangle \langle (n-1)^- | K_{i+1,n-1} \tilde{\epsilon}_i | 1^+ \rangle \langle (n-1)^- | K_{i+1,n} | n^- \rangle}{S_{1,i} S_{i+1,n-1} S_{i,n-1}} \right]$$

(5.2)

$$\mathcal{A}^{\text{tree}}(1, 2, \dots, n) = g^{n-2} \sum_{\text{perms}} \underbrace{\text{Tr}(T_1 T_2 \dots T_n)}_{\text{colour struct.}} \underbrace{\mathcal{A}^{\text{tree}}(1, 2, \dots, n)}_{\text{colour ordered amp.}}$$



NEXT to Maximal Helicity Violating (NMHV)

Helicity amplitude: simplifies!

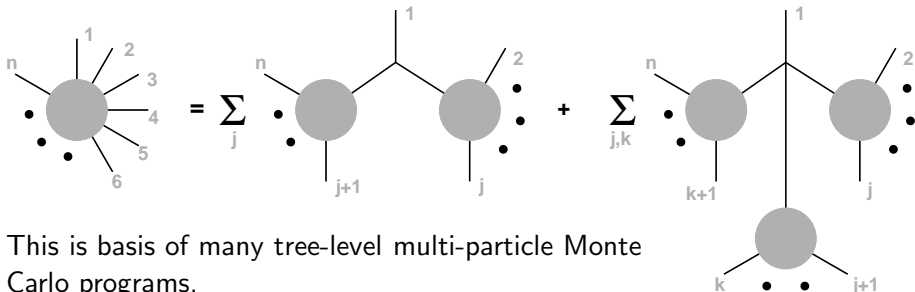
$$\begin{aligned}
 A^{\text{tree}}(- - - + + \dots) = & \\
 \frac{1}{F_{3,1}} \sum_{j=4}^{n-1} \frac{\langle 1 | P_{2,j} P_{j+1,2} | 3 \rangle}{P_{2,j}^2 P_{j+1,2}^2} \times & \\
 \frac{\langle j+1 j \rangle}{[2 | P_{2,j} | j+1 \rangle \langle j | P_{j+1,2} | 2 \rangle]} &
 \end{aligned}$$

Britto et al., hep-th/0503198

Just one of vast array of results obtained with new recursion (*Twistor*) techniques.

Build multi-leg amplitudes by joining sub-amplitudes.

Berends Giele (1988): Join smaller off-shell amplitudes through a (colour-stripped) three or four-gluon vertex:



This is basis of many tree-level multi-particle Monte Carlo programs.

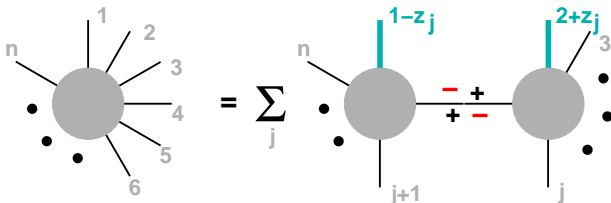
Why powerful?

Sub-amplitudes can be simplified *before* joining them together.

Feynman diagrams, in contrast, can only be simplified after full calculation.

New recursion relations (twistors)

Britto-Cachazo-Feng (BCF): Join smaller sub-amplitudes by a *propagator*.
 Sub-amplitudes made *on-shell* by analytic continuation ($\pm z_j$) of two reference momenta:



Britto, Cachazo & Feng hep-th/0412265; *idem.* + Witten hep-th/0501052
 Earlier (related) rules: Cachazo, Svrcek & Witten hep-th/0403047

Proof based on analytic structure of tree-graphs (they are a sum of poles in complex plane) — *very general*.

Simplicity lies in on-shellness of sub-amplitudes and the need for just a scalar propagator to join them.

Very active field: 200 articles in 2 years (~ 50 by 'QCD people')

Tree level

- ▶ Specific compact results, including NNMHV
 - ▶ Hints of yet deeper simplifications
- ▶ Efficient (recursive) formulations
 - ▶ NB: recall \exists 'standard' numerical methods for tree-level calculations:
- ▶ Massless quarks, gluinos
- ▶ External Higgs boson
- ▶ External weak boson (& fermions)
- ▶ Collinear limits
- ▶ Massive quarks, scalars

Kosower '04; Roiban et al '04
Luo & Wen '05; Britto et al '05

Bena, Bern, Kosower '04

Berends-Giele ('88); 'Alpha' ('95)

Georgiou, Glover & Khoze '04; Wu & Zhu '04

Dixon/Badger, Glover & Khoze '04

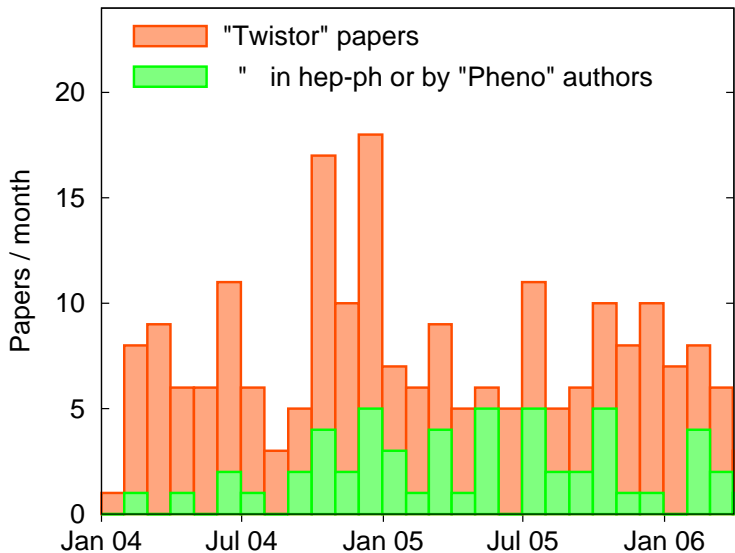
Bern, Forde, Kosower & Mastrolia '04

Birthwright et al '05

Forde & Kosower '05; Schwinn & Weinzierl '06

Ferrario, Rodrigo & Talavera '06; Ozeren & Stirling '06

Amazing progress in short time. . .



Multi-jets Backgrounds to new physics	Tree level Many jets Low accuracy	NLO A few jets Fair accuracy	Other results MC Resummation
NNLO Precision QCD	NNLO Jets Status report	Structure of PT Mining for gold	

Currently available

NLOJET++, MCFM, PHOX, ...

<http://www.cedar.ac.uk/hepcode/>Experimenters' priorities

1. $pp \rightarrow WW + \text{jet}$ Les Houches '05
2. $pp \rightarrow H + 2 \text{ jets}$
 - ▶ Background to VBF Higgs production
3. $pp \rightarrow t\bar{t}b\bar{b}$
4. $pp \rightarrow t\bar{t} + 2 \text{ jets}$
 - ▶ Background to $t\bar{t}H$
5. $pp \rightarrow WW b\bar{b}$
6. $pp \rightarrow VV + 2 \text{ jets}$
 - ▶ Background to $WW \rightarrow H \rightarrow WW$
7. $pp \rightarrow V + 3 \text{ jets}$
 - ▶ General background to new physics
8. $pp \rightarrow VVV + \text{jet}$
 - ▶ Background to SUSY trilepton

Currently available

NLOJET++, MCFM, PHOX, ...

<http://www.cedar.ac.uk/hepcode/>Experimenters' priorities

1. $pp \rightarrow WW + \text{jet}$ Les Houches '05
2. $pp \rightarrow H + 2 \text{ jets}$
 - ▶ Background to VBF Higgs production
3. $pp \rightarrow t\bar{t}b\bar{b}$
4. $pp \rightarrow t\bar{t} + 2 \text{ jets}$
 - ▶ Background to $t\bar{t}H$
5. $pp \rightarrow WW b\bar{b}$
6. $pp \rightarrow VV + 2 \text{ jets}$
 - ▶ Background to $WW \rightarrow H \rightarrow WW$
7. $pp \rightarrow V + 3 \text{ jets}$
 - ▶ General background to new physics
8. $pp \rightarrow VVV + \text{jet}$
 - ▶ Background to SUSY trilepton

Currently available

NLOJET++, MCFM, PHOX, ...

<http://www.cedar.ac.uk/hepcode/>Theorist's list (G. Heinrich)▶ $2 \rightarrow 3$ (OK for a good student!)

