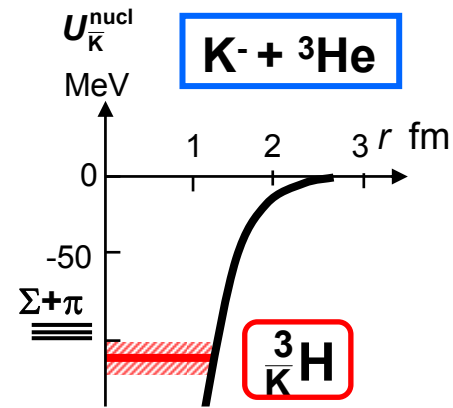
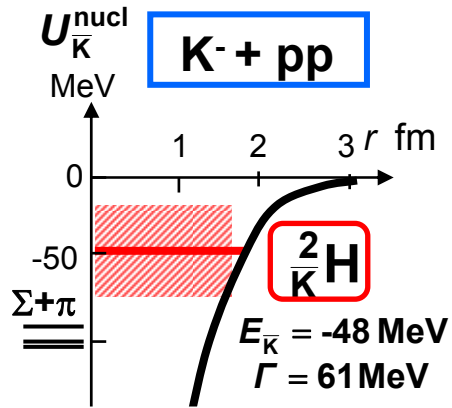
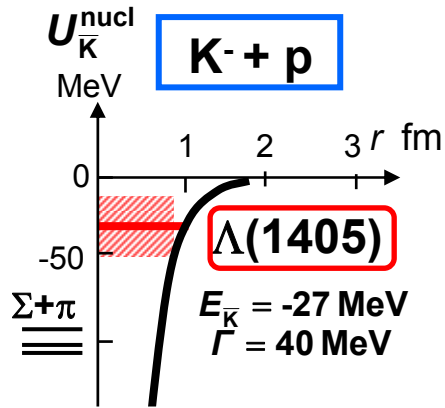


# Single-pole Nature of $\Lambda(1405)$ and Structure of $K^-pp$

Yoshinori AKAISHI & Toshimitsu YAMAZAKI

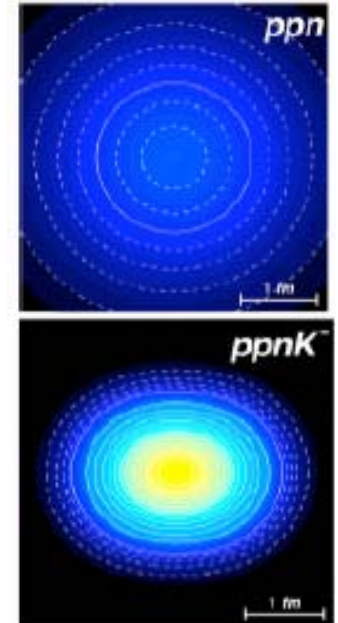
Mitsuaki OBU & Masanobu WADA

# " $\Lambda(1405)$ Ansatz"

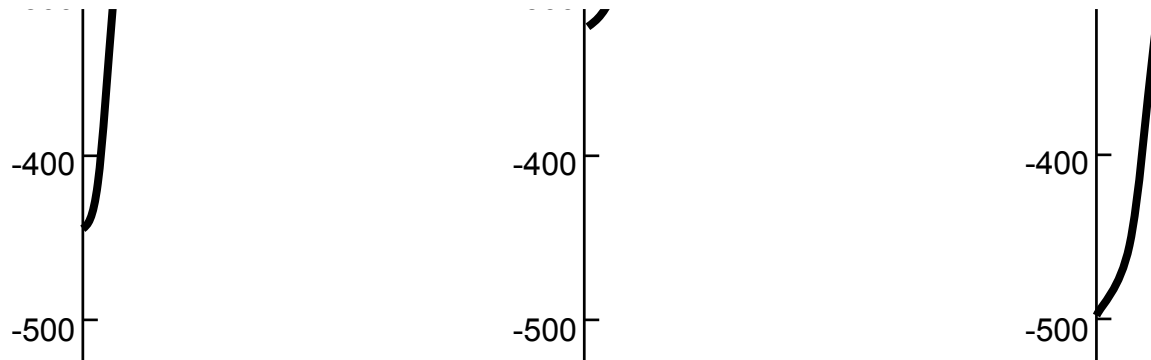


N.V. Shevchenko, A. Gal & J. Mares, Phys. Rev. Lett. 98 (2007) 082301  
 $E = -55 \sim -70 \text{ MeV}$ ,  $\Gamma = 90 \sim 110 \text{ MeV}$   
 Y. Ikeda & T. Sato, Phys. Rev. C 76 (2007) 035203  
 $E = -80 \text{ MeV}$ ,  $\Gamma = 73 \text{ MeV}$

A. Dote et al.



Shrinkage!



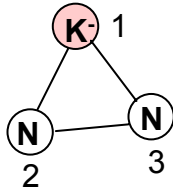
T. Yamazaki & Y. Akaishi, Phys. Lett. B 535 (2002) 70  
 Y. Akaishi & T. Yamazaki, Phys. Rev. C 65 (2002) 044005

# Variational wave function of K<sup>-</sup>pp

## ATMS

Amalgamation of **T**wo-body correlations into **M**ultiple **S**cattering process

$$\Psi = \left[ \left\{ f^{l=0}(r_{12}) \hat{P}_{12}^{l=0} + f^{l=1}(r_{12}) \hat{P}_{12}^{l=1} \right\} f_{NN}(r_{23}) f(r_{31}) + f(r_{12}) f_{NN}(r_{23}) \left\{ f^{l=0}(r_{31}) \hat{P}_{31}^{l=0} + f^{l=1}(r_{31}) \hat{P}_{31}^{l=1} \right\} \right] |T = 1/2\rangle$$



$$\hat{P}_{12}^{l=0} = \frac{1 - \vec{r}_K \vec{r}_N}{4}, \quad \hat{P}_{12}^{l=1} = \frac{3 + \vec{r}_K \vec{r}_N}{4}$$

$$|T = 1/2\rangle = \sqrt{\frac{3}{4}} \left[ \left[ (\bar{K}_1 N_2)^{0,0} p_3 \right] \right] + \sqrt{\frac{1}{4}} \left[ \left[ -\sqrt{\frac{1}{3}} (\bar{K}_1 N_2)^{1,0} p_3 + \sqrt{\frac{2}{3}} (\bar{K}_1 N_2)^{1,1} n_3 \right] \right]$$

$\Lambda^* p$

## Euler-Lagrange equation

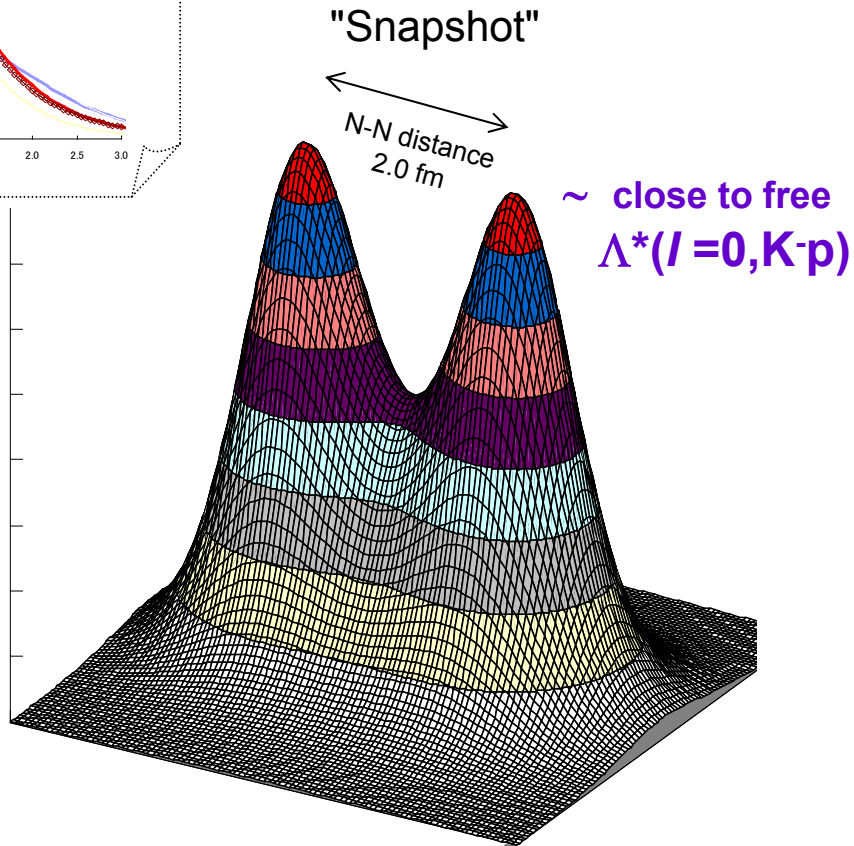
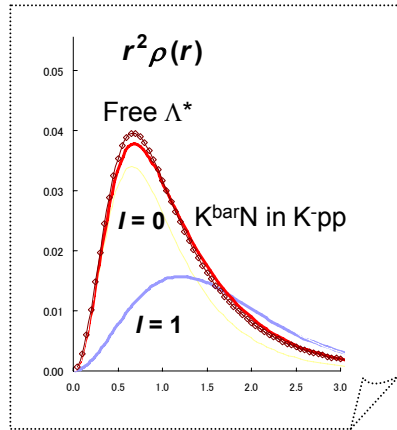
$$\delta_f \{ \langle \Psi | H | \Psi \rangle - \lambda \langle \Psi | \Psi \rangle \} = 0$$

$$v_{KN}^{T=0}(r) = \{ -595 - i83 \}_{\text{MeV}} \exp\left\{ - (r/0.66_{\text{fm}})^2 \right\}$$

