

# The nature of $\Lambda(1405)$ in chiral dynamics

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## $\Lambda(1405)$

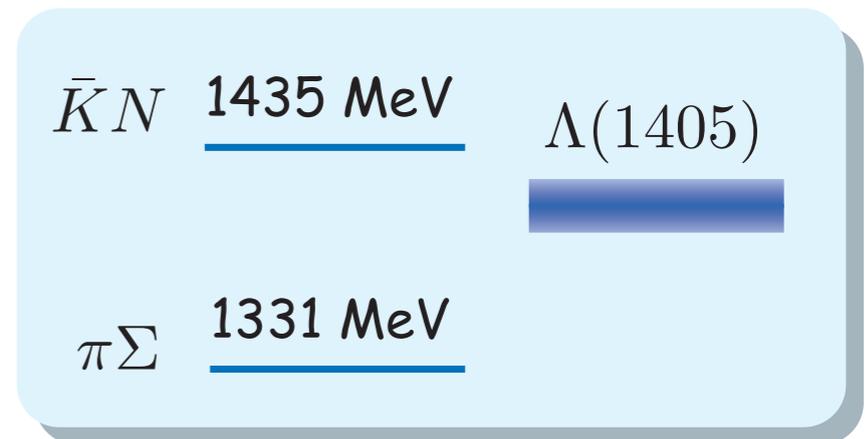
excited state of  $\Lambda$  with  $J^P=(1/2)^-$

just below threshold of  $K^{\text{bar}}N$

important for study of  $K^{\text{bar}}$  at subthreshold, bound kaon

decay to  $\pi\Sigma$

$\pi\Sigma$  and  $K^{\text{bar}}N$  dynamics



# Coupled-channels approach in chiral dynamics

Chiral unitary model

For example,  
Kaiser, Siegel, Weise, NPA594, 325 (95)  
Oset, Ramos, NPA635, 99 (98)  
Oller, Meissner, PLB500, 263 (01)

**a powerful theoretical framework to describe hadronic resonances from hadron dynamics**

**chiral perturbation theory**

low-energy effective theory of QCD

give fundamental interaction of meson-baryon

**scattering theory (N/D method)**

**analyticity** and unitarity

general form of scattering amplitude

Lippmann-Schwinger eq.

$$T = V + VGT$$

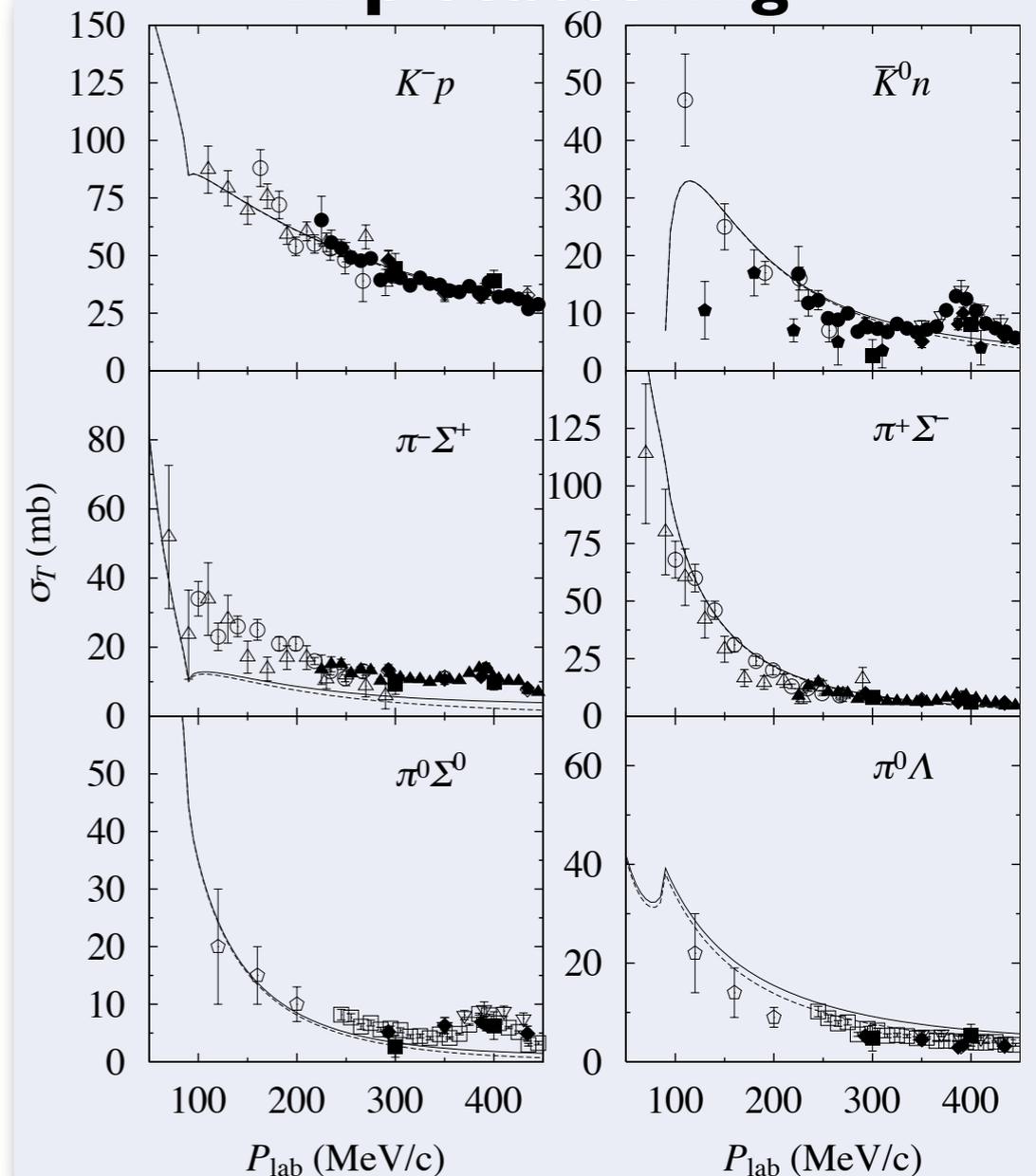
$\bar{K}N, \pi\Sigma, \eta\Lambda, K\Sigma, \pi\Lambda, \eta\Sigma$

**reproduce K-p scattering**

**generate dynamically**

**s-wave  $\Lambda(1405)$  resonances**

## K-p scattering



DJ, Oset, Ramos, PRC66, 055203 (02)