- ▶ $pp \rightarrow WW + \text{jet}$
- ▶ $pp \rightarrow VVV$
- ▶ $pp \rightarrow H + 2 \text{ jets}$

▶ $2 \rightarrow 4$ (Beyond today's means)

- ▶ $pp \rightarrow 4 \text{ jets}$
- ▶ $pp \rightarrow t\bar{t} + 2 \text{ jets}$
- ▶ $pp \rightarrow t\bar{t}b\bar{b}$
- ▶ $pp \rightarrow V + 3 \text{ jets}$
- ▶ $pp \rightarrow VV + 2 \text{ jets}$
- ▶ $pp \rightarrow VVV + \text{jet}$
- ▶ $pp \rightarrow WW b\bar{b}$

Experimenters' priorities

1. $pp \rightarrow WW + \text{jet}$ Les Houches '05
2. $pp \rightarrow H + 2 \text{ jets}$
 - ▶ Background to VBF Higgs production
3. $pp \rightarrow t\bar{t}b\bar{b}$
4. $pp \rightarrow t\bar{t} + 2 \text{ jets}$
 - ▶ Background to $t\bar{t}H$
5. $pp \rightarrow WW b\bar{b}$
6. $pp \rightarrow VV + 2 \text{ jets}$
 - ▶ Background to $WW \rightarrow H \rightarrow WW$
7. $pp \rightarrow V + 3 \text{ jets}$
 - ▶ General background to new physics
8. $pp \rightarrow VVV + \text{jet}$
 - ▶ Background to SUSY trilepton

$$2 \rightarrow 3 \text{ @ NLO} \sim \begin{array}{c} \text{Diagram 1: Tree-level } 2 \rightarrow 4 \end{array} + \begin{array}{c} \text{Diagram 2: 1-loop } 2 \rightarrow 3 \end{array} + \begin{array}{l} \text{Tricks to cancel} \\ \text{divergences} \\ \text{(dipole subtraction)} \end{array}$$

2 → 4 @ Tree
2 → 3 @ 1-loop

Traditionally: 1-loop for $2 \rightarrow 3$ proc. takes 1–2 years

Two ways of doing this more efficiently:

- ▶ Understand field theory better

Enormous progress on this in past two years: ~ 200 articles

- ▶ Get a computer to do most of the work for you

First full $2 \rightarrow 4$ (6-leg) result obtained this way

$$2 \rightarrow 3 \text{ @ NLO} \sim 2 \rightarrow 4 \text{ @ Tree} + \boxed{2 \rightarrow 3 \text{ @ 1-loop}} + \text{Tricks to cancel divergences (dipole subtraction)}$$

Bottleneck

Traditionally: 1-loop for $2 \rightarrow 3$ proc. takes 1–2 years

Two ways of doing this more efficiently:

- ▶ Understand field theory better

Enormous progress on this in past two years: ~ 200 articles

- ▶ Get a computer to do most of the work for you

First full $2 \rightarrow 4$ (6-leg) result obtained this way

$$2 \rightarrow 4 \text{ @ NLO} \sim 2 \rightarrow 5 \text{ @ Tree} + \boxed{2 \rightarrow 4 \text{ @ 1-loop}} + \text{Tricks to cancel divergences (dipole subtraction)}$$

Very hard!

The diagram shows the relationship between tree-level and one-loop calculations for a 2 to 4 process. On the left, a tree-level diagram for a 2 to 5 process is shown, consisting of a vertex with two incoming lines and three outgoing lines. This is added to a one-loop diagram for a 2 to 4 process, which is enclosed in a red box and labeled '2 to 4 @ 1-loop'. The one-loop diagram consists of a loop with two external lines and three other external lines. To the right of the box, the text 'Tricks to cancel divergences (dipole subtraction)' is written. Below the box, the text 'Very hard!' is written in red.

Traditionally: 1-loop for $2 \rightarrow 3$ proc. takes 1–2 years

Two ways of doing this more efficiently:

- ▶ Understand field theory better

Enormous progress on this in past two years: ~ 200 articles

- ▶ Get a computer to do most of the work for you

First full $2 \rightarrow 4$ (6-leg) result obtained this way

$$2 \rightarrow 4 \text{ @ NLO} \sim \begin{array}{c} \text{Diagram: 2 lines merging into 1, then splitting into 5 lines} \\ 2 \rightarrow 5 \text{ @ Tree} \end{array} + \begin{array}{c} \text{Diagram: 1-loop box diagram with 2 external lines and 4 internal lines} \\ 2 \rightarrow 4 \text{ @ 1-loop} \\ \text{Very hard!} \end{array} + \begin{array}{c} \text{Tricks to cancel} \\ \text{divergences} \\ \text{(dipole subtraction)} \end{array}$$

Traditionally: 1-loop for $2 \rightarrow 3$ proc. takes 1–2 years

Two ways of doing this more efficiently:

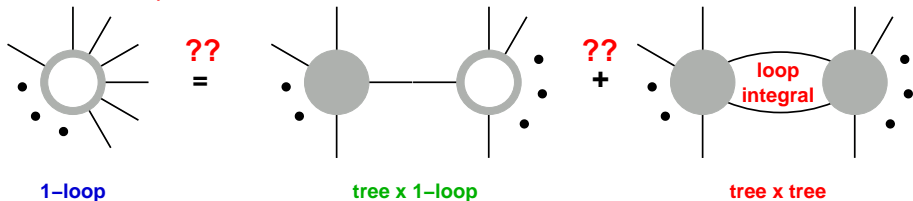
- ▶ Understand field theory better

Enormous progress on this in past two years: ~ 200 articles

- ▶ Get a computer to do most of the work for you

First full $2 \rightarrow 4$ (6-leg) result obtained this way

Would like a relation that avoids need for loop integrations. *Various kinds of recursion possible*

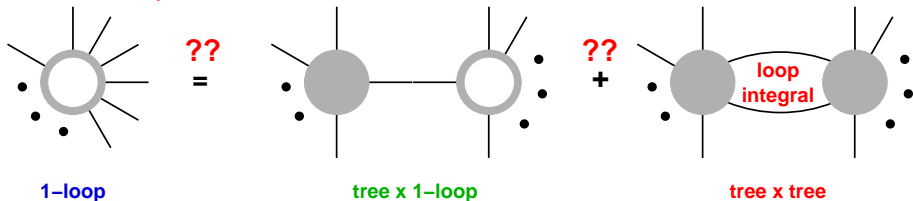


Technically: loop diagrams have more complex analytic properties than trees (*cuts* as well as poles), so BCFW does not apply.

Complex problem, much progress made, many people involved.

Bedford, Bena, Bern, Bidder, Bjerrum-Bohr, Brandhuber, Britto, Cachazo, Del Duca, Dixon, Dunbar, Feng, Forde, Ita, Kosower, McNamara, Mastrolia, Perkins, Roiban, Spence, Travaglini, [...]

Would like a relation that avoids need for loop integrations. *Various kinds of recursion possible*



Technically: loop diagrams have more complex analytic properties than trees (*cuts* as well as poles), so BCFW does not apply.

Complex problem, much progress made, many people involved.

Bedford, Bena, Bern, Bidder, Bjerrum-Bohr, Brandhuber, Britto, Cachazo, Del Duca, Dixon, Dunbar, Feng, Forde, Ita, Kosower, McNamara, Mastrolia, Perkins, Roiban, Spence, Travaglini, [...]

One ingredient of **one** of the “priority processes” ($pp \rightarrow 4$ jets) is the 6-gluon 1-loop amplitude:

$$\mathcal{A}_g = \underbrace{(\mathcal{A}_g + 4\mathcal{A}_f + 3\mathcal{A}_s)}_{\mathcal{N} = 4 \text{ SUSY}} - \underbrace{4(\mathcal{A}_f + \mathcal{A}_s)}_{\mathcal{N} = 1 \text{ chiral SUSY}} + \underbrace{\mathcal{A}_s}_{\text{scalar}}$$

	$\mathcal{N} = 4$	$\mathcal{N} = 1$	$S(c, d, e)$	$S(R)$
$A(- - + + +)$	BDDK94a	BDDK94b	BDDK94b	BDK05
$A(- + - + +)$	BDDK94a	BDDK94b	BBST04	
$A(- + + - +)$	BDDK94a	BDDK94b	BBST04	
$A(- - - + +)$	BDDK94b	BBDD04	BBDI05	Dixon05
$A(- - + - +)$	BDDK94b	BBCF05, BBDP04+5	BFM06	
$A(- + - + -)$	BDDK94b	BBCF05, BBDP04+5	BFM06	

Table adapted from hep-ph/0603187; NB: many results go beyond 6 gluons

Promising + much progress made! But QCD loops are still far from having simplicity of the tree-level results...