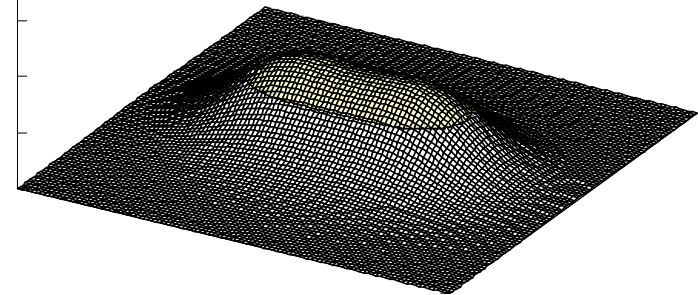
$$v_{KN}^{T=1}(r) = \{ -175 - i105 \}_{\text{MeV}} \exp\left\{ - (r/0.66_{\text{fm}})^2 \right\}$$

$$v_{NN}(r) = 2000_{\text{MeV}} \exp\left\{ - (r/0.447_{\text{fm}})^2 \right\} - 270_{\text{MeV}} \exp\left\{ - (r/0.942_{\text{fm}})^2 \right\} - 5_{\text{MeV}} \exp\left\{ - (r/2.5_{\text{fm}})^2 \right\}$$

# K- distribution in K $\bar{p}$ p



Covalent part

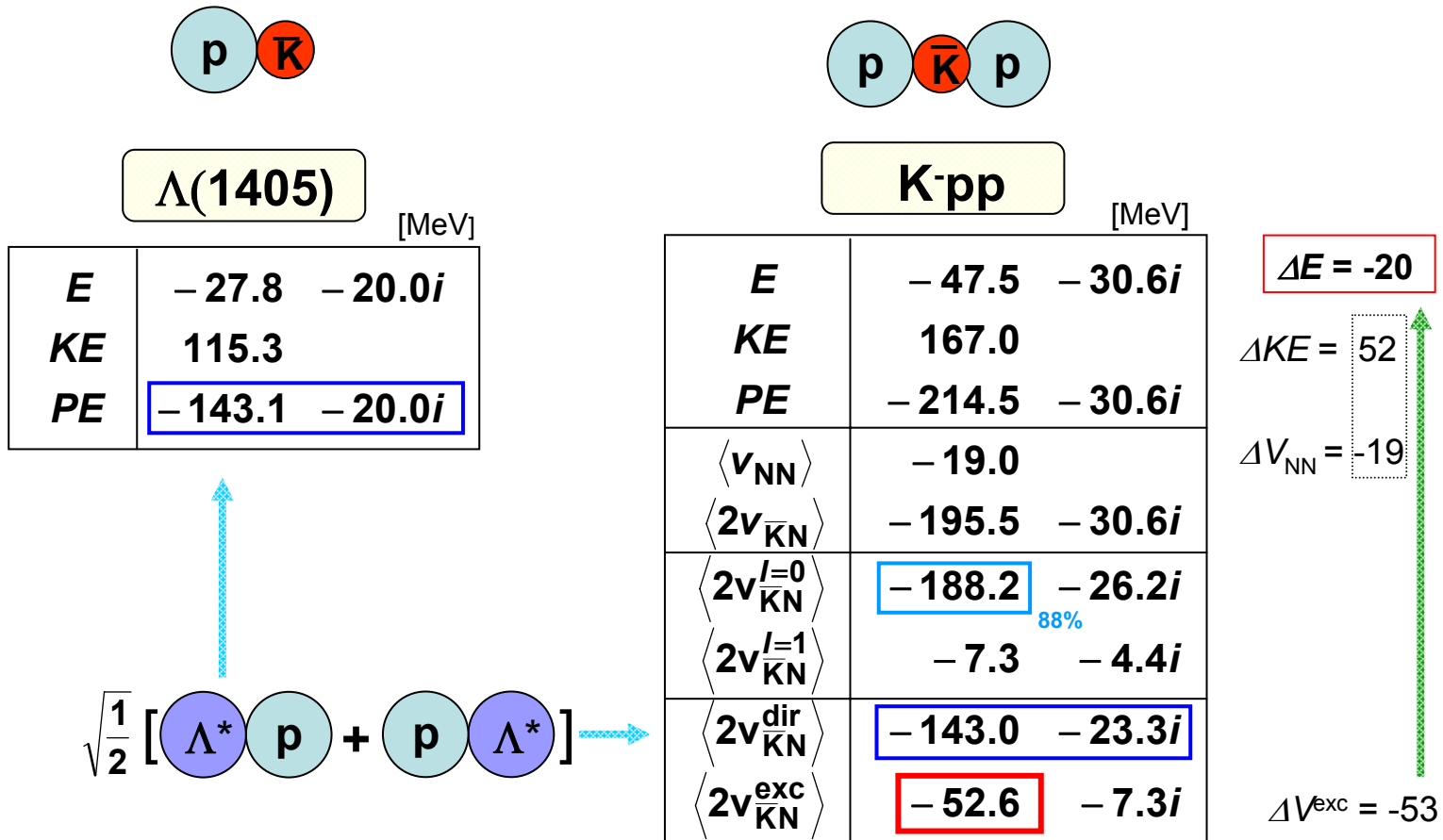


$$\Psi_{\pm} = \phi_a \pm \phi_b$$

$$\int d\vec{r} |\Psi_{\pm}|^2 = \int d\vec{r}_a |\phi_a(\vec{r}_a)|^2 + \int d\vec{r}_b |\phi_b(\vec{r}_b)|^2 \pm \int d\vec{r} \left[ \phi_a^*(\vec{r}_a) \phi_b(\vec{r}_b) + \phi_b^*(\vec{r}_b) \phi_a(\vec{r}_a) \right]$$

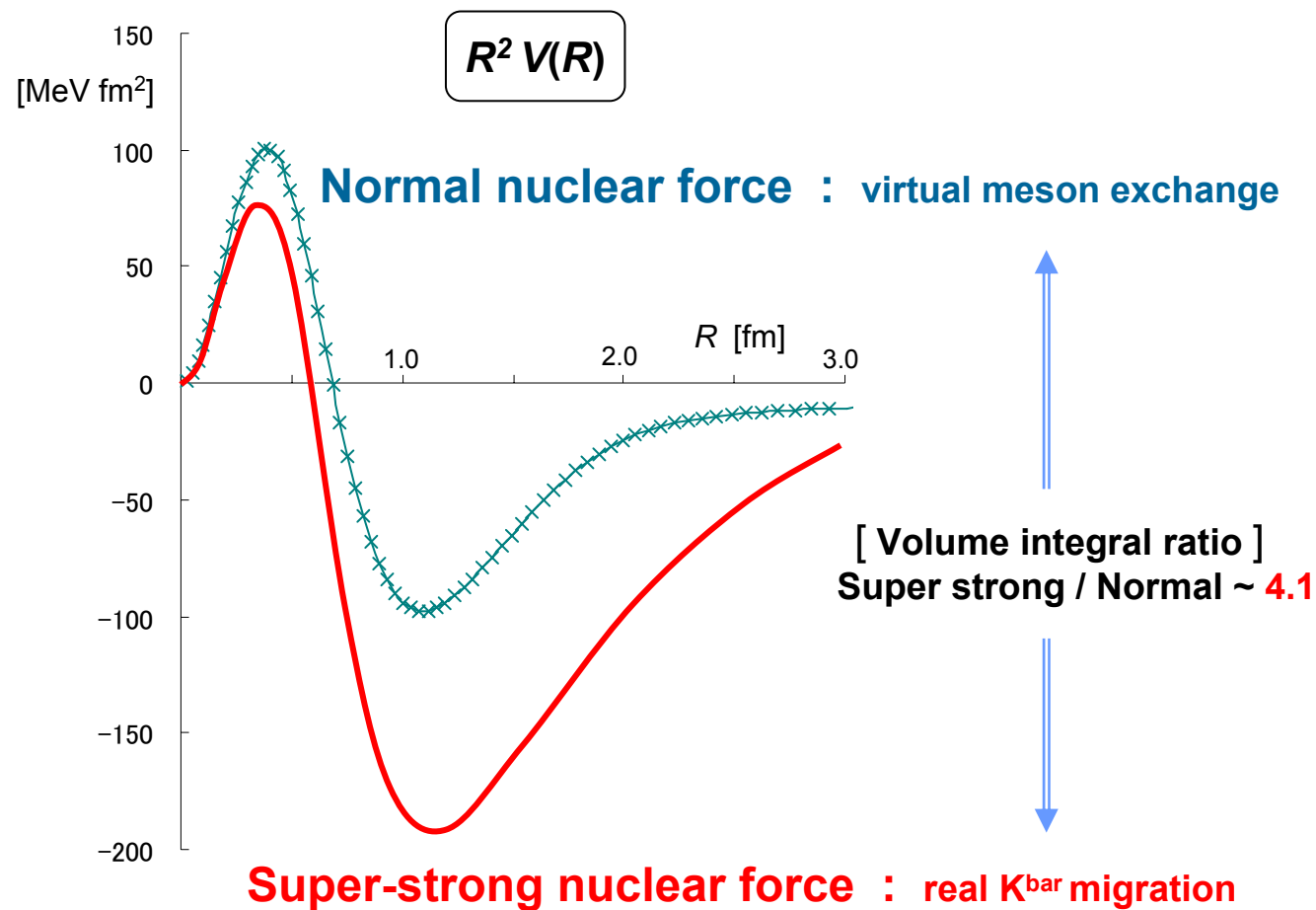
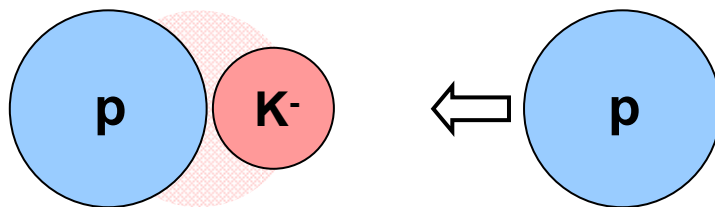
Covalent bonding