# Consequences of chiral unitary model

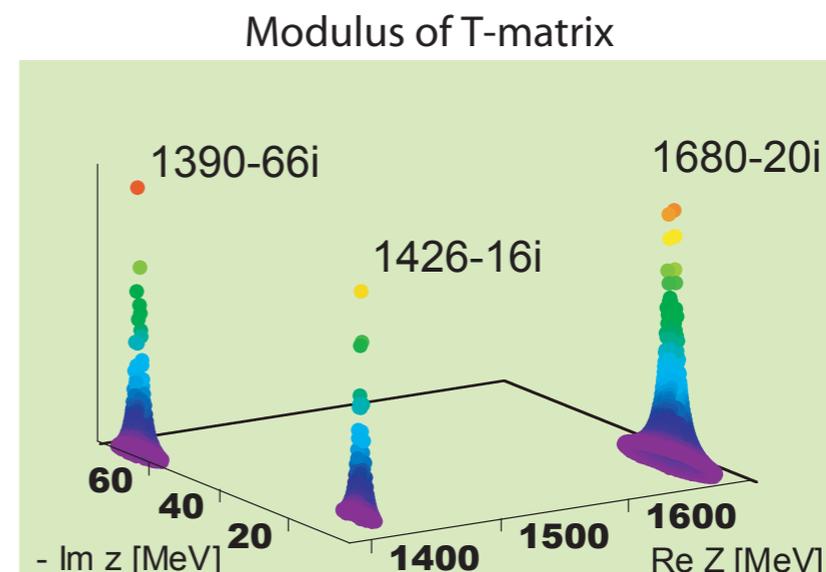
## 1. $\Lambda(1405)$ is a superposition of two states having different properties

pole 1: **1390 MeV**, width **132 MeV**

strongly couples to  **$\pi\Sigma$  state**

pole 2: **1426 MeV**, width **32 MeV**

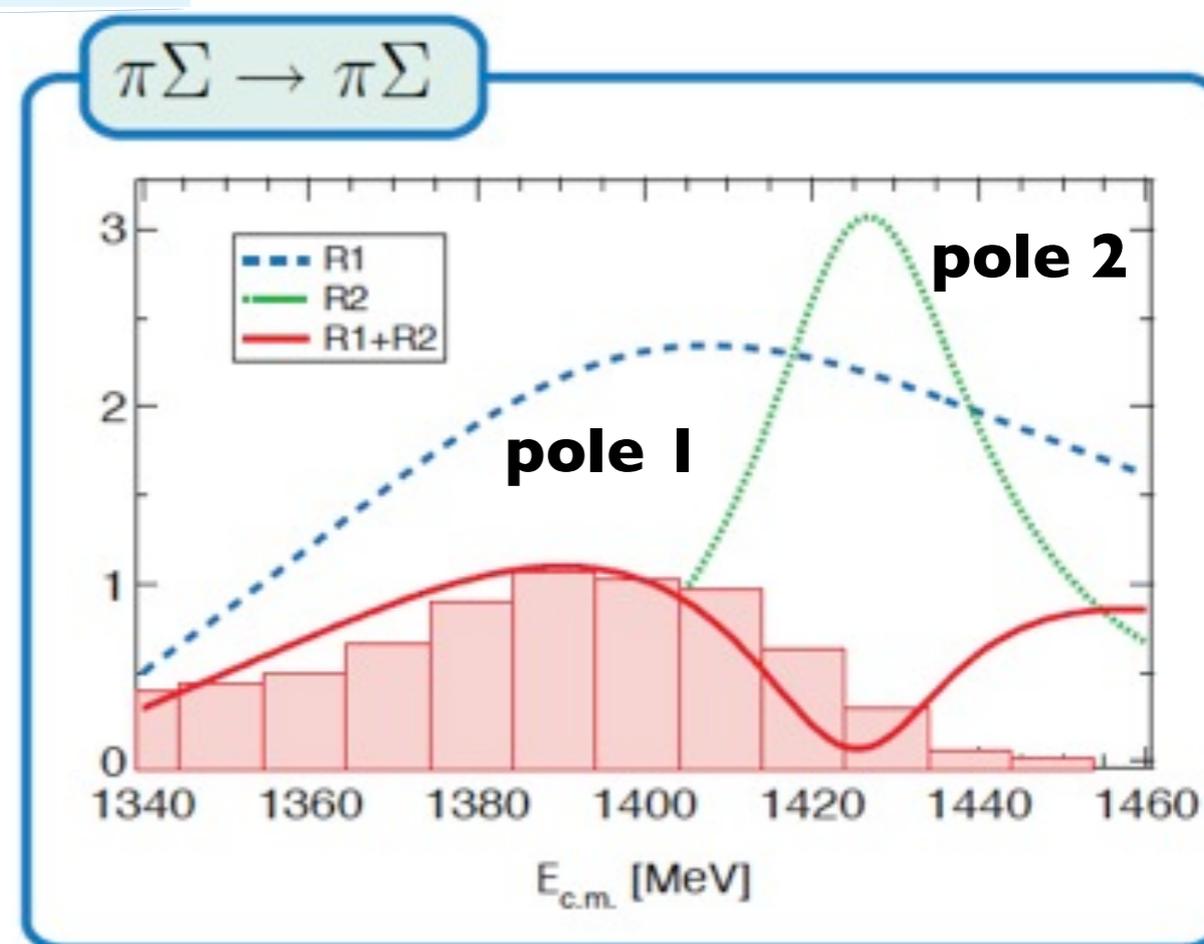
dominantly couples to  **$K^{\text{bar}}N$  state**



**Two state cannot be seen separately due to the widths.**

**Spectrum of  $\Lambda(1405)$  is given by interference of these poles.**

DJ, Oller, Oset, Ramos, Meissner  
NPA725, 181 ('03)



# Double pole structure of $\Lambda(1405)$

reason of existence of two poles: **two attractive channels in  $l=0$**

group theoretically SU(3) singlet and octet

physically  $K^{\text{bar}}N$  and  $\pi\Sigma$

according to ChPT

**$-\Lambda(1405)$  is essentially described by two channels,  $K^{\text{bar}}N$  and  $\pi\Sigma$ .**

- in single channel without channel couplings

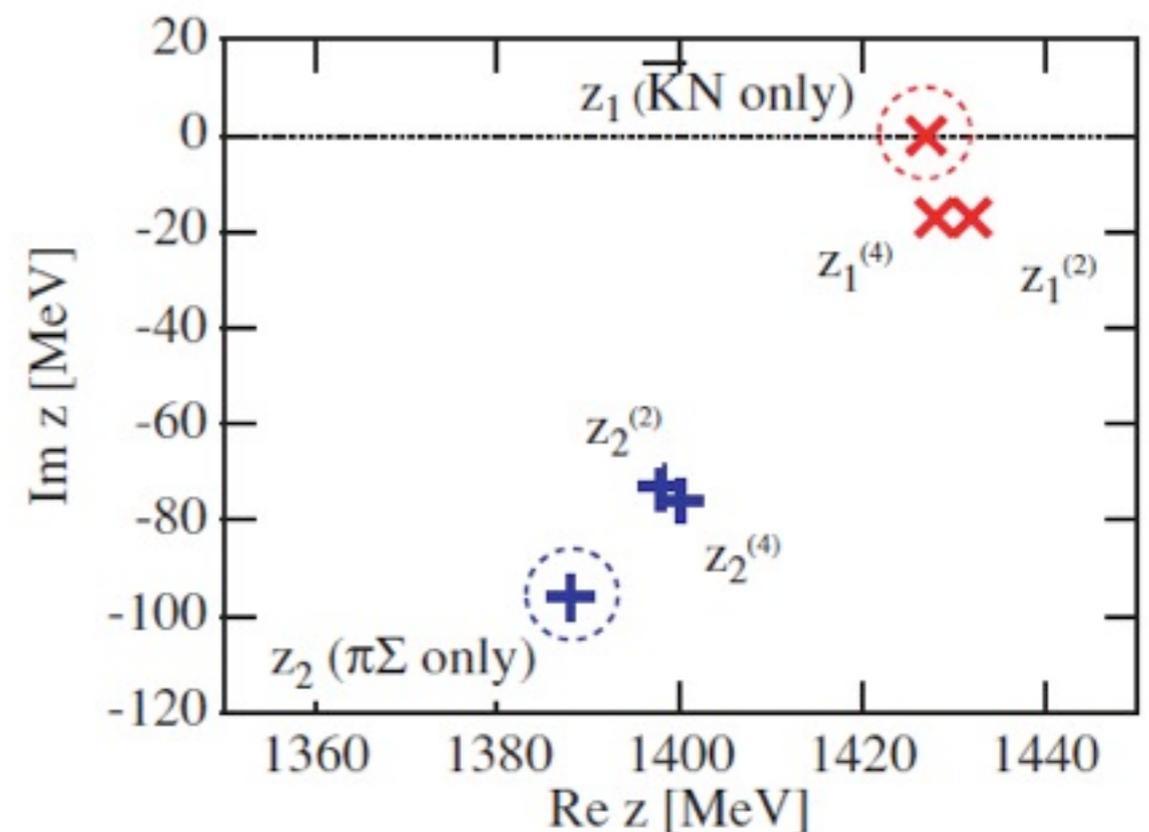
**$K^{\text{bar}}N$  bound state**

**$\pi\Sigma$  resonance**

s-wave resonance in single channel due to energy-dependent potential

- coupled channel gives width to bound state

Hyodo, Weise, PRC77, 035204 ('08)