One ingredient of **one** of the “priority processes” ($pp \rightarrow 4$ jets) is the 6-gluon 1-loop amplitude:

$$\mathcal{A}_g = \underbrace{(\mathcal{A}_g + 4\mathcal{A}_f + 3\mathcal{A}_s)}_{\mathcal{N} = 4 \text{ SUSY}} - \underbrace{4(\mathcal{A}_f + \mathcal{A}_s)}_{\mathcal{N} = 1 \text{ chiral SUSY}} + \underbrace{\mathcal{A}_s}_{\text{scalar}}$$

	$\mathcal{N} = 4$	$\mathcal{N} = 1$	$S(c, d, e)$	$S(R)$
$A(- - + + +)$	BDDK94a	BDDK94b	BDDK94b	BDK05
$A(- + - + +)$	BDDK94a	BDDK94b	BBST04	
$A(- + + - +)$	BDDK94a	BDDK94b	BBST04	
$A(- - - + +)$	BDDK94b	BBDD04	BBDI05	Dixon05
$A(- - + - +)$	BDDK94b	BBCF05, BBDP04+5	BFM06	
$A(- + - + -)$	BDDK94b	BBCF05, BBDP04+5	BFM06	

Table adapted from hep-ph/0603187; NB: many results go beyond 6 gluons

Promising + much progress made! But QCD loops are still far from having simplicity of the tree-level results. . .

Automation of loop calculations with Feynman diagram techniques:

- ▶ Get expressions for all Feynman graphs (**QGRAF**, **FeynArts**). This gives answer in terms of a set of *loop integrals*
- ▶ Use *recursion relations* to reexpress each loop integral in terms of a basis set of known standard integrals NB: recursion for integrals, not amplitudes!
 - ▶ Analytically with algebraic manipulation programs
Binoth, Guillet, Heinrich, Pilon, Schubert '05; + others
 - ▶ Semi-numerically, “on the fly”, Ellis, Giele, Glover, Zanderighi '04-05
 - ▶ Results unstable at special phase-space points (e.g. co-planar momenta): use dedicated strategies there.
- ▶ Alternative integration techniques: e.g. subtract out divergences *before integrating*, do rest numerically.

Automation of loop calculations with Feynman diagram techniques:

- ▶ Get expressions for all Feynman graphs (**QGRAF**, **FeynArts**). This gives answer in terms of a set of *loop integrals*
- ▶ Use *recursion relations* to reexpress each loop integral in terms of a basis set of known standard integrals NB: recursion for integrals, not amplitudes!
 - ▶ Analytically with algebraic manipulation programs
Binoth, Guillet, Heinrich, Pilon, Schubert '05; + others
 - ▶ Semi-numerically, “on the fly”, Ellis, Giele, Glover, Zanderighi '04-05
 - ▶ Results unstable at special phase-space points (e.g. co-planar momenta): use dedicated strategies there.
- ▶ Alternative integration techniques: e.g. subtract out divergences *before integrating*, do rest numerically.

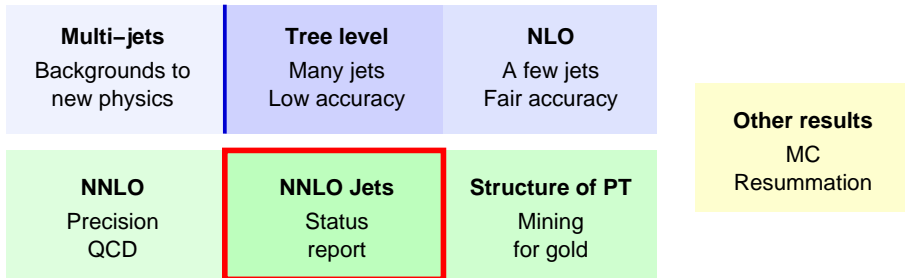
Automation of loop calculations with Feynman diagram techniques:

- ▶ Get expressions for all Feynman graphs (**QGRAF**, **FeynArts**). This gives answer in terms of a set of *loop integrals*
- ▶ Use *recursion relations* to reexpress each loop integral in terms of a basis set of known standard integrals NB: recursion for integrals, not amplitudes!
 - ▶ Analytically with algebraic manipulation programs
Binoth, Guillet, Heinrich, Pilon, Schubert '05; + others
 - ▶ Semi-numerically, “on the fly”, Ellis, Giele, Glover, Zanderighi '04-05
 - ▶ Results unstable at special phase-space points (e.g. co-planar momenta): use dedicated strategies there.
- ▶ Alternative integration techniques: e.g. subtract out divergences *before integrating*, do rest numerically.

- ▶ **Full 6-gluon 1-loop amplitude!** Ellis, Giele, Zanderighi '06
Only fully known 2 → 4 1-loop amplitude in QCD
- ▶ $pp \rightarrow H + 2 \text{ jets}$: amplitudes done, implementation into MCFM in progress Ellis, Campbell, Giele, Zanderighi, '05-06
- ▶ $gg \rightarrow WW$ via quark loop Binoth, Ciccolini, Kauer, Krämer '05
- ▶ Similar techniques in EW: $e^+e^- \rightarrow 4 \text{ fermions}$ Denner, Dittmaier, Roth, Wieders '05

Automated techniques have advantage of flexibility
 But: speed can be issue in numerical variants.

NB: more 'traditional' NLO methods still important, ➡ talk by Oleari



- ▶ Processes **with two QCD partons** @ LO are mostly done
 - ▶ $e^+e^- \rightarrow$ hadrons, $\tau \rightarrow \nu +$ hadrons
 - ▶ DIS coeff. fns., sum rules
 - ▶ $pp \rightarrow W, Z, \gamma^*, H, WH, ZH$ (many including spin correl.)

- ▶ **Next in line:** $e^+e^- \rightarrow 3$ jets?
 - ▶ simplest!
 - ▶ α_s & other measurements at LEP are theory limited
 theory uncertainty $\sim 3 - 4 \times$ exp. error
 - ▶ useful for studying perturbative/ non-perturbative interface.

- ▶ Then DIS $\rightarrow 2 + 1$ and $pp \rightarrow 2$ jets...

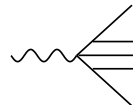
$$1 \rightarrow 3 \text{ @ NNLO} \sim \underbrace{\text{Diagram 1}}_{1 \rightarrow 5 \text{ @ Tree}} + \underbrace{\text{Diagram 2}}_{1 \rightarrow 4 \text{ @ 1-loop}} + \underbrace{\text{Diagram 3}}_{1 \rightarrow 3 \text{ @ 2-loop}} + \text{Tricks to cancel divergences}$$

$1 \rightarrow 5 \text{ @ Tree}$ $1 \rightarrow 4 \text{ @ 1-loop}$ $1 \rightarrow 3 \text{ @ 2-loop}$

$4 + 2\epsilon$ dim:
J is observable

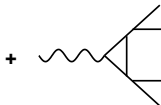
$1 \rightarrow 3$ @ NNLO

$$\int d\Phi_5 J(p_{1..5})$$



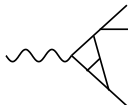
$1 \rightarrow 5$ @ Tree

$$\int d\Phi_4 \epsilon^{-2} J(p_{1..4})$$



$1 \rightarrow 4$ @ 1-loop

$$\int d\Phi_3 \epsilon^{-4} J(p_{1..3})$$



$1 \rightarrow 3$ @ 2-loop

+ Tricks to cancel divergences

$4 + 2\epsilon$ dim:
 J is observable

$1 \rightarrow 3$ @ NNLO

$$\int d\Phi_5 J(p_{1..5}) \quad \int d\Phi_4 \epsilon^{-2} J(p_{1..4}) \quad \int d\Phi_3 \epsilon^{-4} J(p_{1..3})$$

$1 \rightarrow 5$ @ Tree $1 \rightarrow 4$ @ 1-loop $1 \rightarrow 3$ @ 2-loop

Tricks to cancel divergences

Bottleneck

$4 + 2\epsilon$ dim:
J is observable

$1 \rightarrow 3$ @ NNLO

$$\int d\Phi_5 J(p_{1..5}) \quad \int d\Phi_4 \epsilon^{-2} J(p_{1..4}) \quad \int d\Phi_3 \epsilon^{-4} J(p_{1..3})$$

The diagram shows three Feynman diagrams representing different orders of perturbation theory. The first diagram is a tree-level process with one incoming wavy line and five outgoing straight lines. The second diagram is a one-loop process with one incoming wavy line and four outgoing straight lines, featuring a triangular loop. The third diagram is a two-loop process with one incoming wavy line and three outgoing straight lines, featuring a more complex loop structure. The diagrams are separated by plus signs.