# Heitler-London picture of $K^-pp$

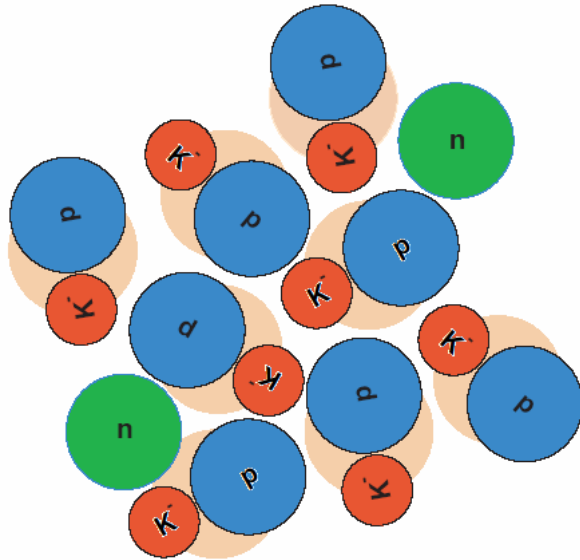


Real  $K^{\text{bar}}$  migrating attraction

# Adiabatic p-p potential in $K^-pp$



$\Lambda^* = (K-p)_{I=0}$  condensed matter



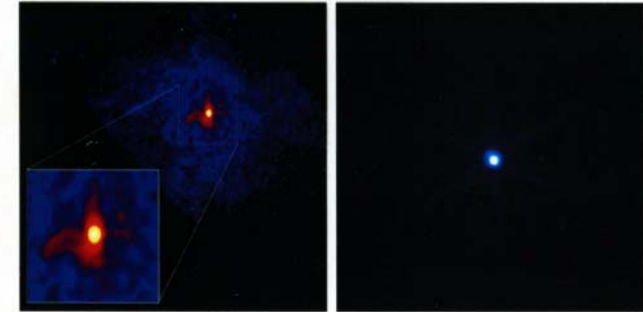
**Real kaons are migrating!**  
— a new paradigm —

**New!**

DISTO data on K<sup>-</sup>pp

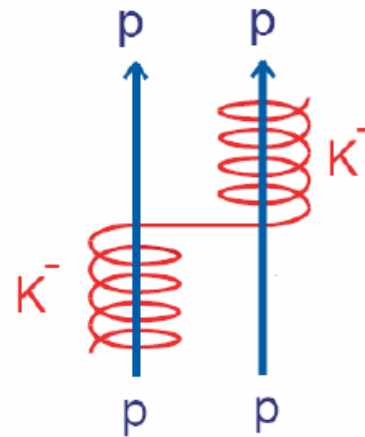
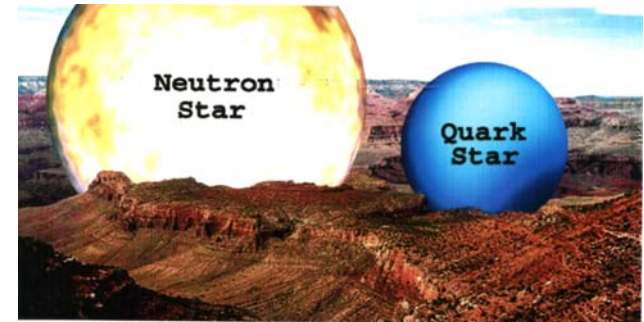
M. Maggiora et al.

NASA's Chandra X-ray

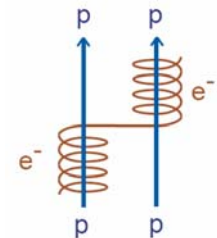


3C59

RX J1856



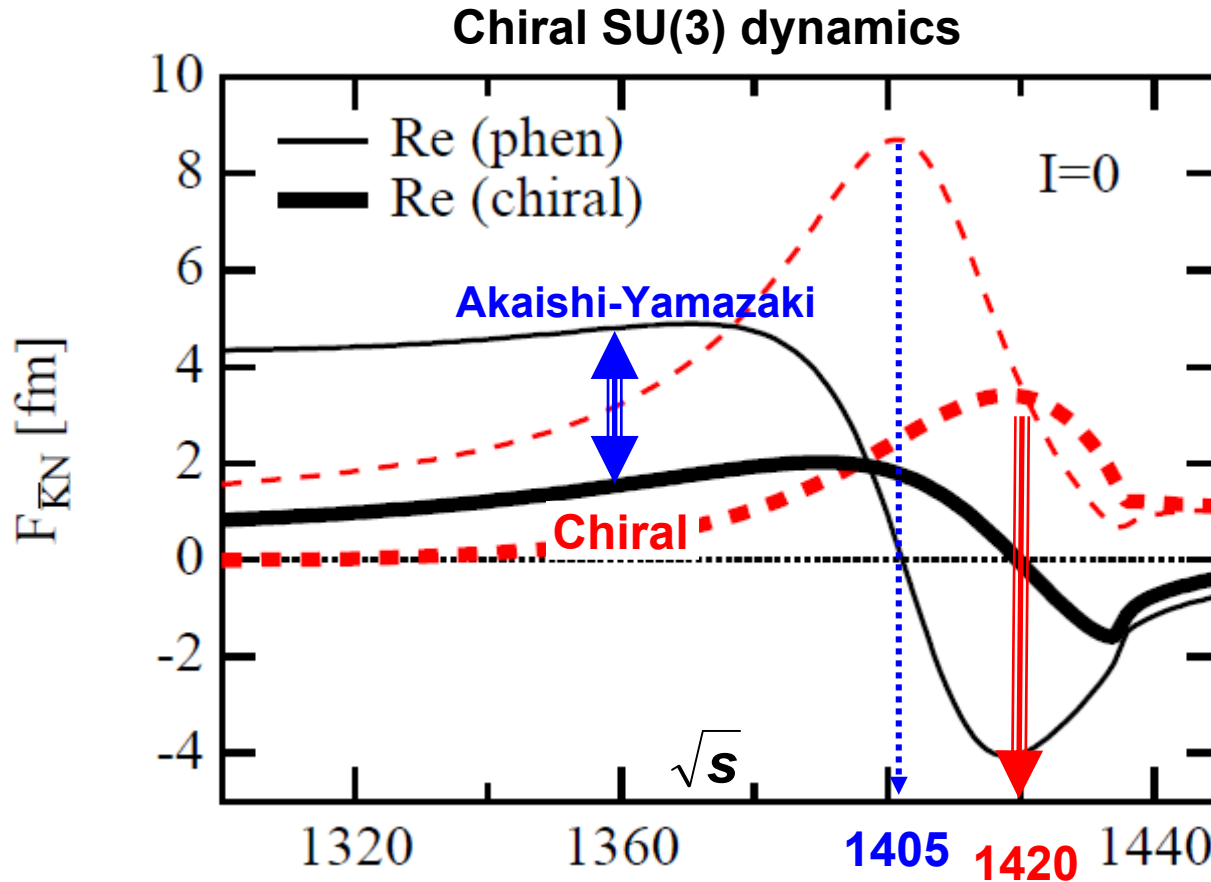
Revival of Heitler-London-Heisenberg picture