comp. of single, two-, four-channel

for physics on the real axis

**$K^{\text{bar}}N$  bound state and  $\pi\Sigma$  strong correlation**

are the ingredients of  $\Lambda(1405)$ . **Double pole structure**

# Consequences of chiral unitary model

## 2. Resonance position of $\Lambda(1405)$ depends on channels

pole 1: **1390 MeV, width 132 MeV**

strongly couples to  **$\pi\Sigma$  state**

pole 2: **1426 MeV, width 32 MeV**

dominantly couples to  **$K^{\text{bar}}N$  state**

Due to the presence of two poles having different properties, the  $\pi\Sigma$  invariant mass spectrum (peak position of  $\Lambda(1405)$ ) depends on the initial channel.

**$\pi\Sigma$  invariant mass spectrum ( $I=0$ )**

$$\frac{d\sigma}{dM_{\pi\Sigma}} = A|T|^2 q_{\text{c.m.}}^{\pi\Sigma}$$

The resonance positions depend on the channels by which  $\Lambda(1405)$  is produced.

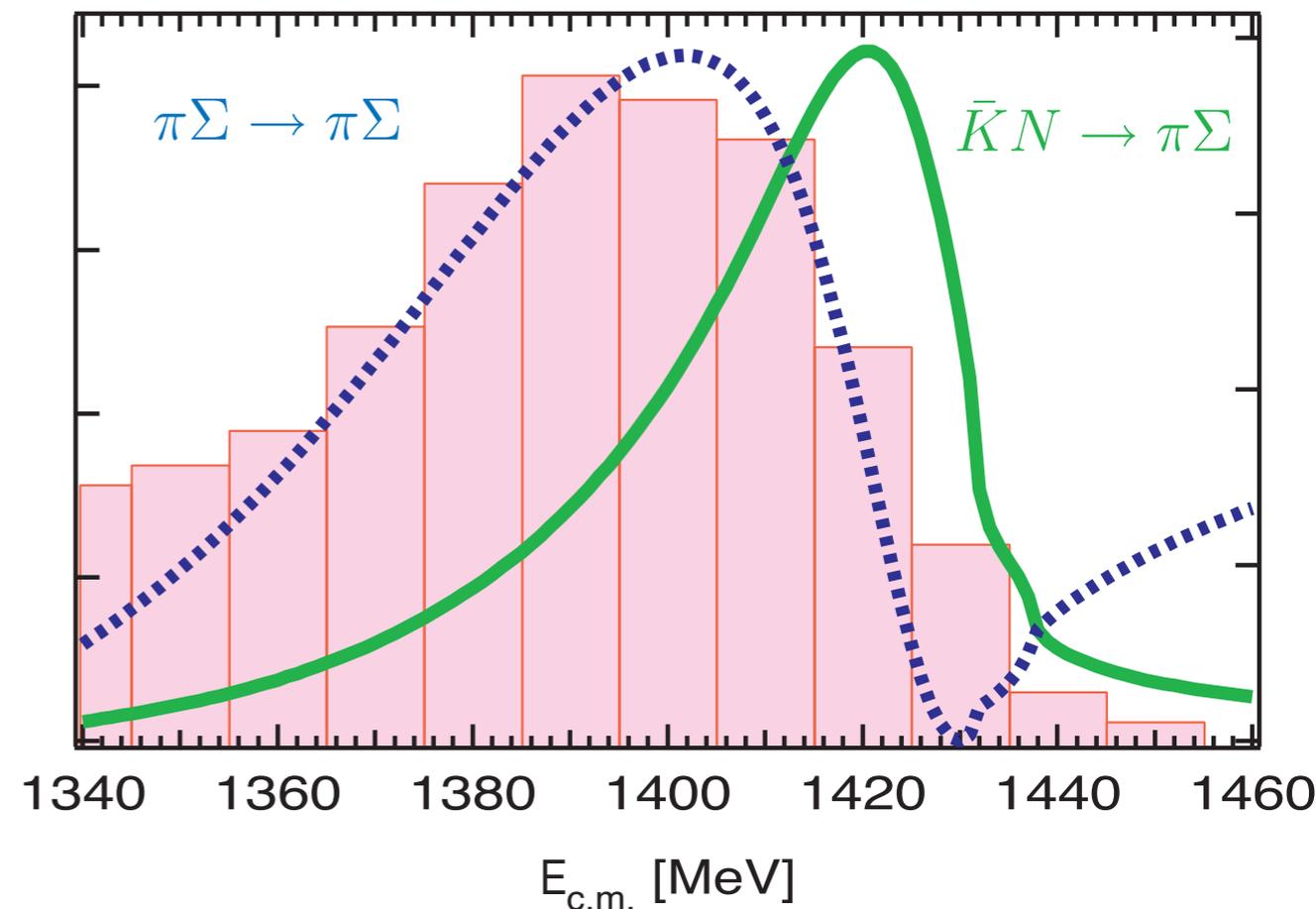
$\Lambda(1405)$  as  $K^{\text{bar}}N$  quasibound state

1420 MeV  $\rightarrow$  binding energy 15 MeV

1405 MeV  $\rightarrow$  binding energy 30 MeV

the heights are adjusted

$\pi\Sigma$  Mass distribution



**peak at 1420 MeV  
in  $K^{\text{bar}}N$  channel**

DJ, Oller, Oset, Ramos, Meissner, NPA725, 181 ('03)