$1 \rightarrow 5$ @ Tree $1 \rightarrow 4$ @ 1-loop $1 \rightarrow 3$ @ 2-loop

Tricks to cancel
divergences

Bottleneck

“You have to do the integral, but you don’t know the integrand”

Anastasiou (KITP LoopFest III)

Subtraction:

Catani, Seymour '96 + earlier authors

- find an integrable function with same divergences as amplitudes
- subtract it from real
- add integrated version to virtuals.

Sector decomposition:

Binoth, Heinrich '00

- split phase space into regions with at most one divergence each
- Introduce plus-prescription (i.e. as in splitting functions) to allow separate extraction of $\epsilon^{-4}, \dots, \epsilon^0$.

$4 + 2\epsilon$ dim:
 J is observable

$1 \rightarrow 3$ @ NNLO

$$\int d\Phi_5 J(p_{1..5}) + \int d\Phi_4 \epsilon^{-2} J(p_{1..4}) + \int d\Phi_3 \epsilon^{-4} J(p_{1..3})$$

$1 \rightarrow 5$ @ Tree $1 \rightarrow 4$ @ 1-loop $1 \rightarrow 3$ @ 2-loop

Tricks to cancel
 divergences

Bottleneck

“You have to do the integral, but you don't know the integrand”

Anastasiou (KITP LoopFest III)

► **Subtraction:**

Catani, Seymour '96 + earlier authors

- find an integrable function with same divergences as amplitudes
- subtract it from real
- add integrated version to virtuals.

► **Sector decomposition:**

Binoth, Heinrich '00

- split phase space into regions with at most one divergence each
- Introduce plus-prescription (i.e. as in splitting functions) to allow separate extraction of $\epsilon^{-4}, \dots, \epsilon^0$.

$4 + 2\epsilon$ dim:
 J is observable

$1 \rightarrow 3$ @ NNLO

$$\int d\Phi_5 J(p_{1..5}) + \int d\Phi_4 \epsilon^{-2} J(p_{1..4}) + \int d\Phi_3 \epsilon^{-4} J(p_{1..3})$$

$1 \rightarrow 5$ @ Tree $1 \rightarrow 4$ @ 1-loop $1 \rightarrow 3$ @ 2-loop

Tricks to cancel
 divergences

Bottleneck

“You have to do the integral, but you don't know the integrand”

Anastasiou (KITP LoopFest III)

► Subtraction:

Catani, Seymour '96 + earlier authors

- find an integrable function with same divergences as amplitudes
- subtract it from real
- add integrated version to virtuals.

► Sector decomposition:

Binoth, Heinrich '00

- split phase space into regions with at most one divergence each
- Introduce plus-prescription (i.e. as in splitting functions) to allow separate extraction of $\epsilon^{-4}, \dots, \epsilon^0$.

Subtraction

[Standard approach @ NLO]

- ▶ Applied to C_F^3 colour part of $e^+e^- \rightarrow 3$ jets

$$(\alpha_s C_F/2\pi)^3 \text{ piece of } \langle 1 - T \rangle = -20.4 \pm 4$$

Gehrmann-de Ridder, Gehrmann & Glover '04

- ▶ **New:** Full 'antenna' subtraction formulae recently published idem. '05

- ▶ talk by Del Duca for alternative subtractions

Sector decomposition

- ▶ Applied to $pp \rightarrow W, Z, H$ (fully differential, spin correlations)

Anastasiou, Dixon, Melnikov, Petriello '03–06

- ▶ **New:** partial $e^+e^- \rightarrow 3$ jets

Heinrich '06

Expect first full $e^+e^- \rightarrow 3$ jet results soon (end 2006)

Subtraction

[Standard approach @ NLO]

- ▶ Applied to C_F^3 colour part of $e^+e^- \rightarrow 3 \text{ jets}$

$$(\alpha_s C_F / 2\pi)^3 \text{ piece of } \langle 1 - T \rangle = -20.4 \pm 4$$

Gehrmann-de Ridder, Gehrmann & Glover '04

- ▶ **New:** Full 'antenna' subtraction formulae recently published idem. '05

↳ talk by Del Duca for alternative subtractions

Sector decomposition

- ▶ Applied to $pp \rightarrow W, Z, H$ (fully differential, spin correlations)

Anastasiou, Dixon, Melnikov, Petriello '03–06

- ▶ **New:** partial $e^+e^- \rightarrow 3 \text{ jets}$

Heinrich '06

Expect first full $e^+e^- \rightarrow 3 \text{ jet}$ results soon (end 2006)

Subtraction

[Standard approach @ NLO]

- ▶ Applied to C_F^3 colour part of $e^+e^- \rightarrow 3$ jets

$$(\alpha_s C_F/2\pi)^3 \text{ piece of } \langle 1 - T \rangle = -20.4 \pm 4$$

Gehrmann-de Ridder, Gehrmann & Glover '04

- ▶ **New:** Full 'antenna' subtraction formulae recently published idem. '05

↳ talk by Del Duca for alternative subtractions

Sector decomposition

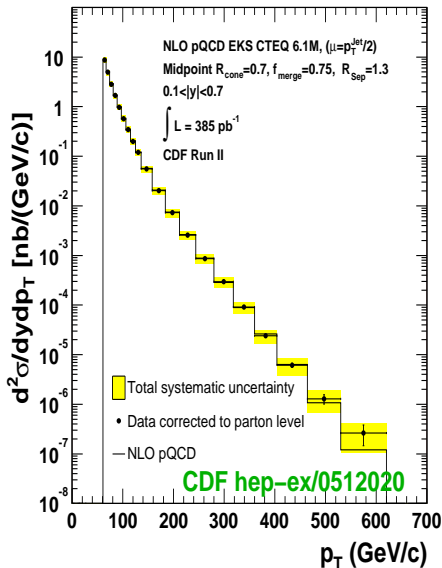
- ▶ Applied to $pp \rightarrow W, Z, H$ (fully differential, spin correlations)

Anastasiou, Dixon, Melnikov, Petriello '03–06

- ▶ **New:** partial $e^+e^- \rightarrow 3$ jets

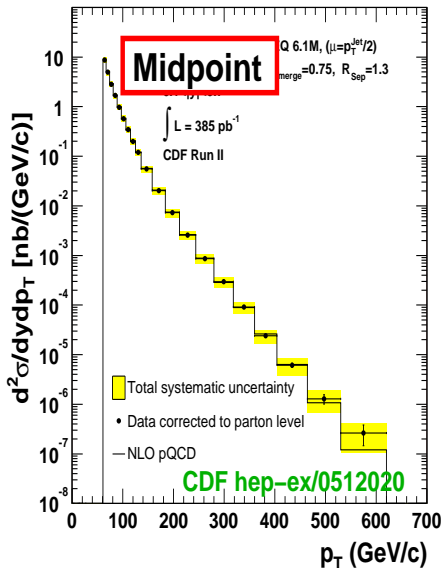
Heinrich '06

Expect first full $e^+e^- \rightarrow 3$ jet results soon (end 2006)



(N)NLO is **useless** if

- ▶ Jet-algo is not IR safe
 - CDF has *modified* midpoint cone
 - New 'search-cone step' IR unsafe [discovered by Wobisch]
 - ▶ Theory and experiment use different algorithms
 - R_{sep} in NLO theory, but not data
 - ▶ NB: 'NNLO-NLL' – rough approx. of NNLO, ignorant of jet-algo
- Good news:
- ▶ CDF also has k_T -algo result
 - ▶ Progress in making k_T -algo faster/friendlier
 - Cacciari [talk] & GPS '05–06