T. Yamazaki & Y. Akaishi,  
Proc. Japan Academy, B 83 (2007) 144

# $K^{\text{bar}}N$ scattering amplitude

T. Hyodo and W. Weise, Phys. Rev. C 77 (2008) 035204



"ORB": E. Oset, A. Ramos & C. Bennhold, Phys. Lett. B 527 (2002) 99

"HNJH": T. Hyodo, S.I. Nam, D. Jido & A. Hosaka, Phys. Rev. C 68 (2003) 018201

"BNW": B. Borasoy, R. Nissler & W. Weise, Eur. Phys. J. A 25 (2005) 79

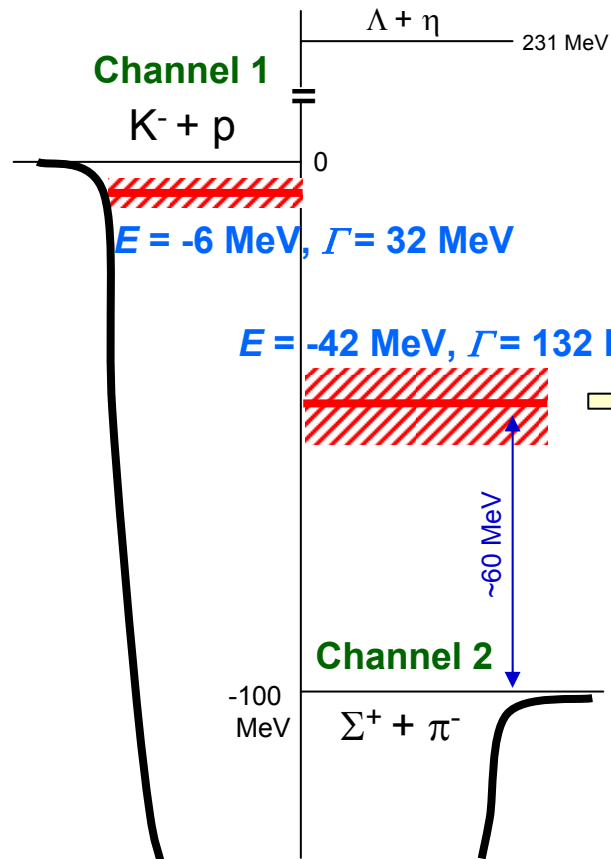
"BMN": B. Borasoy, U.G. Meissner & R. Nissler, Phys. Rev. C 74 (2006) 055201

"JOORM": D. Jido, J.A. Oller, E. Oset, A. Ramos & U.G. Meissner, Nucl. Phys. A 725 (2003) 181



# Double pole structure of $\Lambda(1405)$

D. Jido, J.A. Oller, E. Oset, A. Ramos & U.G. Meissner, Nucl. Phys. A 725 (2003) 181



$\Lambda(1405)$  consists of two poles,  
one of which is not  $K\bar{p}$  but  $\Sigma\pi$  pole.

What effects on experimental observables?

# Chiral SU(3) dynamics

## Weinberg-Tomozawa term

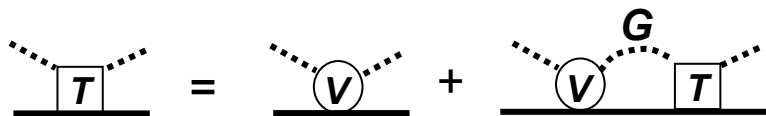
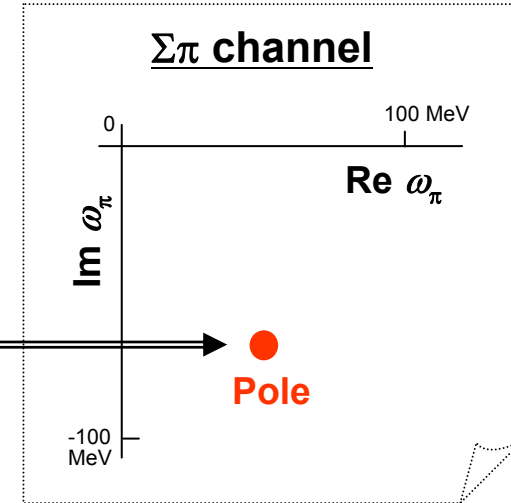
$$-c_{ij} \frac{\omega_i + \omega_j}{4f^2}$$

$$c_{ij}^{l=0} = \begin{bmatrix} 3 & -\sqrt{\frac{3}{2}} \\ -\sqrt{\frac{3}{2}} & 4 \end{bmatrix} \begin{array}{l} \bar{K}N \\ \pi\Sigma \end{array}$$

$$\omega_i \leftarrow \sqrt{s} - M_i$$

Energy dependent interaction

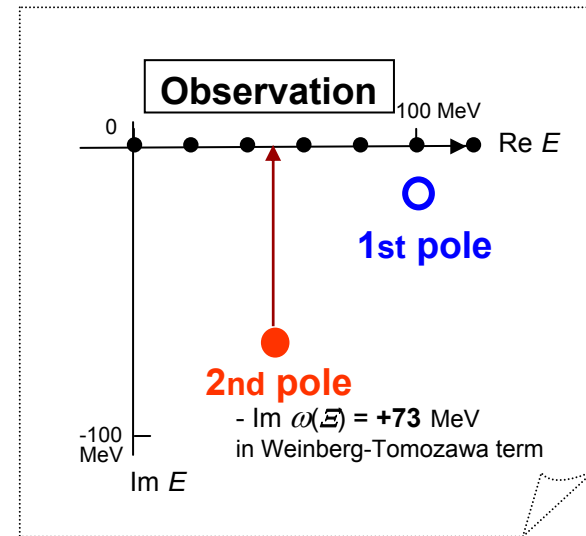
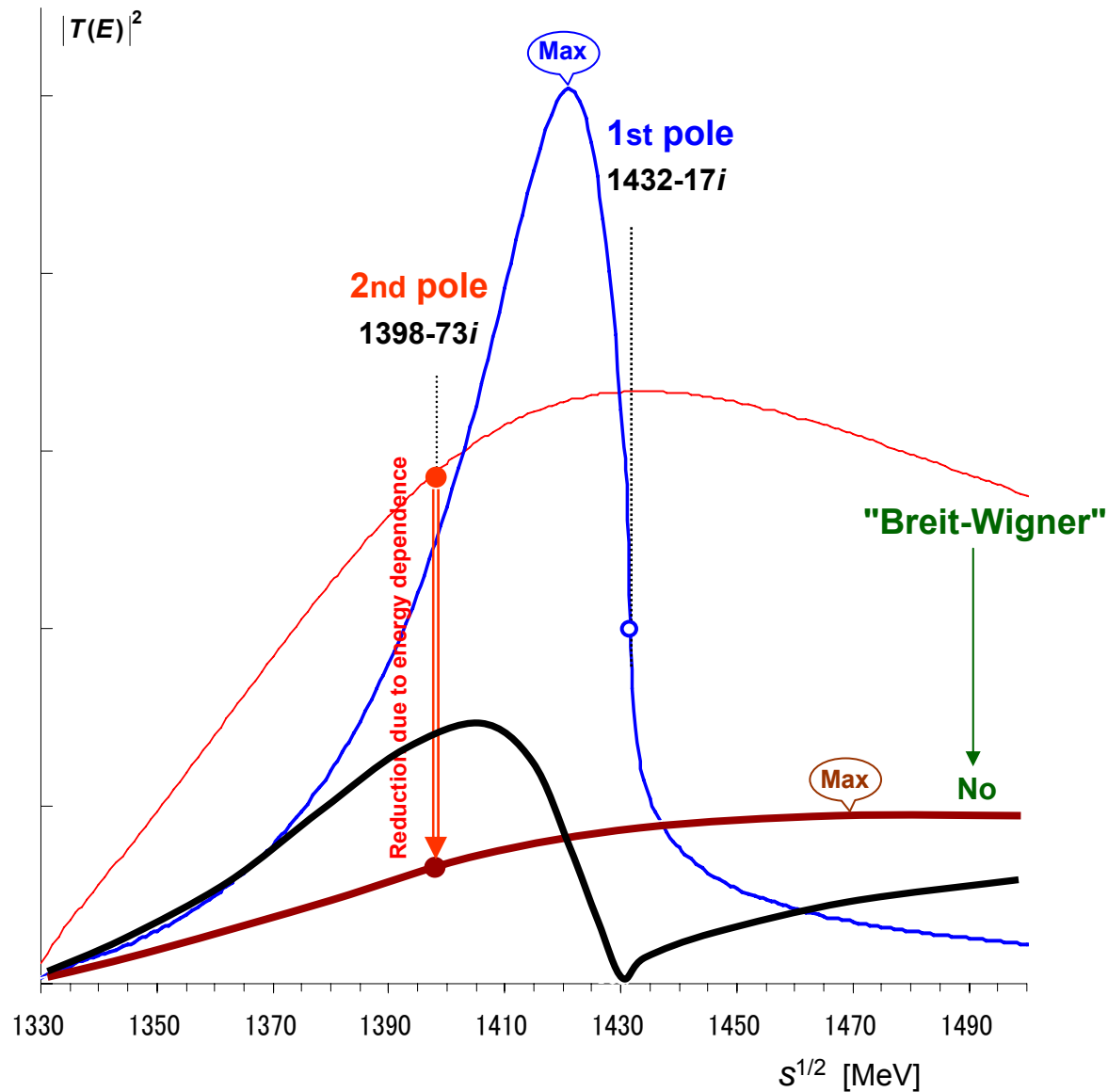
with a large positive imaginary part