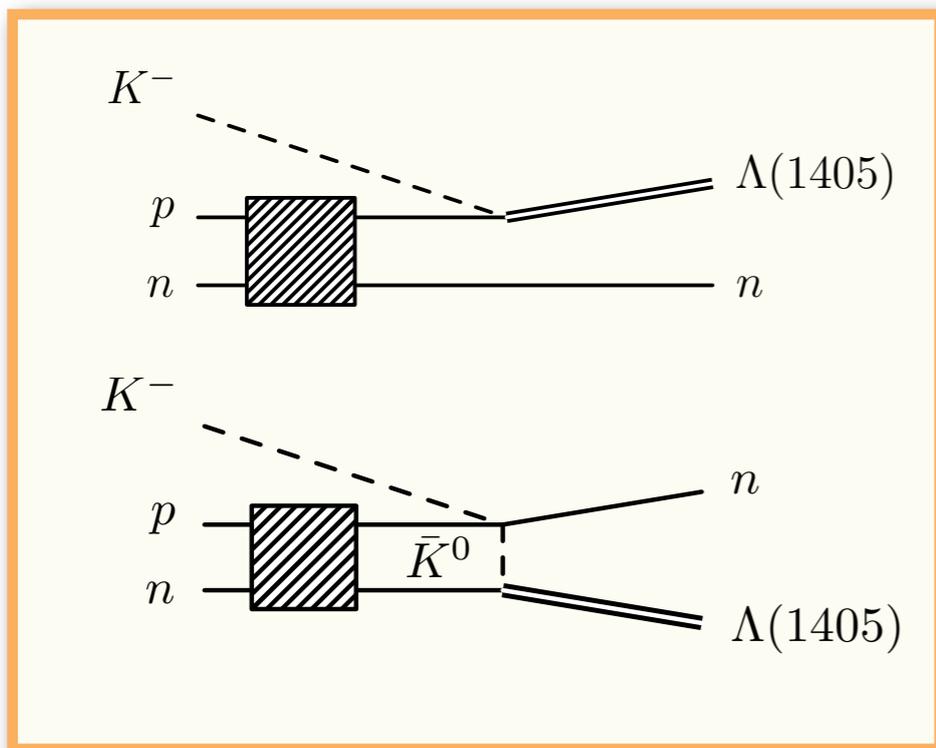
# $\Lambda(1405)$ in $K^{\text{bar}}N$ channel

## Want to see $\Lambda(1405)$ produced by $K^{\text{bar}}N$ !!

$\Lambda(1405)$  is located below the  $K^{\text{bar}}N$  threshold

cannot be produced by direct reaction  $\bar{K}N \rightarrow \Lambda(1405)$

### indirect reaction



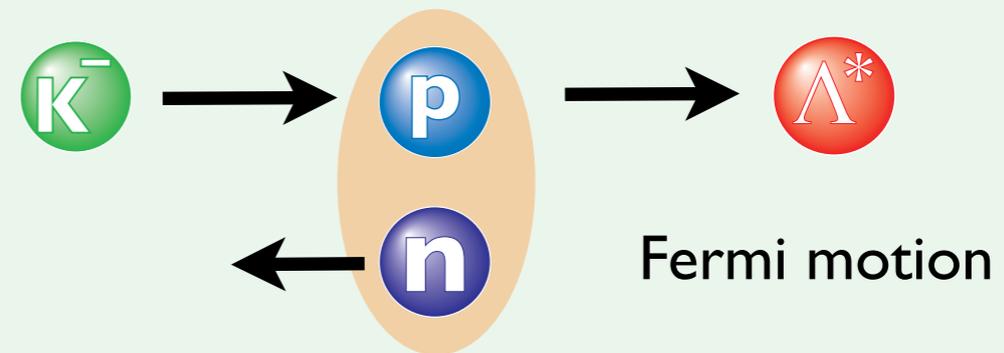
**flow of strangeness is clear**

$\Lambda(1405)$  is produced by  $K^{\text{bar}}N$



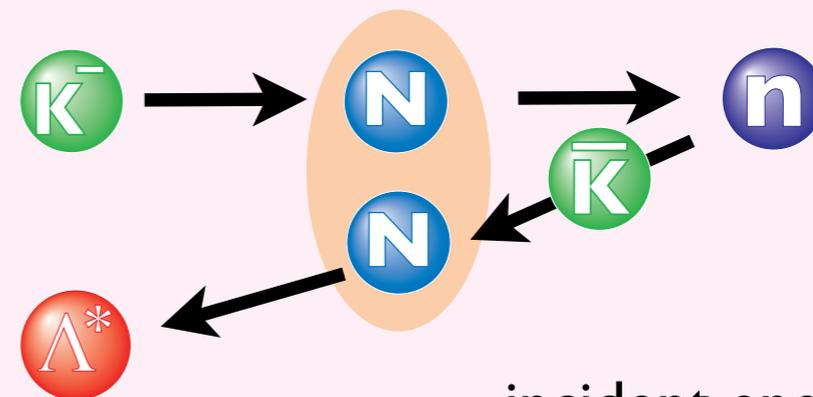
- single step

$\Lambda^*$  forward



- double step

$\Lambda^*$  backward



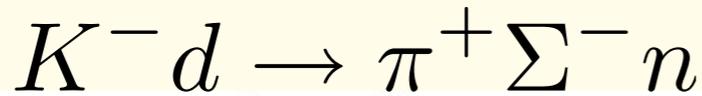
incident energy is shared by two nucleons

**dominant**

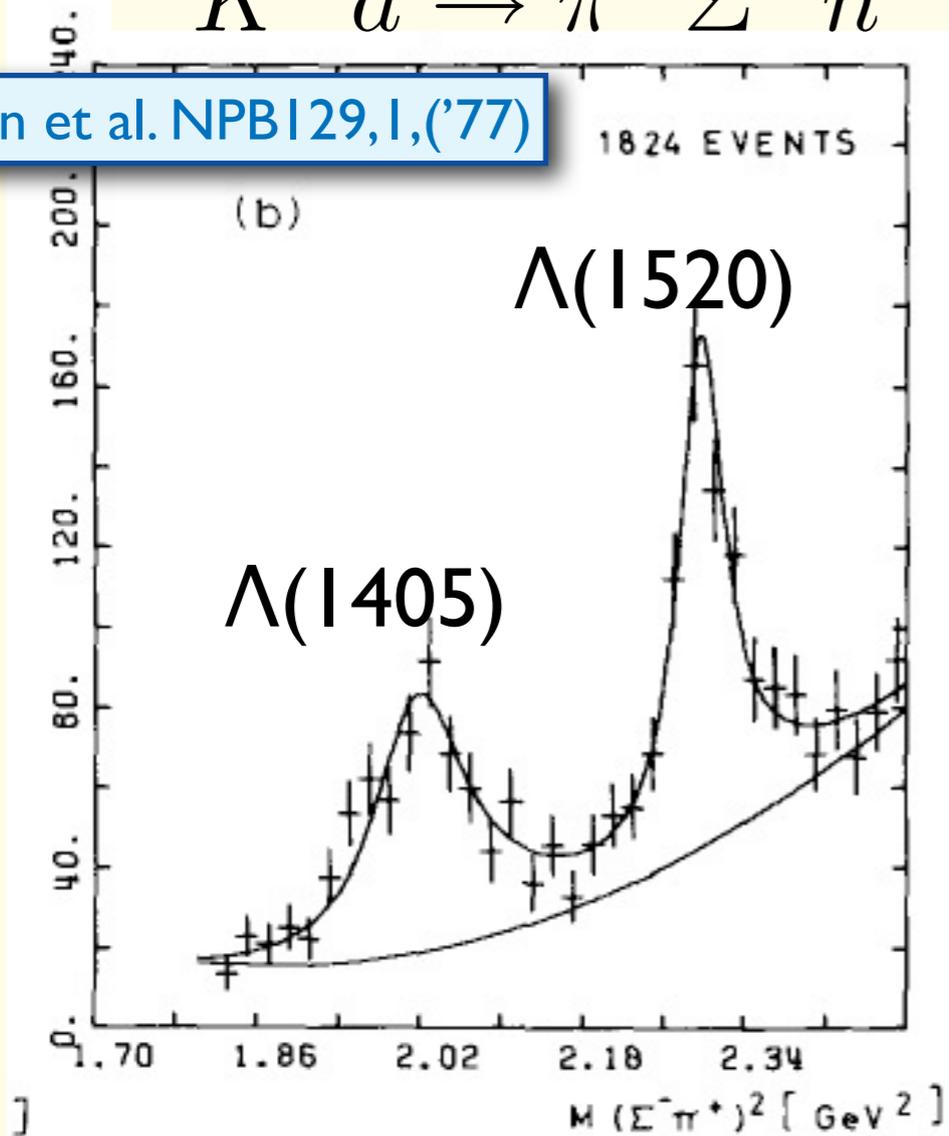
# $\Lambda(1405)$ in $K^{\text{bar}}N$ channel