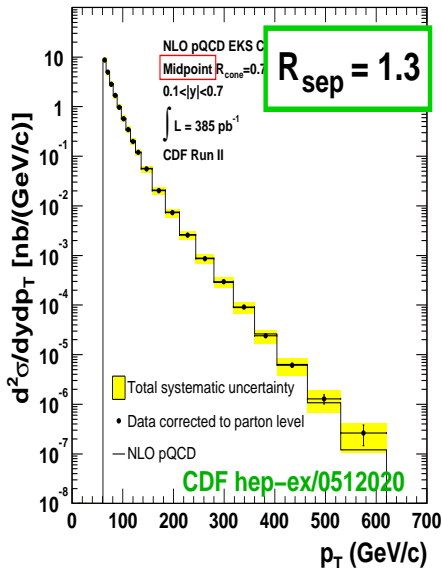


(N)NLO is **useless** if

- ▶ Jet-algo is not IR safe
 - CDF has *modified* midpoint cone
 - New 'search-cone step' IR unsafe
[discovered by Wobisch]
- ▶ Theory and experiment use different algorithms
 - R_{sep} in NLO theory, but not data
- ▶ NB: 'NNLO-NLL' – rough approx. of NNLO, ignorant of jet-algo

Good news:

- ▶ CDF also has k_t -algo result
- ▶ Progress in making k_t -algo faster/friendlier
Cacciari [talk] & GPS '05–06

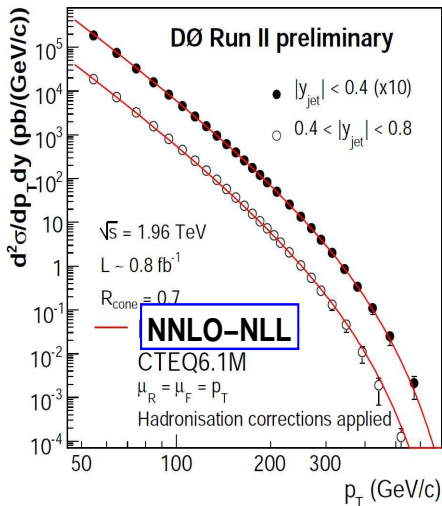


(N)NLO is **useless** if

- ▶ Jet-algo is not IR safe
 - CDF has *modified* midpoint cone
 - New 'search-cone step' IR unsafe [discovered by Wobisch]
- ▶ Theory and experiment use different algorithms
 - R_{sep} in NLO theory, but not data
- ▶ NB: 'NNLO-NLL' – rough approx. of NNLO, ignorant of jet-algo

Good news:

- ▶ CDF also has k_t -algo result
 - ▶ Progress in making k_t -algo faster/friendlier
- Cacciari [→talk] & GPS '05–06

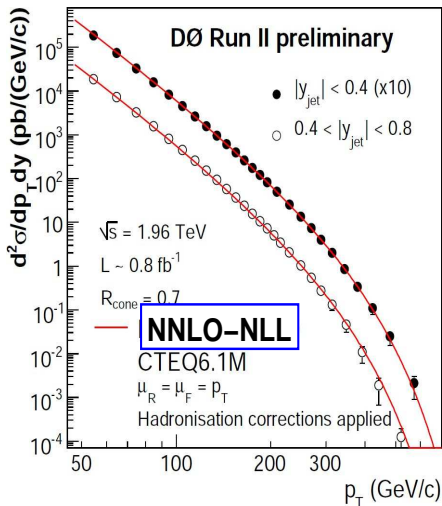


(N)NLO is **useless** if

- ▶ Jet-algo is not IR safe
 - CDF has *modified* midpoint cone
 - New 'search-cone step' IR unsafe [discovered by Wobisch]
- ▶ Theory and experiment use different algorithms
 - R_{sep} in NLO theory, but not data
- ▶ NB: 'NNLO-NLL' – rough apprx. of NNLO, ignorant of jet-algo

Good news:

- ▶ CDF also has k_t -algo result
 - ▶ Progress in making k_t -algo faster/friendlier
- Cacciari [→talk] & GPS '05–06



(N)NLO is **useless** if

- ▶ Jet-algo is not IR safe
 - CDF has *modified* midpoint cone
 - New 'search-cone step' IR unsafe [discovered by Wobisch]
- ▶ Theory and experiment use different algorithms
 - R_{sep} in NLO theory, but not data
- ▶ NB: 'NNLO-NLL' – rough appr. of NNLO, ignorant of jet-algo

Good news:

- ▶ CDF also has k_t -algo result
- ▶ Progress in making k_t -algo faster/friendlier
 - Cacciari [→talk] & GPS '05–06

Multi-jets Backgrounds to new physics	Tree level Many jets Low accuracy	NLO A few jets Fair accuracy	Other results MC Resummation
NNLO Precision QCD	NNLO Jets Status report	Structure of PT Mining for gold	

1. Next-to-Next-to-Leading Order Evolution of Non-Singlet Fragmentation Functions. By A. Mitov, S. Moch, A. Vogt. [hep-ph/0604053] DESY-06-036 (Apr 2006) 10p.
2. Higher-order soft corrections to lepton pair and Higgs boson production. By S. Moch & A. Vogt. Phys.Lett.B631:48-57,2005. [hep-ph/0508265]
3. Three-loop results for quark and gluon form-factors. By S. Moch, J.A.M. Vermaseren, A. Vogt. Phys.Lett.B625:245-252,2005. [hep-ph/0508055]
4. The Quark form-factor at higher orders. By S. Moch, J.A.M. Vermaseren, A. Vogt. JHEP 0508:049,2005. [hep-ph/0507039]
5. Higher-order corrections in threshold resummation. By S. Moch, J.A.M. Vermaseren, A. Vogt. Nucl.Phys.B726:317-335,2005. [hep-ph/0506288]
6. The Third-order QCD corrections to deep-inelastic scattering by photon exchange. By J.A.M. Vermaseren, A. Vogt, S. Moch. Nucl.Phys.B724:3-182,2005. [hep-ph/0504242]
7. The Longitudinal structure function at the third order. By S. Moch, J.A.M. Vermaseren, A. Vogt. Phys.Lett.B606:123-129,2005. [hep-ph/0411112]
8. Three loop universal anomalous dimension of the Wilson operators in N=4 SUSY Yang-Mills model. By A.V. Kotikov, L.N. Lipatov, A.I. Onishchenko, V.N. Velizhanin. Phys.Lett.B595:521-529,2004, Erratum-ibid.B632:754-756,2006. [hep-th/0404092]

Various unexpected structures in MVV results. E.g. at large x , can write

$$P_{ij}(x) = \frac{A}{(1-x)_+} + B \delta(1-x) + C \ln(1-x) + \mathcal{O}(1), \quad A = \sum A_n (\alpha_s/4\pi)^n, \text{ etc.}$$

Remarkably, different coefficients seem to be interrelated:

▶ $C_2 = A_1^2$

Curci, Furmanski, Petronzio '80

▶ $C_3 = 2A_1 A_2$

MVV '04

∃ a proposal that there is a more fundamental evolution equation with a *universal splitting function* appendix of Dokshitzer, Khoze, Troian '96

$$\partial_{\ln Q^2} D(x, Q^2) = \int_0^1 \frac{dz}{z} \mathcal{P}(x, \alpha_s(\frac{Q^2}{z})) D(\frac{x}{z}, z^\sigma Q^2) \quad \left\{ \begin{array}{l} \sigma = 1 : \text{ time-like} \\ \sigma = -1 : \text{ space-like} \end{array} \right.$$

Postulate new universal splitting function \mathcal{P} to be classical at large $x \Rightarrow$

$C = A^2$ at all orders; get most of NNLL $\mathcal{O}(1)$ term too!

Various unexpected structures in MVV results. E.g. at large x , can write

$$P_{ij}(x) = \frac{A}{(1-x)_+} + B \delta(1-x) + C \ln(1-x) + \mathcal{O}(1), \quad A = \sum A_n (\alpha_s/4\pi)^n, \text{ etc.}$$

Remarkably, different coefficients seem to be interrelated:

▶ $C_2 = A_1^2$

Curci, Furmanski, Petronzio '80

▶ $C_3 = 2A_1 A_2$

MVV '04

∃ a proposal that there is a more fundamental evolution equation with **a universal splitting function** appendix of Dokshitzer, Khoze, Troian '96

$$\partial_{\ln Q^2} D(x, Q^2) = \int_0^1 \frac{dz}{z} \mathcal{P}(x, \alpha_s(\frac{Q^2}{z})) D(\frac{x}{z}, z^\sigma Q^2) \quad \begin{cases} \sigma = 1 : & \text{time-like} \\ \sigma = -1 : & \text{space-like} \end{cases}$$

Postulate new universal splitting function \mathcal{P} to be classical at large $x \Rightarrow$

$C = A^2$ at all orders; get most of NNLL $\mathcal{O}(1)$ term too!