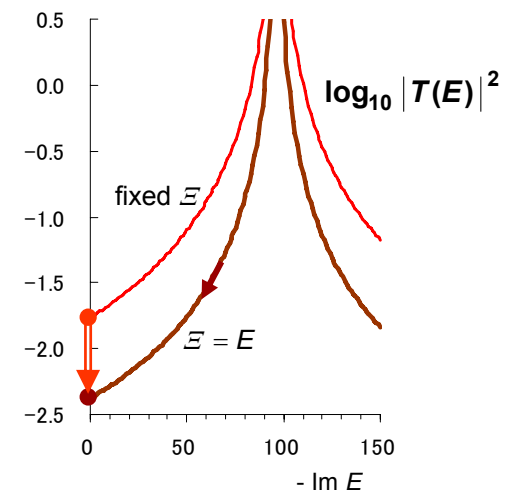
$$T = \frac{1}{1 - VG} V \quad ; \quad V = \{V_{ij}^{WT}\} \quad \text{Total 10 channels}$$

$$1 - VG = 0 \quad \Rightarrow \quad \text{Resonance (bound) states}$$

# $\Sigma\pi$ invariant mass spectrum of $\bar{K}N$ - $\Sigma\pi$ coupled system



$$|T(E)|^2 \approx \frac{\Gamma(\mathcal{E})^2 / 4}{(E - E_R(\mathcal{E}))^2 + \Gamma(\mathcal{E})^2 / 4}$$

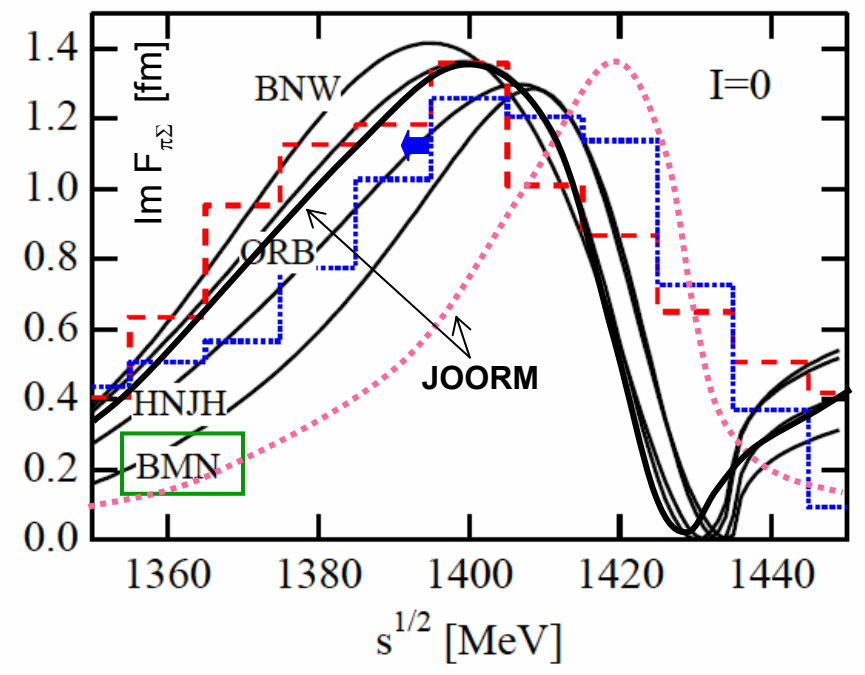
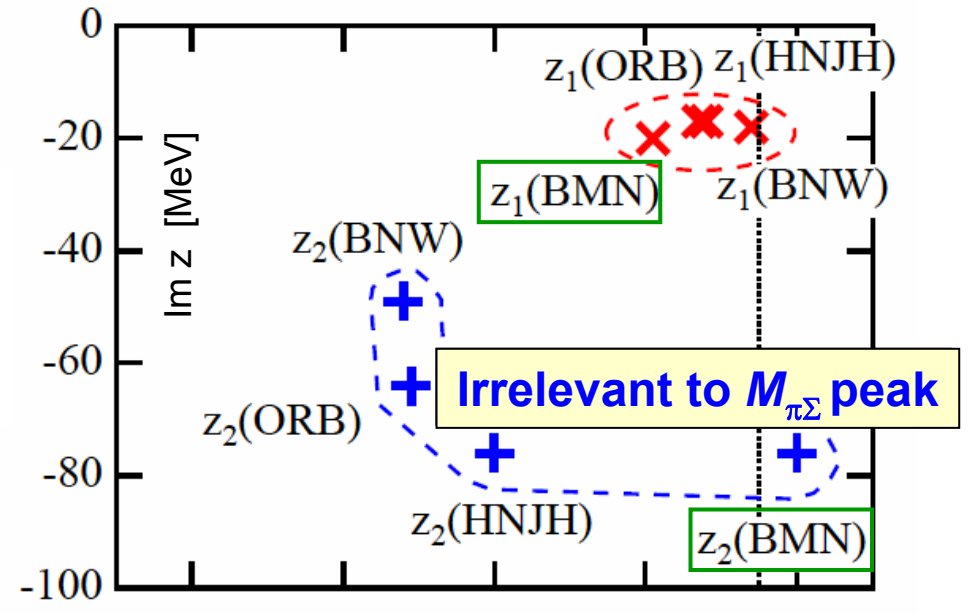
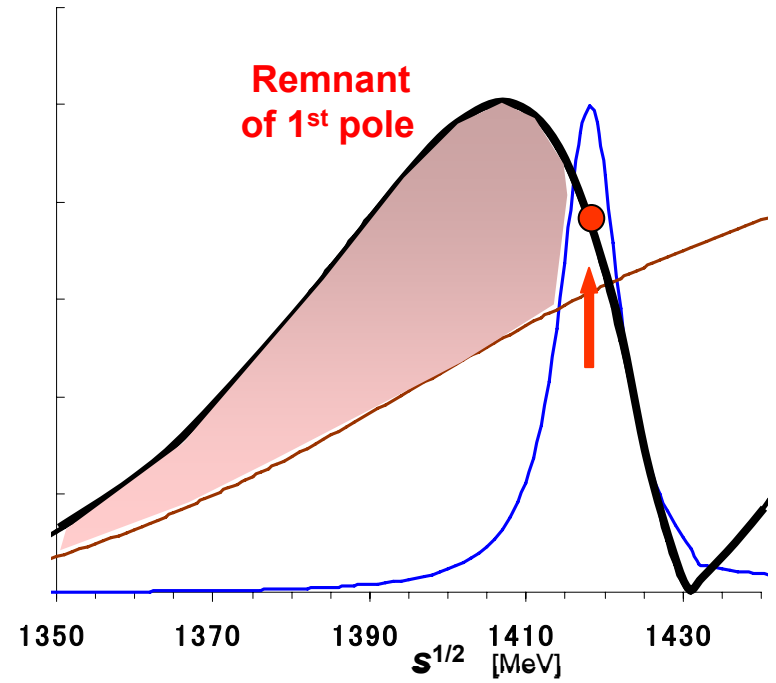


# Feshbach theory

1;  $\bar{K}N$   
2;  $\Sigma\pi$

$$|T_{22}(E)|^2 = \left| e^{i\delta_c} \sin \delta_c + \frac{\Gamma/2}{E - E_R + i\Gamma/2} \right|^2$$