**Experiment** bubble chamber initial K momentum 686 ~ 844 MeV/c

$\pi\Sigma$  invariant mass spectrum



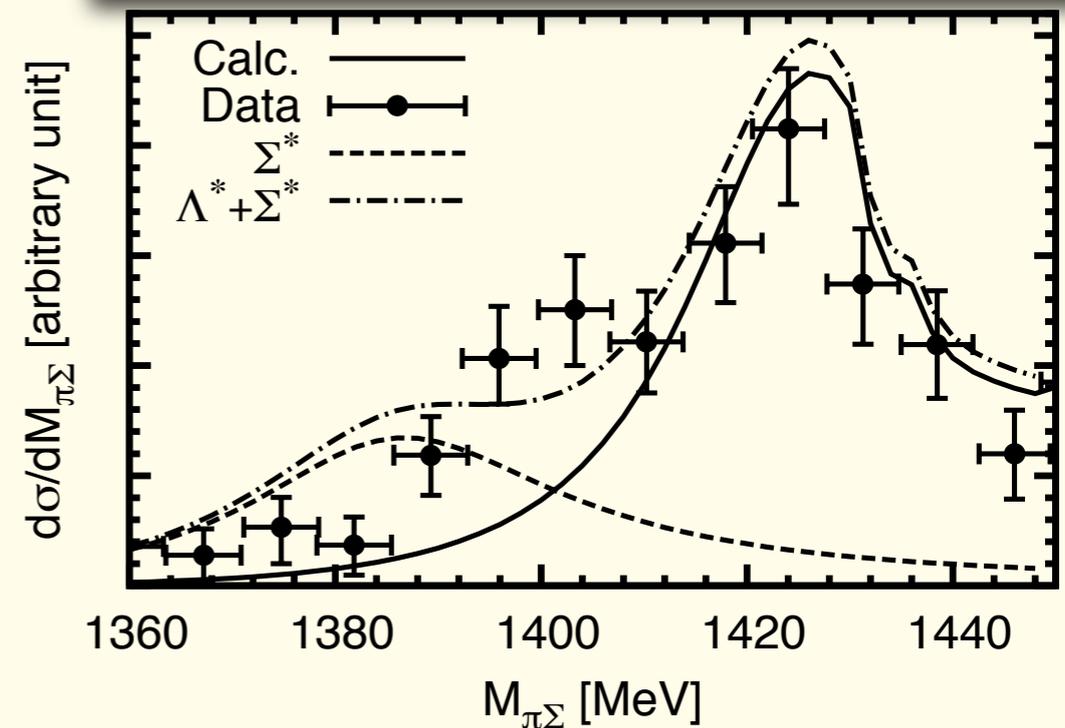
Braun et al. NPB129,1,(77)



**peak position 1420 MeV**

**theoretical calculation in ChUM**

DJ, Oset, Sekihara, accepted in Eur.Phys.J.A.



**production cross section of  $\Lambda(1405)$**

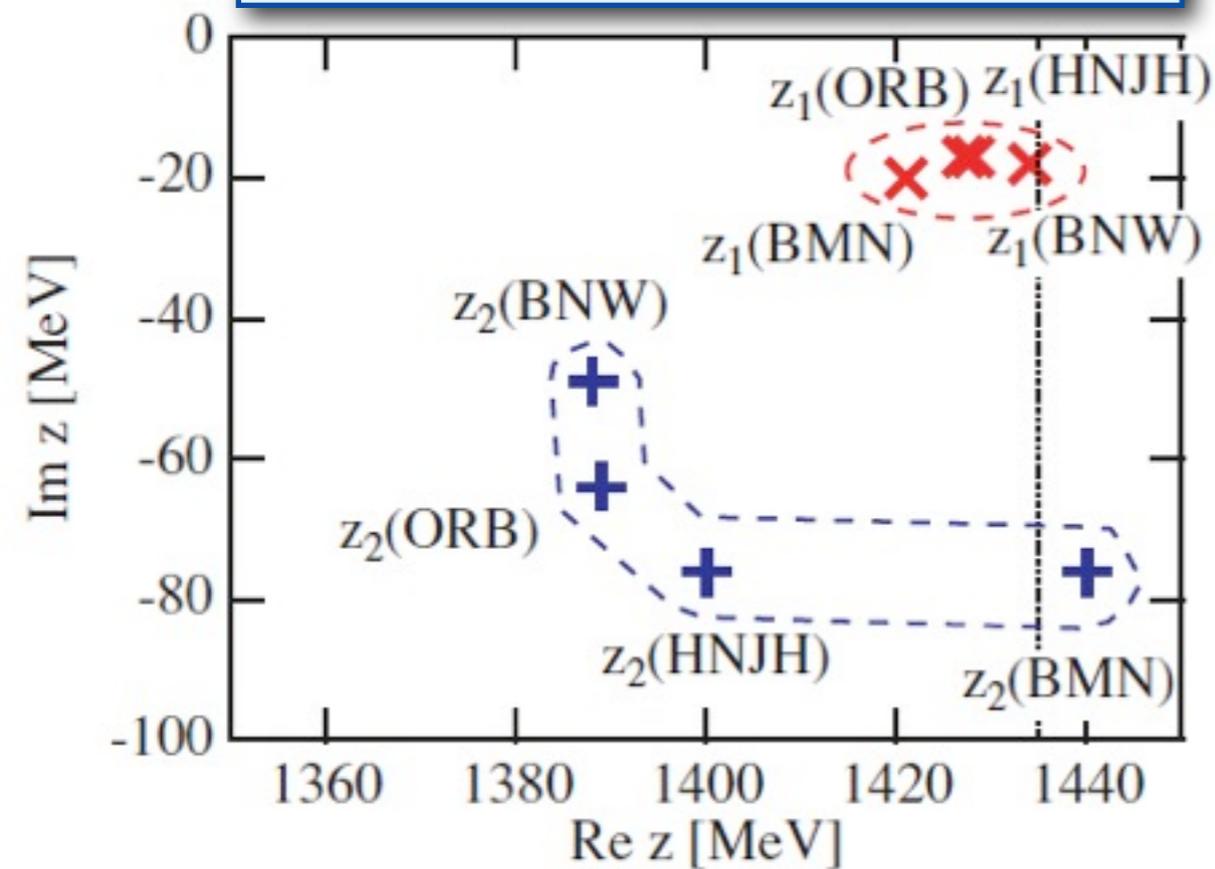
385  $\mu\text{b}$  @ 800MeV/c (exp.  $410 \pm 100 \mu\text{b}$ )  
agrees with data in shape and size  
bump around 1385 MeV is found to be from  $\Sigma^*$   
inclusion of  $\Sigma^*$  does not distort the shape.