Dokshitzer, Marchesini, GPS '05

Original aim of Dokshitzer was to understand difference between time-like ($\sigma = +1$) and space-like ($\sigma = -1$) splitting functions.

i.e. fragmentation function and splitting function evolution

Normally related at order n via:

$$P_{\sigma=+1}^{(n)}(z) \iff P_{\sigma=-1}^{(n)}(1/z)$$

Curci, Furmanski, Petronzio '80

Stratmann & Vogelsang '97

New universality: get difference at order n from result at order $n - 1$

$$P_{\sigma=\pm 1}^{(n-1)}(z) \implies P_{\sigma=+1}^{(n)}(z) - P_{\sigma=-1}^{(n)}(z)$$

For non-singlet NNLO: both approaches give same prediction for time-like case

Original aim of Dokshitzer was to understand difference between time-like ($\sigma = +1$) and space-like ($\sigma = -1$) splitting functions.

i.e. fragmentation function and splitting function evolution

Normally related at order n via:

$$P_{\sigma=+1}^{(n)}(z) \iff P_{\sigma=-1}^{(n)}(1/z)$$

Curci, Furmanski, Petronzio '80

Stratmann & Vogelsang '97

New universality: get difference at order n from result at order $n - 1$

$$P_{\sigma=\pm 1}^{(n-1)}(z) \implies P_{\sigma=+1}^{(n)}(z) - P_{\sigma=-1}^{(n)}(z)$$

For non-singlet NNLO: both approaches give same prediction for time-like case

Original aim of Dokshitzer was to understand difference between time-like ($\sigma = +1$) and space-like ($\sigma = -1$) splitting functions.

i.e. fragmentation function and splitting function evolution

Normally related at order n via:

$$P_{\sigma=+1}^{(n)}(z) \iff P_{\sigma=-1}^{(n)}(1/z)$$

Curci, Furmanski, Petronzio '80

Stratmann & Vogelsang '97

New universality: get difference at order n from result at order $n - 1$

$$P_{\sigma=\pm 1}^{(n-1)}(z) \implies P_{\sigma=+1}^{(n)}(z) - P_{\sigma=-1}^{(n)}(z)$$

For non-singlet NNLO: both approaches give same prediction for time-like case

- ▶ Many other “goodies” in the MVV papers (even more in supersymmetric limit). . .
- ▶ In $\mathcal{N} = 4$ SUSY Yang-Mills amplitudes, planar n -loop seems to be reducible just to powers of 1-loop:

$$M_{n\text{-leg}}^{(2\text{-loop})} = \frac{1}{2}(M_n^{(1)})^2 + f(\epsilon)M_n^{(1)}(2\epsilon) - \frac{\pi^4}{72} + \mathcal{O}(\epsilon)$$

4-legs: Anastasiou, Bern, Dixon, Kosower '03

5-legs: Cachazo, Spradlin, Volovich '06; Bern et al '06

NB: numerical loop calcs: Anastasiou & Daleo [→talk] '05 ; Czakon '05

$$M_{n\text{-leg}}^{(3\text{-loop})} = -\frac{1}{3}(M_n^{(1)})^3 + M_n^{(1)}(\epsilon)M_n^{(2)}(\epsilon) + f^{(3)}(\epsilon)M_n^{(1)}(3\epsilon) + C^{(3)} + \mathcal{O}(\epsilon)$$

4-legs: Bern, Dixon, Smirnov '05

- ▶ In large-angle soft-colour resummation (‘fifth form factor’) for $2 \rightarrow 2$ scattering, symmetry in exch. of kinematic variables and $\#$ of colours:

$$\frac{\ln \frac{s^2}{ut} - 2\pi i}{\ln \frac{u}{t}} \iff N_C$$

Dokshitzer & Marchesini '05 (see also Seymour '05)

Multi-jets Backgrounds to new physics	Tree level Many jets Low accuracy	NLO A few jets Fair accuracy
NNLO Precision QCD	NNLO Jets Status report	Structure of PT Mining for gold

Other results
MC
Resummation

Fixed order calculations

- ▶ 4-loop decoupling relations for α_s (i.e. heavy-quark thresholds)
 - Schroder & Steinhauser '05; Chetyrkin, Kuhn, Sturm, '05
- ▶ IR safety for jet flavour
 - Banfi, GPS [→ talk], Zanderighi '06

MC calculations

- ▶ Herwig++ → adolescence ($pp \rightarrow DY$), Ariadne++
 - Lönnblad's talk
- ▶ Steady progress in matching MC & NLO
 - MC@NLO: → Frixione's talk
 - alternative methods: → Soper's, Nason's talks
- ▶ Using NNLL and NNLO for reweighting of event generators
 - Davatz et al '04; Davatz et al '06

Analytical resummations:

- ▶ Collinear region (and threshold): $MVV \Rightarrow \alpha_s^n L^{n-2}$
- ▶ Generic large angle region, even $\alpha_s^n L^n$ *much less well understood*
 - ▶ Gaps-between-jets phenomenology
 - Forshaw, Kyrieleis, Seymour '05-'06
 - ▶ Non-global: unanticipated new $\alpha_s^n L^n$ for jets
 - Banfi & Dasgupta [→ talk] '05

Fixed order calculations

- ▶ 4-loop decoupling relations for α_s (i.e. heavy-quark thresholds)
 - Schroder & Steinhauser '05; Chetyrkin, Kuhn, Sturm, '05
- ▶ IR safety for jet flavour
 - Banfi, GPS [→ talk], Zanderighi '06

MC calculations

- ▶ Herwig++ → adolescence ($pp \rightarrow DY$), Ariadne++
 - Lönnblad's talk
- ▶ Steady progress in matching MC & NLO
 - MC@NLO: → Frixione's talk
 - alternative methods: → Soper's, Nason's talks
- ▶ Using NNLL and NNLO for reweighting of event generators
 - Davatz et al '04; Davatz et al '06

Analytical resummations:

- ▶ Collinear region (and threshold): $MVV \Rightarrow \alpha_s^n L^{n-2}$
- ▶ Generic large angle region, even $\alpha_s^n L^n$ *much less well understood*
 - ▶ Gaps-between-jets phenomenology
 - Forshaw, Kyrieleis, Seymour '05-'06
 - ▶ Non-global: unanticipated new $\alpha_s^n L^n$ for jets
 - Banfi & Dasgupta [→ talk] '05

Fixed order calculations

- ▶ 4-loop decoupling relations for α_s (i.e. heavy-quark thresholds)
 - Schroder & Steinhauser '05; Chetyrkin, Kuhn, Sturm, '05
- ▶ IR safety for jet flavour
 - Banfi, GPS [→ talk], Zanderighi '06

MC calculations

- ▶ Herwig++ → adolescence ($pp \rightarrow DY$), Ariadne++
 - Lönnblad's talk
- ▶ Steady progress in matching MC & NLO
 - MC@NLO: → Frixione's talk
 - alternative methods: → Soper's, Nason's talks
- ▶ Using NNLL and NNLO for reweighting of event generators
 - Davatz et al '04; Davatz et al '06

Analytical resummations:

- ▶ Collinear region (and threshold): $MVV \Rightarrow \alpha_s^n L^{n-2}$
- ▶ Generic large angle region, even $\alpha_s^n L^n$ *much less well understood*
 - ▶ Gaps-between-jets phenomenology
 - Forshaw, Kyrieleis, Seymour '05-'06
 - ▶ Non-global: unanticipated new $\alpha_s^n L^n$ for jets
 - Banfi & Dasgupta [→ talk] '05

- ▶ Twistors / amplitude-recursion: major theory advance — starting to give very non-trivial results, especially for loops
 - Many string theorists now thinking about QCD
 - Some phenomenologists diverted into strings
- ▶ Automated 1-loop calculations are important complementary development.
 - More flexible; crucial for cross-checks
- ▶ $e^+e^- \rightarrow 3$ jets at NNLO is on final stretch
 - How much longer before DIS 2+1 and pp 2 \rightarrow 2?
- ▶ Once NNLO is available, comparison to data is not the only thing to be done with it.
 - Learn about structure in QCD
- ▶ Steady progress also for MC, resummations