Interference

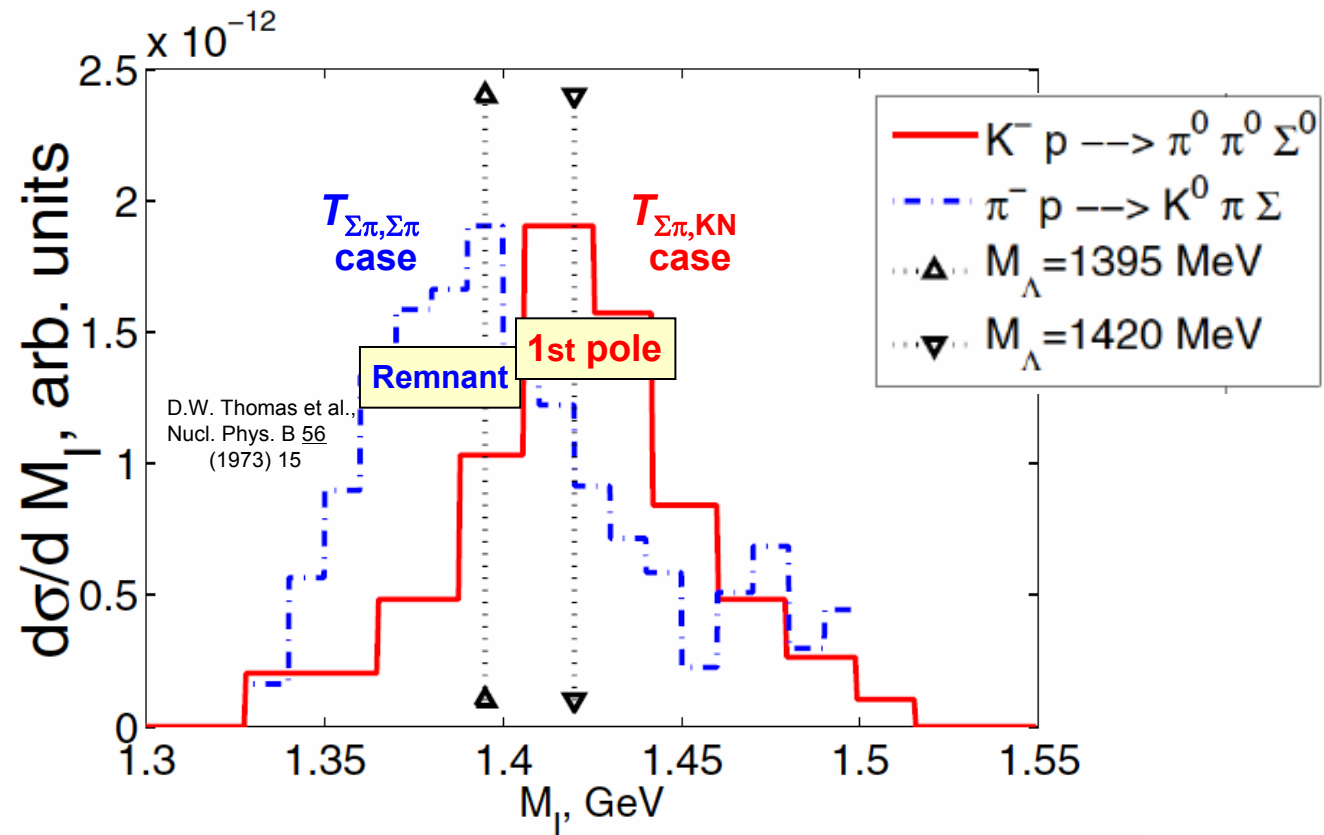
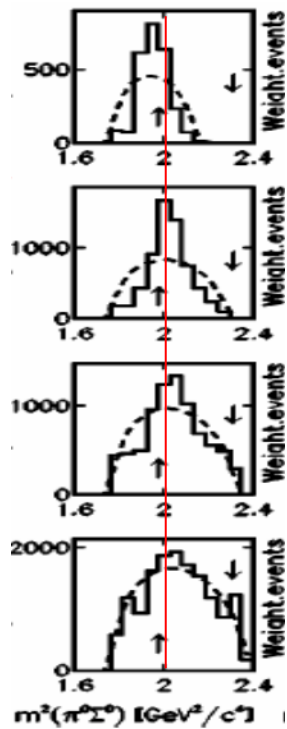


Logically impossible !



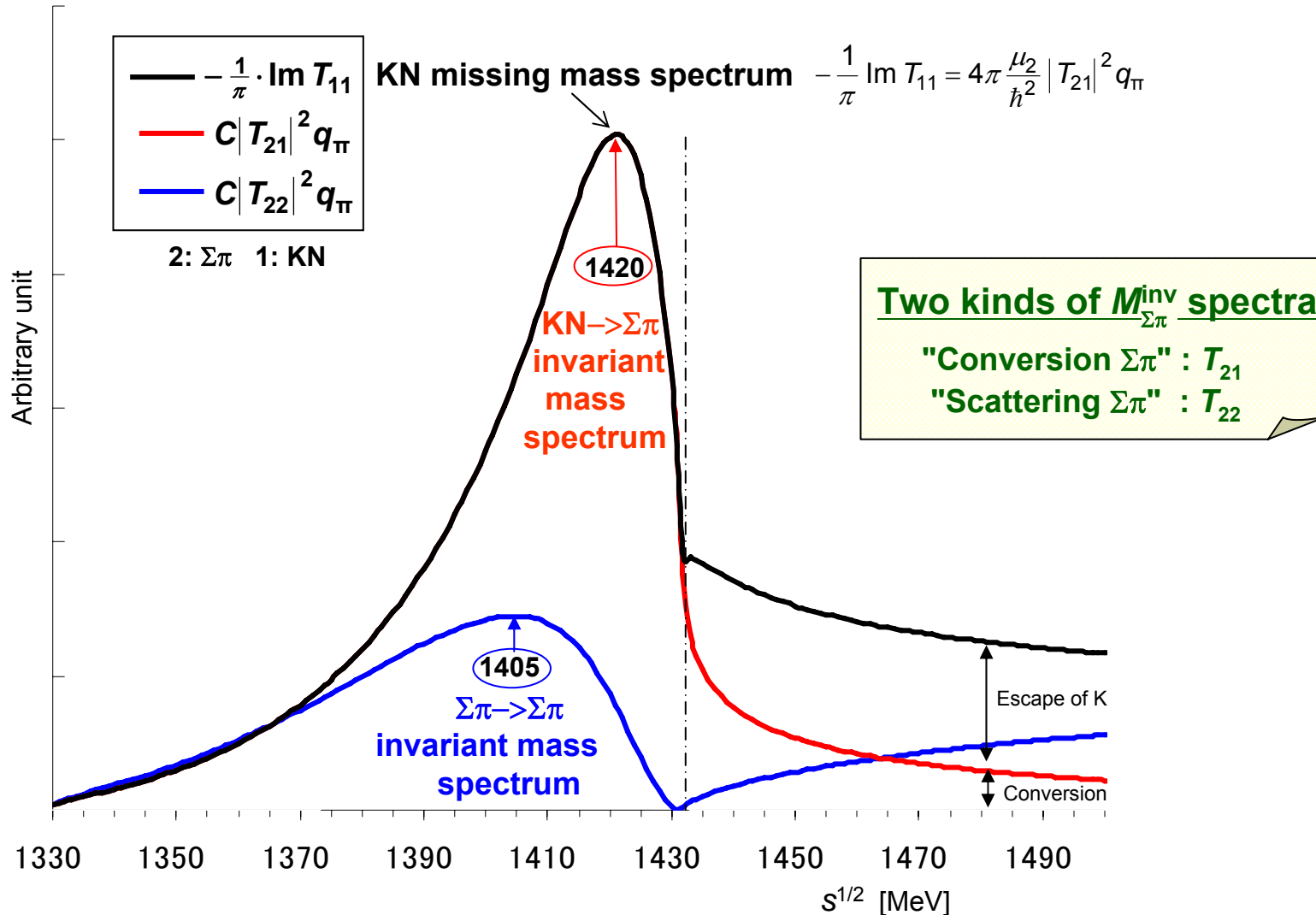
# Evidence for two-pole structure of $\Lambda(1405)$