Proposal of J-PARC experiment  
H. Noumi, Poster M15, today

# Influence of $\pi\Sigma$ correlation to $\Lambda(1405)$

Model dependence of pole positions in chiral coupled channels approach

Hyodo, Weise, PRC77, 035204 ('08)



**higher pole**

dominantly couples to  $K^{\text{bar}}N$

less model-dependent  
around 1420 MeV

**constrained by  $K^{\text{bar}}N$  scatt. data**

**lower pole**

strongly couples to  $\pi\Sigma$

strongly model-dependent  
lack of  $\pi\Sigma$  scattering data

**need any data for  $\pi\Sigma$  scattering**

Model calculation of  $I=0$

$\pi\Sigma$  scattering length (fm)

	BNW	ORB	HNJH	BMN
$\pi\Sigma$ scattering length (fm)	<b>0.517</b>	<b>0.789</b>	<b>0.692</b>	<b>0.770</b>

Ikeda, Hyodo, DJ, et al. in preparation

# How strong $\pi\Sigma$ interaction ?

## More interesting question

$\pi \Sigma$ : resonance or virtual state ??

$\pi\Sigma$  interaction is attractive.

If attraction is unexpectedly enough strong, there could be a  **$\pi \Sigma$  virtual state** below threshold. In this case,  $\Lambda(1405)$  consists of single pole.

Energy-independent potentials also provide  $\pi\Sigma$  virtual states.

If virtual state exists

**$\pi\Sigma$  scattering length  $\sim 5$  fm**

for  $\pi\Sigma$  resonance case

Model calculation of $I=0$	BNW	ORB	HNJH	BMN
$\pi\Sigma$ scattering length (fm)	<b>0.517</b>	<b>0.789</b>	<b>0.692</b>	<b>0.770</b>

More systematic study will be public soon.

Ikeda, Hyodo, DJ, et al. in preparation

# Consequences of chiral unitary model

## 3. $\Lambda(1405)$ is a quasibound state of meson-baryon

a theoretical indication

the details are given in Hyodo' talk  
(Sep. 17, parallel session 2A)

### chiral unitary model

model parameters tuned so as to

a) reproduce scattering data

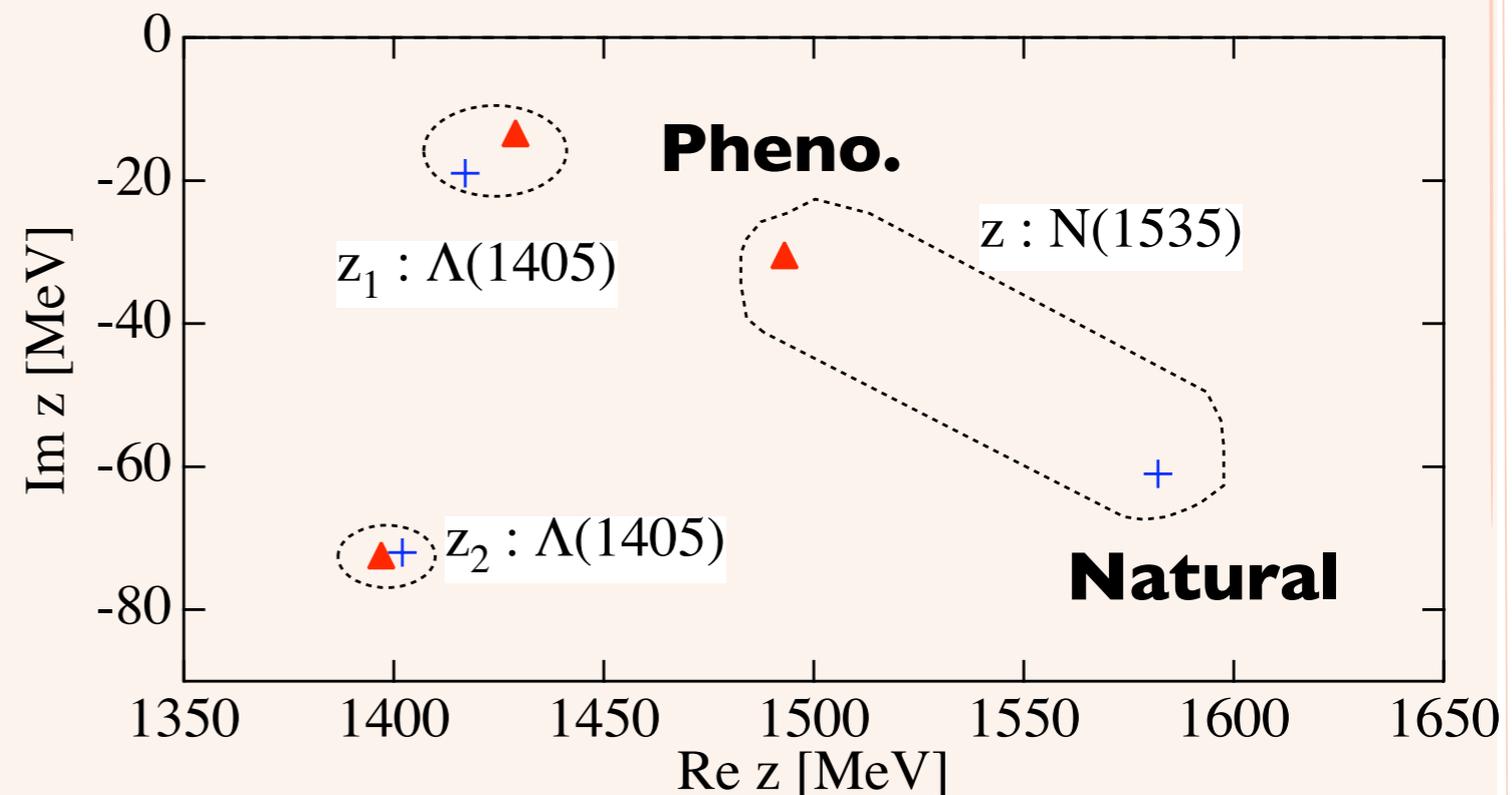
▲ **Pheno.**

b) exclude quark-originated  
states theoretically

+ **Natural**

**V : WT term**

**$\Lambda(1405)$  has mostly meson-baryon components.**



Hyodo, Jido, Hosaka, PRC78, 025203 ('08)

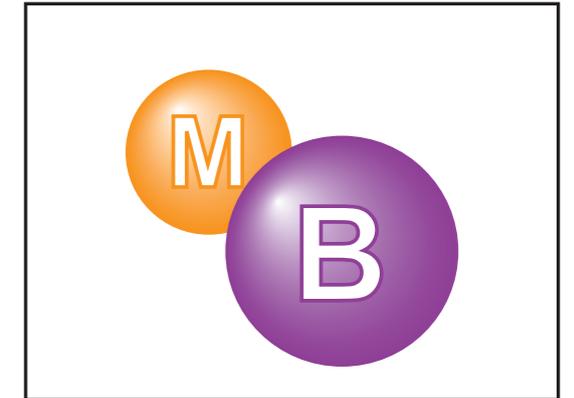
# $\Lambda(1405)$ as quasibound state of hadrons

## 1. large $N_c$ behavior

Hyodo, DJ, Roca, PRD77, 056010 (2008)  
Roca, Hyodo, DJ, NPA809,65, (2008)

different scaling of the width from quark model

**T. Hyodo, poster M19, today**



## 2. Electromagnetic radii

Sekihara, Hyodo, DJ, PLB669, 133 (2008)

$\Lambda(1405)$ : quasibound state of  $K^{\text{bar}}N$  with 10~30 MeV

**spatially extended**

almost real Kaon surrounding nucleon  
larger radius than neutron charge radius

**negative charge radius**

$$\langle r^2 \rangle_E = -0.12 \text{ [fm}^2\text{]} \quad \text{virtual pion cloud}$$

$K^-$  spreads widely around proton

electromagnetic form factor of  $\Lambda(1405)$  in chiral unitary model

$$\langle r^2 \rangle_E = -0.13 + 0.30i \text{ [fm}^2\text{]}$$

$$\text{moduls } |\langle r^2 \rangle_E| = 0.33 \text{ [fm}^2\text{]}$$

complex number

$$\text{remove decay chan. } \langle r^2 \rangle_E = -0.52 \text{ [fm}^2\text{]}$$

# Potential model for $\Lambda(1405)$

$\Lambda(1405)$  : quasibound state with a small binding energy  $\sim 10\text{-}30$  MeV

$\Lambda(1405)$  can be described by single-channel potential model with  $\pi\Sigma$  decay channel in the imaginary part

$\pi\Sigma$  coupled channel effect will be important, if binding energy is large.