Thanks to: Bern, Butterworth, P. Ciafaloni, Comelli, Dokshitzer, R.K. Ellis, Kosower, Lönnblad, Marchesini, Moretti, Seymour, Vogt, Webber

- ▶ Twistors / amplitude-recursion: major theory advance — starting to give very non-trivial results, especially for loops
 - Many string theorists now thinking about QCD
 - Some phenomenologists diverted into strings
- ▶ Automated 1-loop calculations are important complementary development.
 - More flexible; crucial for cross-checks
- ▶ $e^+e^- \rightarrow 3$ jets at NNLO is on final stretch
 - How much longer before DIS $2+1$ and $pp\ 2 \rightarrow 2$?
- ▶ Once NNLO is available, comparison to data is not the only thing to be done with it.
 - Learn about structure in QCD
- ▶ Steady progress also for MC, resummations

Thanks to: Bern, Butterworth, P. Ciafaloni, Comelli, Dokshitzer, R.K. Ellis, Kosower, Lönnblad, Marchesini, Moretti, Seymour, Vogt, Webber

EXTRA SLIDES

CPU time in seconds for the computation of the n gluon amplitude on a standard PC (2 GHz Pentium IV), summed over all helicities.

n	4	5	6	7	8	9	10	11	12
Berends-Giele	0.00005	0.00023	0.0009	0.003	0.011	<i>0.030</i>	<i>0.09</i>	<i>0.27</i>	<i>0.7</i>
CSW	<i>0.00001</i>	0.00040	0.0042	0.033	0.24	1.77	13	81	—
BCF	<i>0.00001</i>	<i>0.00007</i>	<i>0.0003</i>	<i>0.001</i>	<i>0.006</i>	0.037	0.19	0.97	5.5

Dinsdale, Ternick & Weinzierl '06

Gain a factor of ~ 4 for moderate n — useful, not overwhelming.

Slowly making it into phenomenological work

NB: trees in MadEvent, ALPGEN, HELAC/PHEGAS, CompHEP, GRACE, Amegic

↳ talk by Worek

But: real progress here is in discovery of new analytical structures in field theory (helpful also elsewhere, e.g. loops).

Supersymmetric decomposition (allow gluons, fermions and scalars in loops)

$$\mathcal{A}_g = \underbrace{(\mathcal{A}_g + 4\mathcal{A}_f + 3\mathcal{A}_s)}_{\mathcal{N} = 4 \text{ SUSY}} - \underbrace{4(\mathcal{A}_f + \mathcal{A}_s)}_{\mathcal{N} = 1 \text{ chiral SUSY}} + \underbrace{\mathcal{A}_s}_{\text{scalar}}$$

SUSY gives many cancellations. Most difficult piece is *scalar*.

Analytical structure involves coefficients (c, d, e) of standard boxes (l_4), triangles (l_3) and bubbles (l_2), and rational terms (R):

$$\mathcal{A}_s = \sum_i c_i l_4^i + \sum_i d_i l_3^i + \sum_i e_i l_2^i + R$$

- ▶ coefficients (c, d, e) can be (i) read off by merging trees (*cut constructibility*) (ii) obtained recursively (à la BCFW)
- ▶ rational parts can be obtained recursively

Example of Giele-Glover method

$$\int \frac{d^D \ell \ell^{\mu_1} \ell^{\mu_2}}{(\ell + q_1)^2 (\ell + q_2)^2 (\ell + q_3)^2 (\ell + q_4)^2}$$

$$= \frac{1}{2} g^{\mu_1 \mu_2} I(D+2; 1, 1, 1, 1) + 2q_1^{\mu_1} 2q_1^{\mu_2} I_4(D+4; 3, 1, 1, 1) + \dots$$

Then

$$2I_4(8; 3, 1, 1, 1) = -2 \left(\sum_i S_{1i}^{-1} \right) I_4(8; 2, 1, 1, 1)$$

$$- S_{11}^{-1} I_4(6; 1, 1, 1, 1) - S_{12}^{-1} I_3(6; 1, 0, 1, 1)$$

$$- S_{13}^{-1} I_3(6; 1, 1, 0, 1) - S_{14}^{-1} I_3(6; 1, 1, 1, 0)$$

The $I_n(D; 1, 1, 1, 1)$ etc. are the basis integrals. S_{ij} is kinematical matrix, $S_{ij} = (q_i - q_j)^2$.

Reduction procedure done numerically for each kinematic configuration.

4 + 2ε dim:

J is observable

1 → 3 @ NNLO

$$\int d\Phi_5 J(p_{1..5}) + \int d\Phi_4 \varepsilon^{-2} J(p_{1..4}) + \int d\Phi_3 \varepsilon^{-4} J(p_{1..3})$$

1 → 5 @ Tree 1 → 4 @ 1-loop 1 → 3 @ 2-loop

Tricks to cancel
divergences

Bottleneck

How to get cancellations?

1. Subtraction method:

$$\int d^D\Phi_5 M_5 J(p_{1..5}) + \int d^D\Phi_4 M_4 J(p_{1..4}) + \dots$$

Applied to $e^+e^- \rightarrow 2$ jets and C_F^3 colour part of $e^+e^- \rightarrow 3$ jets:

$$(\alpha_s C_F / 2\pi)^3 \text{ piece of } \langle 1 - T \rangle = -20.4 \pm 4$$

Gehrmann-De Ridder, Gehrmann & Glover '04

In principle all $e^+e^- \rightarrow 3$ jet 'antenna' subtraction pieces are ready —
'just' need to be coded!

idem. '05

4 + 2ε dim:

J is observable

1 → 3 @ NNLO

$$\int d\Phi_5 J(p_{1..5}) + \int d\Phi_4 \varepsilon^{-2} J(p_{1..4}) + \int d\Phi_3 \varepsilon^{-4} J(p_{1..3})$$

1 → 5 @ Tree 1 → 4 @ 1-loop 1 → 3 @ 2-loop

Tricks to cancel
divergences

Bottleneck

How to get cancellations?

1. Subtraction method:

$$\int d^4\Phi_5 [M_5 J(p_{1..5}) - S_5 J(\tilde{p}_{1..3})] + \int d^4\Phi_4 [M_4 J(p_{1..4}) + S_4 J(\tilde{p}_{1..3})] + \dots$$

Applied to $e^+e^- \rightarrow 2$ jets and C_F^3 colour part of $e^+e^- \rightarrow 3$ jets:

$$(\alpha_s C_F / 2\pi)^3 \text{ piece of } \langle 1 - T \rangle = -20.4 \pm 4$$

Gehrmann-De Ridder, Gehrmann & Glover '04

In principle all $e^+e^- \rightarrow 3$ jet 'antenna' subtraction pieces are ready —
'just' need to be coded!

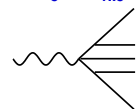
idem. '05

4 + 2ε dim:

J is observable

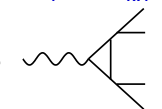
1 → 3 @ NNLO

$$\int d\Phi_5 J(p_{1..5})$$



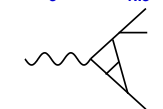
1 → 5 @ Tree

$$\int d\Phi_4 \varepsilon^{-2} J(p_{1..4})$$



1 → 4 @ 1-loop

$$\int d\Phi_3 \varepsilon^{-4} J(p_{1..3})$$



1 → 3 @ 2-loop

 Tricks to cancel
 divergences

Bottleneck

How to get cancellations?

1. Subtraction method:

$$\int d^4\Phi_5 [M_5 J(p_{1..5}) - S_5 J(\tilde{p}_{1..3})] + \int d^4\Phi_4 [M_4 J(p_{1..4}) + S_4 J(\tilde{p}_{1..3})] + \dots$$

 Applied to $e^+e^- \rightarrow 2$ jets and C_F^3 colour part of $e^+e^- \rightarrow 3$ jets:

$$(\alpha_s C_F / 2\pi)^3 \text{ piece of } \langle 1 - T \rangle = -20.4 \pm 4$$

Gehrmann-De Ridder, Gehrmann & Glover '04

In principle all $e^+e^- \rightarrow 3$ jet 'antenna' subtraction pieces are ready —
 'just' need to be coded!