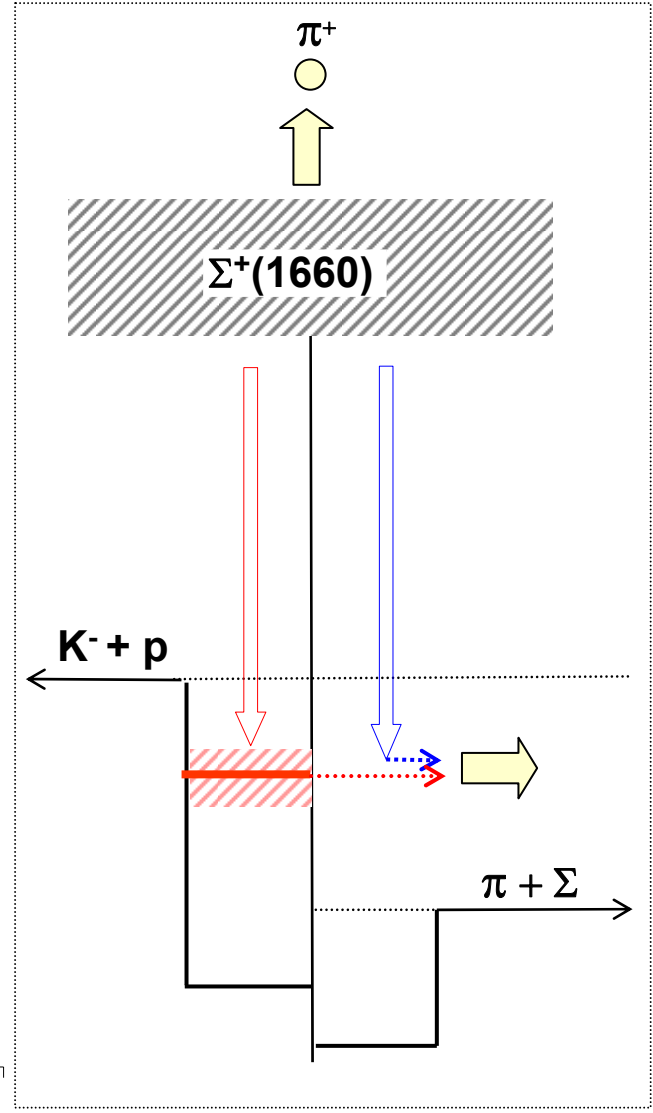
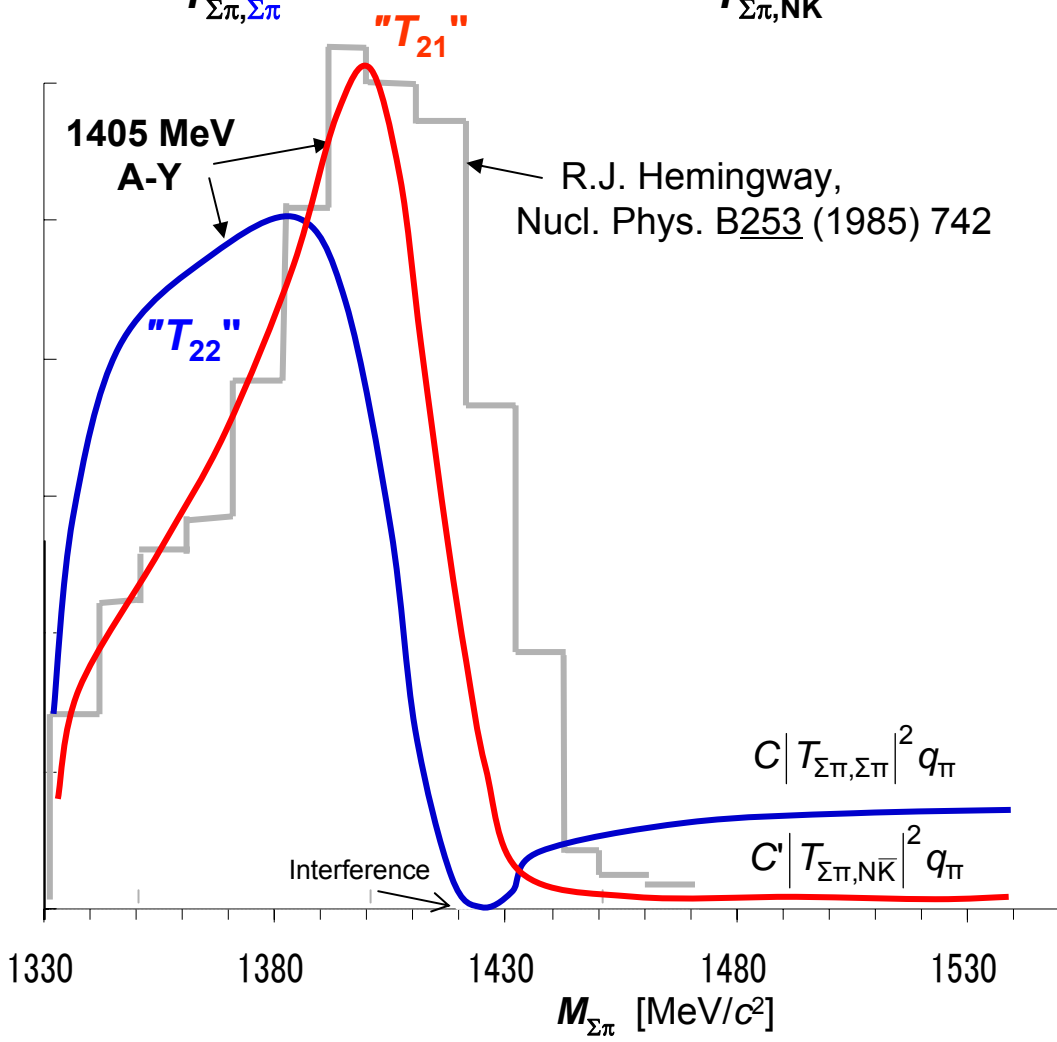
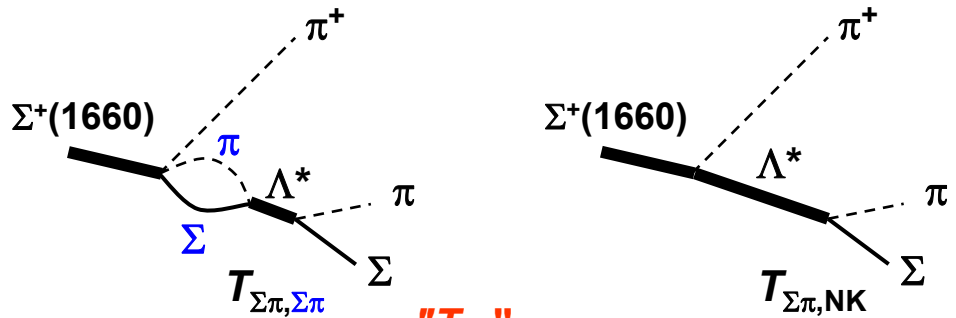
V.K. Magas, E. Oset and A. Ramos, Phys. Rev. Lett. **95** (2005) 052301



# Observables of $\bar{K}N-\Sigma\pi$ coupled system

Hyodo-Weise's chiral SU(3) dynamics





# $\Sigma\pi$ invariant-mass spectrum

1;  $\bar{K}N$   
2;  $\Sigma\pi$

" $T_{21}/T_{22}$ " problem

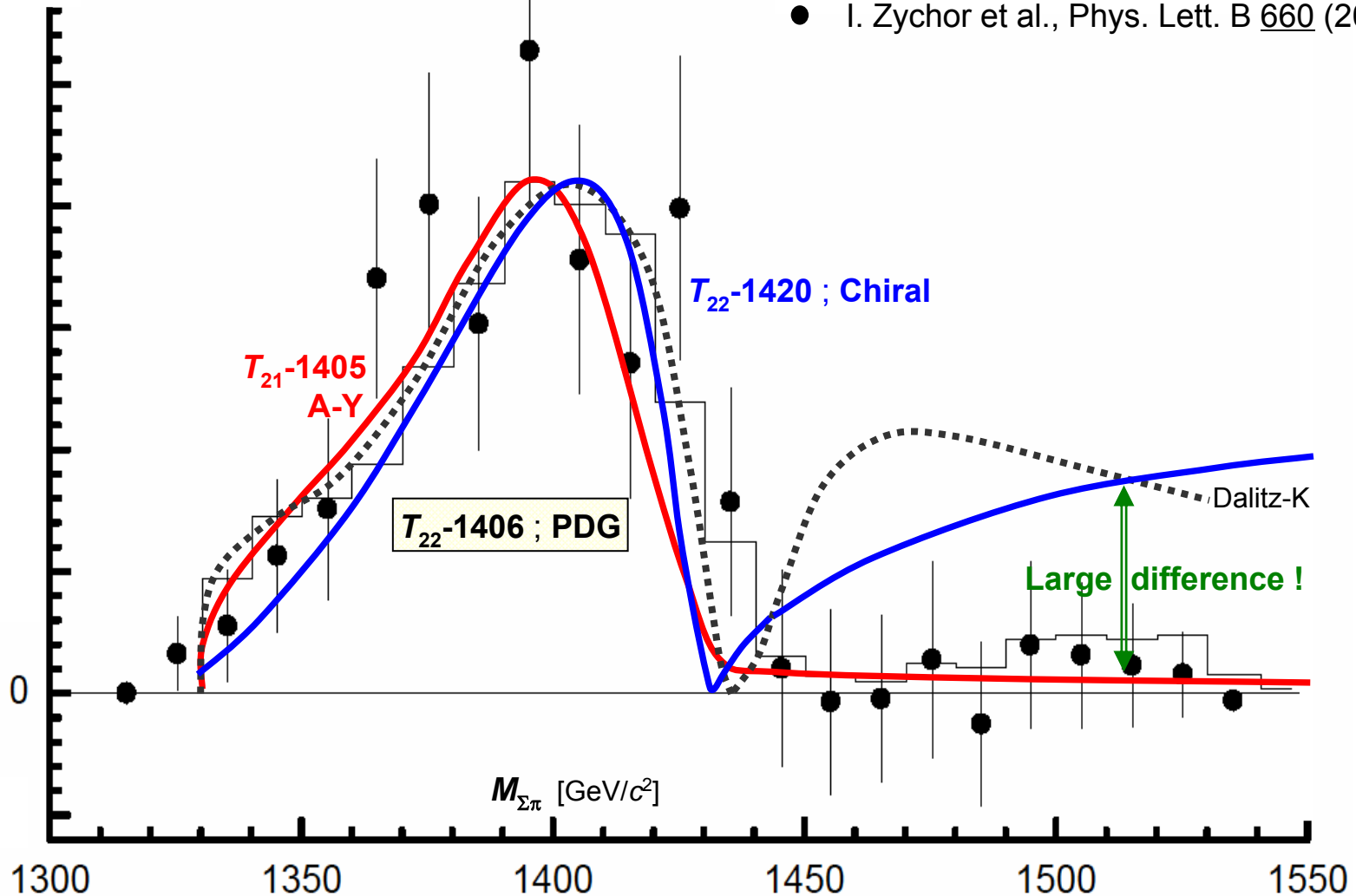
$$|\Sigma^+\pi^+\rangle = \sqrt{\frac{1}{3}}|I=0\rangle + \sqrt{\frac{1}{2}}|I=1\rangle + \sqrt{\frac{1}{6}}|I=2\rangle$$