## Hyodo-Weise potential (HW-HNJH)

derived from chiral dynamics

energy dependent, but small in energy of interest

resonance position  $\sim 1420$  MeV

PRC77,035204 (08)

## Akaishi-Yamazaki potential (AY)

obtained phenomenologically

$l=0$  : reproduce  $\Lambda(1405)$  as quasi-bound state of  $K^{\text{bar}}N$

mass: **1405 MeV**, width: 40 MeV

PRC64,044005 (02)

$K^{\text{bar}}NN$  one of the simplest nuclear system

Akaishi-Yamazaki

later various models applied to this system

present achievement in theory : bound with large width

we have controversy over the binding energy and width

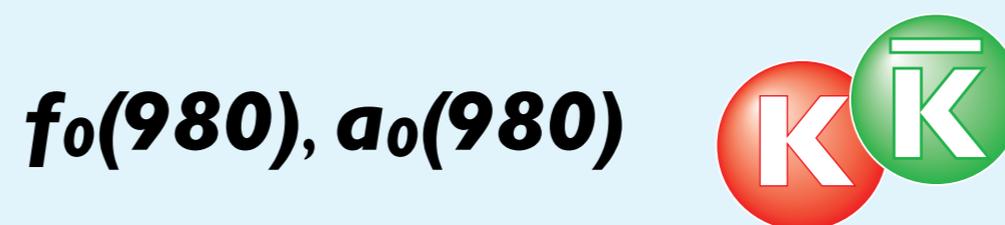
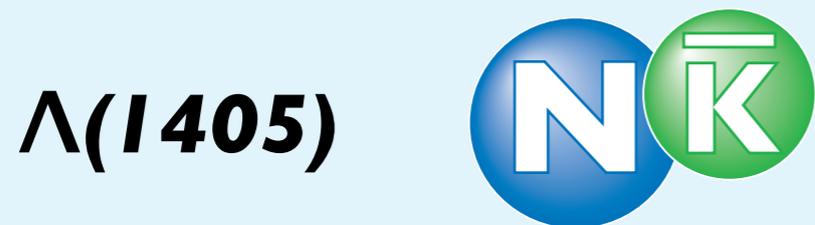
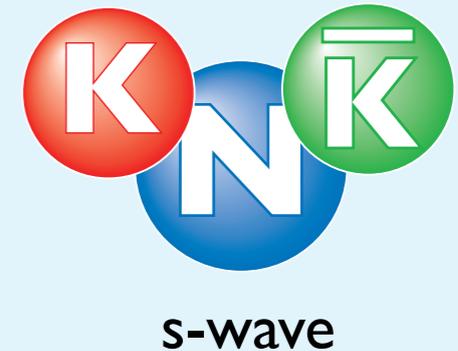
# $K \bar{K} N$ system with $I=1/2, J^P=1/2^+$

## assumption

non-relativistic potential model

$\Lambda(1405)$  is a quasi-bound state of  $K^{\text{bar}}N$

$f_0(980)$  and  $a_0(980)$  are quasi-bound states of  $KK^{\text{bar}}$



## Interactions in $KK^{\text{bar}}N$ system

	$I=0$	$I=1$	threshold
$\bar{K}N$ <span style="color: red;">attraction</span>	$\Lambda(1405)$	weak attraction	1434.6 MeV
$K\bar{K}$	$f_0(980)$	$a_0(980)$	991.4 MeV
$KN$ <span style="color: purple;">repulsion</span>	very weak	strong repulsion	1434.6 MeV

# Result $KK^{\text{bar}}N$

$N^*$  at 1910 MeV

$K\bar{K}N$  is bound below thresholds of  $\Lambda(1405)+K$ ,  $a_0(f_0)+N$

## - loosely bound system

B.E. from  $KK^{\text{bar}}N$

width

HW: **19 MeV**

**88 MeV**

AY: **39 MeV**

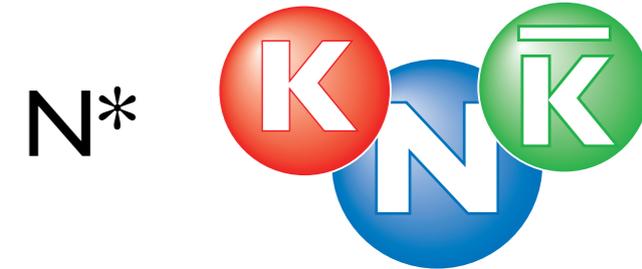
**98 MeV**

sum of those of isolated two-particle systems

DJ, Y. Kanada-En'yo, **PRC78, 035203 (2008)**

Faddeev calculation also obtains this resonance  
A.M. Torres's Talk, 15 Sep.

Martinez Torres, Khemchandani, Oset, **PRC79, 065207 (2009)**



## spatial structure

$\Lambda(1405)$

$a_0(980)$



2.8 fm  
(2.3)



2.1 fm  
(1.4)



2.3 fm  
(2.1)

r.m.s radius: **1.7 fm**

hadron-hadron distances are comparable with nucleon-nucleon distances in nuclei

**main decay modes**  $\pi\Sigma K$  from  $\Lambda(1405)$   
 $\pi\eta N$  from  $a_0(980)$

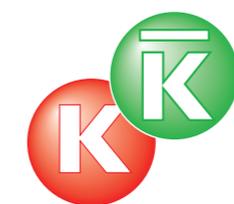
- **coexistence of two quasi-bound states keeping their characters**  $\Lambda(1405)+K$   
 $a_0(980)+N$