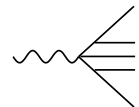
idem. '05

4 + 2ε dim:

J is observable

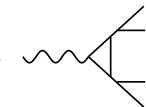
1 → 3 @ NNLO

$$\int d\Phi_5 J(p_{1..5})$$



1 → 5 @ Tree

$$\int d\Phi_4 \epsilon^{-2} J(p_{1..4})$$



1 → 4 @ 1-loop

$$\int d\Phi_3 \epsilon^{-4} J(p_{1..3})$$



1 → 3 @ 2-loop

+

 Tricks to cancel
 divergences

Bottleneck

How to get cancellations?

2. Sector decomposition for isolating divergences

Binoth & Heinrich '00

$$\int d^D\Phi_5 M_5 J(p_{1..5}) = \epsilon^{-4} \int d^4\Phi_5 f_{-4} M_5 J(p_{1..5}) + \dots + \int d^4\Phi_5 f_0 M_5 J(p_{1..5})$$

The f_{-i} involve plus-distributions of kinematic invariants. Each integral finite.

Applied to

▶ $e^+e^- \rightarrow 2 \text{ jets}$

Binoth & Heinrich '04; Anastasiou, Melnikov & Petriello '04

▶ $e^+e^- \rightarrow 3 \text{ jets (partial)}$

Heinrich '06

▶ $pp \rightarrow W, Z, H$ (fully exclusive)

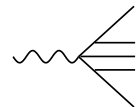
Anastasiou, Dixon, Melnikov & Petriello '04-06

4 + 2ε dim:

J is observable

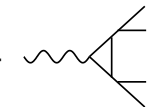
1 → 3 @ NNLO

$$\int d\Phi_5 J(p_{1..5})$$



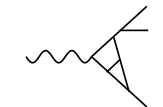
1 → 5 @ Tree

$$\int d\Phi_4 \epsilon^{-2} J(p_{1..4})$$



1 → 4 @ 1-loop

$$\int d\Phi_3 \epsilon^{-4} J(p_{1..3})$$



1 → 3 @ 2-loop

+

 Tricks to cancel
 divergences

Bottleneck

How to get cancellations?

2. Sector decomposition for isolating divergences

Binoth & Heinrich '00

$$\int d^D\Phi_5 M_5 J(p_{1..5}) = \epsilon^{-4} \int d^4\Phi_5 f_{-4} M_5 J(p_{1..5}) + \dots + \int d^4\Phi_5 f_0 M_5 J(p_{1..5})$$

The f_{-i} involve plus-distributions of kinematic invariants. Each integral finite.

Applied to

▶ $e^+e^- \rightarrow 2 \text{ jets}$

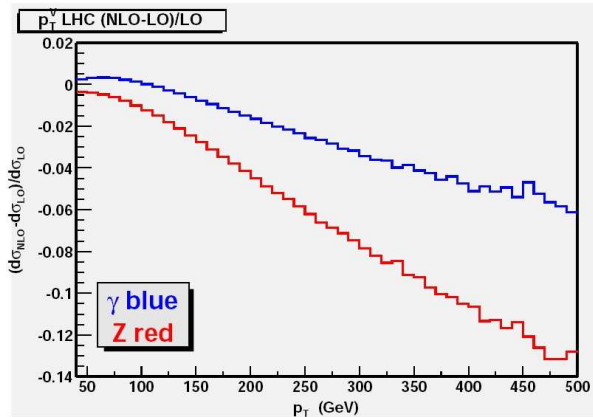
Binoth & Heinrich '04; Anastasiou, Melnikov & Petriello '04

▶ $e^+e^- \rightarrow 3 \text{ jets (partial)}$

Heinrich '06

▶ $pp \rightarrow W, Z, H$ (fully exclusive)

Anastasiou, Dixon, Melnikov & Petriello '04-06



Widely discussed for ILC. How about pp ?

e.g. NLO EW corrections to $pp \rightarrow Z + \text{jet}$

These are *significant* (even NNLO \sim few %)

Maina Moretti Ross '04

Kulesza et al '04

QED effects \lesssim 1%

Martin et al.

Glosser et al

Fortran

- ▶ Matching to multi-parton LO matrix elements now widespread CKKW
- ▶ New, better shower in Pythia (k_{\perp} ordered)
- ▶ Underlying event models much improved / more practical
- ▶ Reaching end of line soon! Pythia, Jimmy (Herwig)

C++

based on ThePEG			Independent	
Herwig++	Pythia 7	Ariadne/LDC	Pythia 8	Sherpa
ready for e^+e^- $pp \rightarrow DY$ ready	<i>cancelled</i>	see talk by Lönblad	being coded	ready for e^+e^- and pp

Includes new,
improved angular-
ordered shower

New player!
Dresden group

Resummation ingredients summary

order	Soft + collinear		Hard Collinear		Soft large angle		
	incl.	hadr.	incl.	hadr.	incl.	global	NG
$\alpha_s^n L^{n+1}$	✓	[✓]+BSZ04	N.A.			N.A.	
$\alpha_s^n L^n$	✓	[✓]+BSZ04	✓	✓	✓	[✓]	$[N_c \rightarrow \infty]$
$\alpha_s^n L^{n-1}$	MVV04	[FG04]	BCFG03	[FG04]	CFG+HK01	—	—
$\alpha_s^n L^{n-2}$	—	—	MVV05	—	MVV05	—	—

Large angle global | 2 → 2 BKOS89–98; generic: BCMN03; 2 → 3 [partial] KS05

Large angle NG | hemisph./patch: DS01–02; k_t algo: AS02, BD05

✓ ≡ historical results/techniques (< '01) [...] ≡ only for special cases
 hadr. ≡ anything measuring hadrons NG ≡ non-global

AS: Appleby & Seymour; BCFG: Bozzi & CFG; BCMN: Bonciani, Catani, Mangano, Nason; BD: Banfi & Dasgupta; BKOS: Botts, Kidonakis, Oderda, Sterman; BSZ: Banfi, GPS & Zanderighi; CFG: Catani & FG; DS: Dasgupta & GPS; FG: de Florian & Grazzini; KS: Kyrieleis & Seymour; HK: Harlander & Kilgore.

Best accuracies (NNLL) for most inclusive observables

Higgs transverse-mom. distribution.

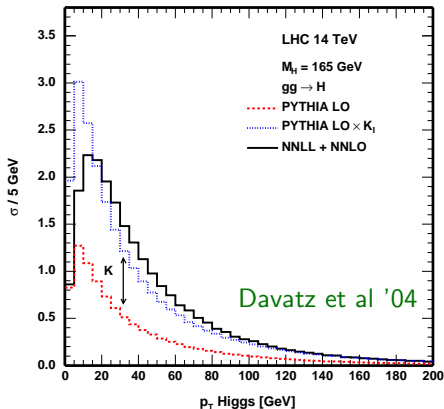
*First differential NNLL resummation*Resums $L \simeq \ln \frac{M_H}{Q_t}$ (for $gg \rightarrow H$)

$$\exp[\alpha_s^n L^{n+1} + \alpha_s^n L^n + \alpha_s^n L^{n-1}]$$

Bozzi et al '03

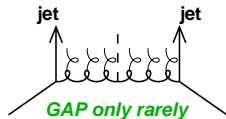
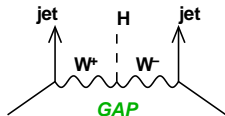
- ▶ NNLL uncertainty $\sim 7\%$
($\sim \text{NLL}/2$)
- ▶ Shape quite different from plain parton showering (Pythia)
— relevant for Higgs searches
($gg \rightarrow H \rightarrow WW \rightarrow l\nu l\nu$)?

Davatz et al '04



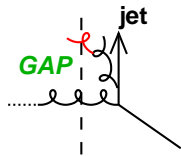
How rare are gaps in
 $pp \rightarrow 2 \text{ jets with big } \Delta Y$?

Answer needs advanced tools



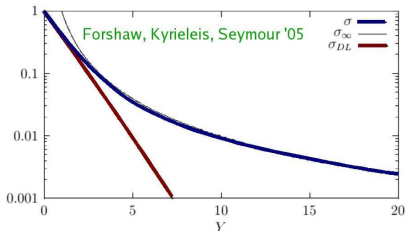
Non-global logarithms

- ▶ Appear for measurements of *part* of phase space
Also e.g. dijet properties, Banfi & Dasgupta '03
- ▶ Only in large- N_c limit! Not automated!
Connections to BFKL: Marchesini-Mueller '03; Weigert '03



Multi-jet structure

- ▶ Stony Brook soft-colour evolution
- ▶ Breakdown of 'probabilistic radiation'



Are Monte Carlos up to the job?

Unknown...