$$|\Sigma^0\pi^0\rangle = -\sqrt{\frac{1}{3}}|I=0\rangle + \sqrt{\frac{2}{3}}|I=2\rangle$$

$$|\Sigma^-\pi^+\rangle = \sqrt{\frac{1}{3}}|I=0\rangle - \sqrt{\frac{1}{2}}|I=1\rangle + \sqrt{\frac{1}{6}}|I=2\rangle$$

— R.J. Hemingway, Nucl. Phys. B253 (1985) 742

● I. Zychor et al., Phys. Lett. B 660 (2008) 167





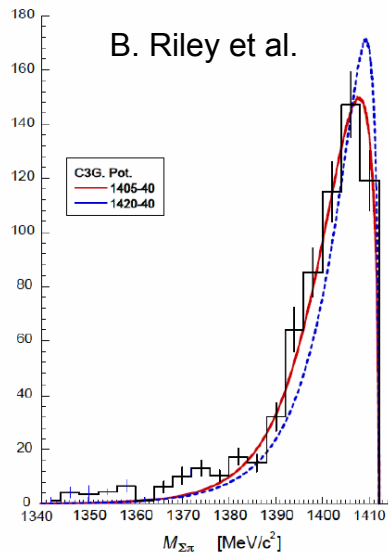
# Concluding remarks

The  $\Lambda^*$  resonance forms the basic structure of  $K\text{-pp}$ .

The  $\Lambda(1405)$  is regarded to be of single-pole nature, since the 2<sup>nd</sup> pole is irrelevant to any experimental peak.

It is virtually important to distinguish the mass of  $\Lambda(1405)$ , 1405 MeV or 1420 MeV, by considering " $T_{21}/T_{22}$ " problem.

## Stopped $K^-$ on $^4\text{He}$



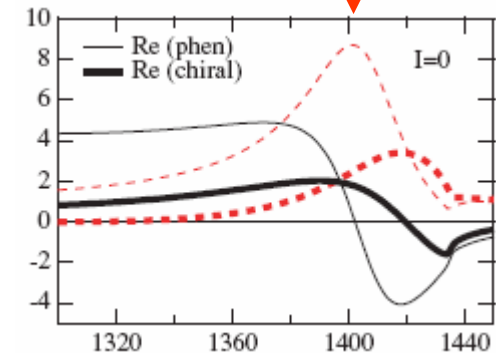
This  $T_{21}$  datum supports the "1405 Ansatz".

$$M_{\Lambda} c^2 = 1405.5_{-1.0}^{+1.4} \text{ MeV and } \Gamma = 25.6_{-3}^{+4} \text{ MeV}$$

J. Esmaili et al.; arXiv:0906.0505

$$|T_{21}|^2 \text{Im}G_2 = \text{Im}T_{11}$$

Stopped  $K^-$  on D T. Suzuki et al.  
would provide a decisive datum.

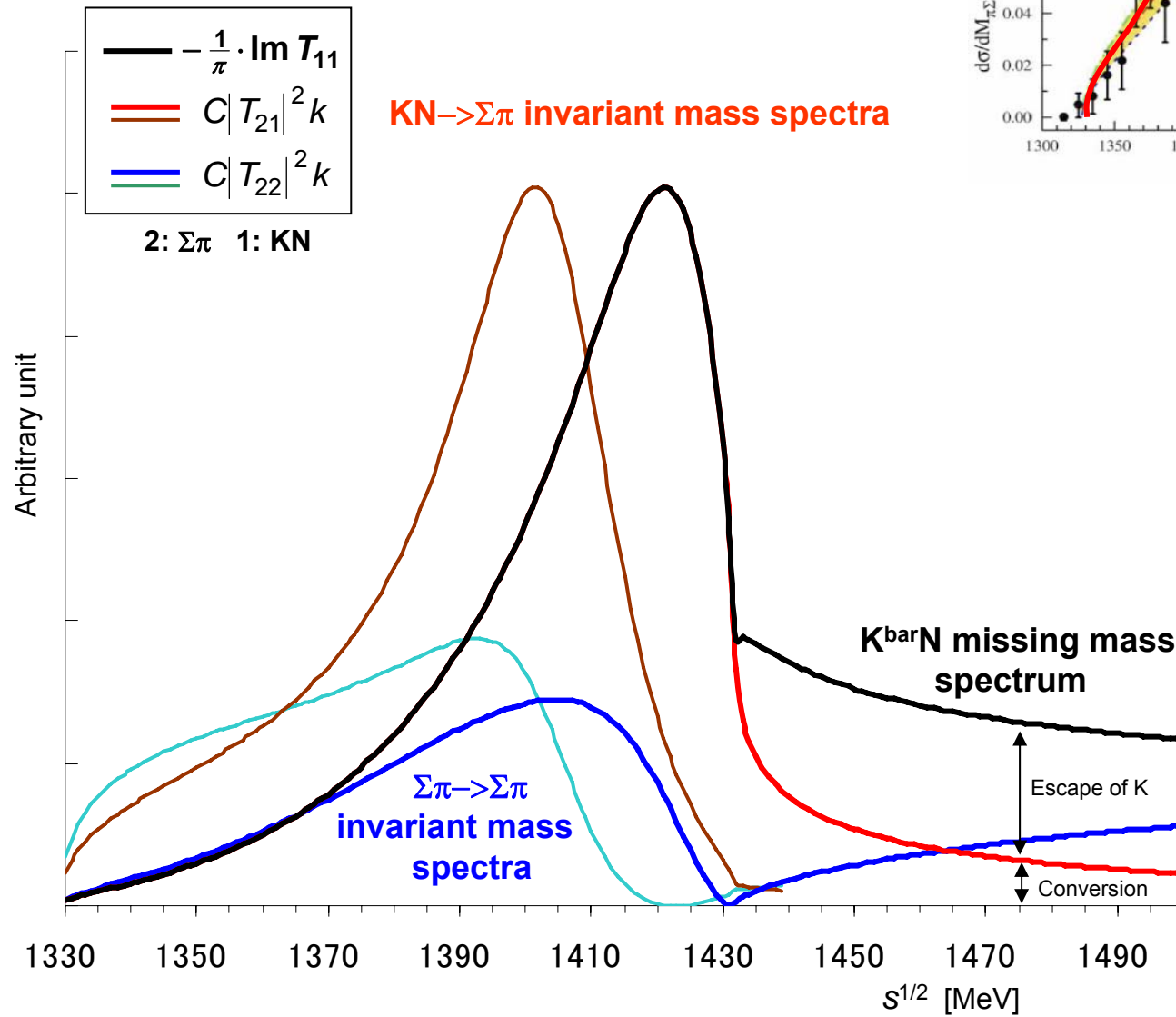
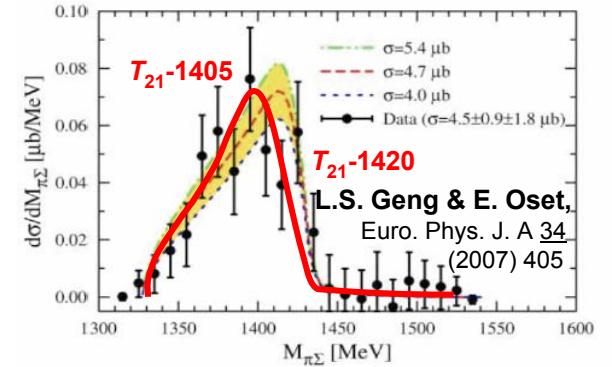


## **Acknowledgments**

**J. Esmaili, K.S. Myint**

**M. Kawai, O. Morimatsu**

**Thank you very much!**



# $K^-d \rightarrow \Sigma^- \pi^+ n$

$P_K = 760 \text{ MeV}/c$

O. Braun et al.,  
Nucl. Phys. B129 (1977) 1