$\Lambda(1405)$

$a_0(980)$



HW: 1.9 fm  
AY: 1.4 fm



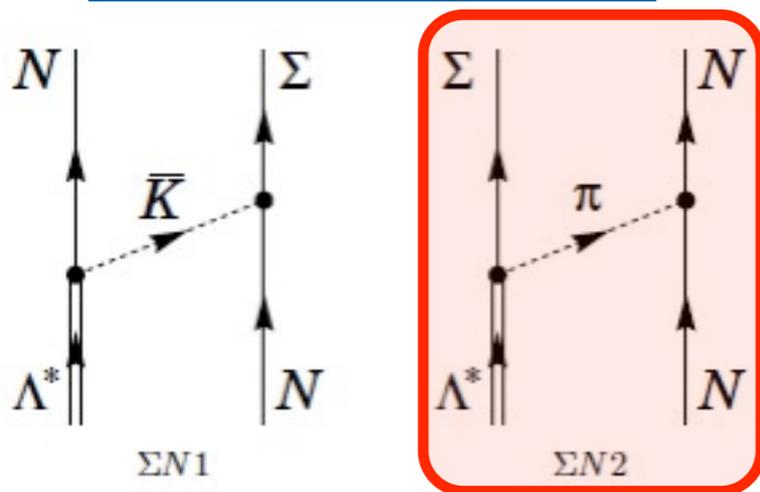
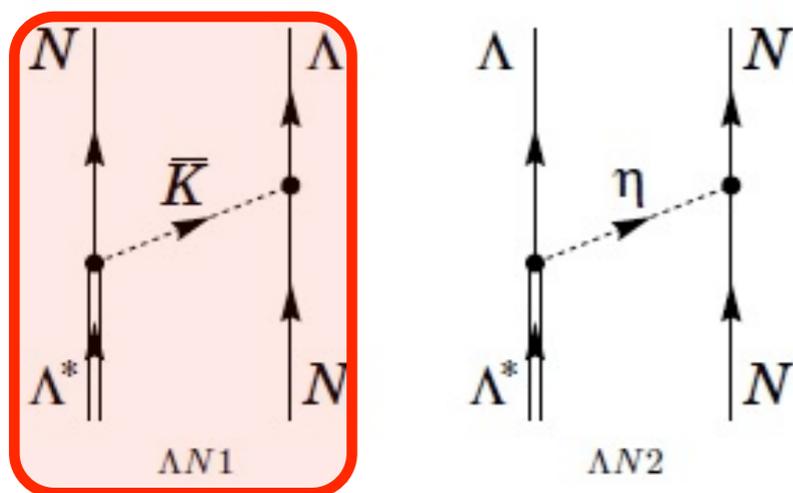
HW: 2.1 fm  
AY: 2.2 fm

# Nonmesonic decay of $\Lambda(1405)$ in nuclear matter

Kaonic nuclei = hadronic excitation of hypernuclei

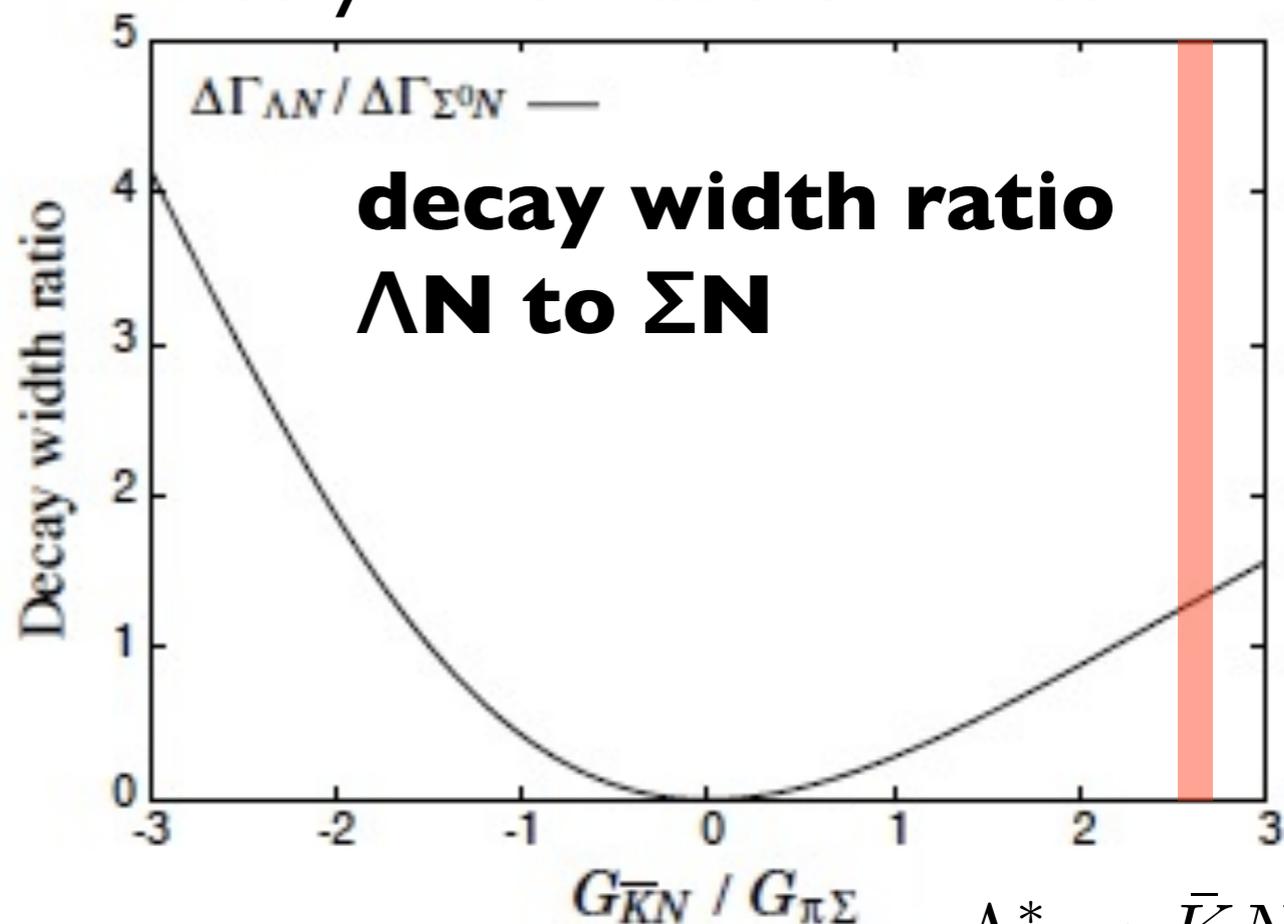
**$\Lambda^*$  can be doorway of kaon absorption**

Sekihara, DJ, Y. Kanada-En'yo,  
PRC79, 062201(R) (2009);  
Sekihara, parallel session 2-A

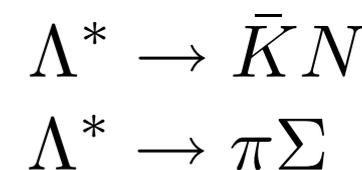


one-meson exchange model

Decay width ratio of  $\Lambda N$  to  $\Sigma N$



ratio of couplings



**total nonmesonic decay of  $\Lambda^*$ :  
20 MeV at  $\rho_0$**

# Summary

## **$\Lambda(1405)$ is the gift of meson-baryon dynamics**

- quasibound state of meson and baryon  
large spatial size
- double pole structure  
strong attraction in  $K^{\text{bar}}N$  and  $\pi \Sigma$  channels  
 $K^{\text{bar}}N$  bound state and  $\pi \Sigma$  resonance
- resonance position depends on channels

$K^- d \rightarrow \Lambda(1405)n$       any information of  $\pi \Sigma (l=0)$  scattering

## **$\Lambda(1405)$ in few-body systems**

another example of kaon bound system:  $K^{\text{bar}}KN$

a new  $N^*$  resonance  $N(1910)$       coexistence of  $\Lambda(1405)$ -K and  $a_0(980)$ -N

doorway state of K absorption

# Collaborators

**Sekihara** (Kyoto)

**Hyodo** (Tokyo Tech.)

**Kanada-En'yo** (YITP, Kyoto)

**Hosaka** (RCNP, Osaka)

**Ikada** (RIKEN & Tokyo)

**Roca** (Murcia)

**Oller** (Murcia)

**Oset** (Valencia)

**Ramos** (Barcelona)

**Meißner** (Bonn, Jülich)