Equation of state and magnetic monopoles in SU(2) gluon plasma

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## Outline

- Thermodynamics of Yang-Mills theory
- Models of color confinement at T < Tc
  - Abelian monopoles
- In deconfinement (gluon plasma at T > Tc)
  - Are they (still) alive as real object?
- Contribution to (trace of) energy-momentum tensor from Abelian monopoles

## Thermodynamics

• Free Energy (T is temperature and V is spatial volume)

 $F = -T \log Z(T, V)$ 

Pressure

$$p = \frac{T}{V} \frac{\partial \log Z(T, V)}{\partial \log V} = -\frac{F}{V} = \frac{T}{V} \log Z(T, V)$$

• Energy density  

$$\varepsilon = \frac{T}{V} \frac{\partial \log Z(T, V)}{\partial \log T}$$

• Entropy density 
$$s(T) = \frac{\varepsilon + p}{T} = \frac{\partial p(T)}{\partial T}$$

### **Thermodynamics: Trace Anomaly**

• Trace anomaly of the energy-momentum tensor  $T_{\mu\nu}$  $\theta(T) = \langle T^{\mu}_{\mu} \rangle \equiv \varepsilon - 3p = T^5 \frac{\partial}{\partial T} \frac{p(T)}{T^4}$  Trace anomaly is vanishing,  $\Box$  if excitations are massive,  $M \gg T$   $\varepsilon \sim p \sim \exp\{-M/T\}$  $\Box$  if particles are massless and non-interacting,  $\varepsilon = 3p$ Pressure via trace anomaly  $p(T) = T^4 \int^T \frac{dT_1}{T_1} \frac{\theta(T_1)}{T_1^4}$ Energy density via trace anomaly  $\varepsilon(T) = 3T^4 \int^T \frac{dT_1}{T_1} \frac{\theta(T_1)}{T_1^4} + \theta(T)$ Trace anomaly is a key quantity

### Trace Anomaly for SU(2) pure gluons

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#### Partition Function

$$Z(T,V) = \int DU \exp\{-\beta \sum_{P} S_{P}[U]\}, S_{P}[U] = (1 - \frac{1}{2}TrU_{P})$$

$$T = 1/(N_t a), \quad V = (N_s a)$$

Trace anomaly on the lattice

$$\frac{\theta(T)}{T^4} = 6 N_t^4 \left( \frac{\partial \beta(a)}{\partial \log a} \right) \cdot \left( \langle S_P \rangle_T - \langle S_P \rangle_0 \right)$$

## Mechanisms of color confinement

#### Dual superconductor picture

['t Hooft, Mandelstam, Nambu, '74-'76]

- Based on existence of special gluonic configurations, called ``magnetic monopoles''
- □ Monopoles are classified with respect to the Cartan subgroup  $[U(1)]^{N-1}$  of the SU(N) gauge group

Confinement is due to monopole condensation

- Monopole dominance for various quantities
  - □ String tension
  - Polyakov loop behaviors
  - Critical exponents

# Confinement (T<Tc) and plasma (T>Tc)

- The monopoles are percolating and condensed in confining vacuum
- The percolating monopole cluster disappears and monopole condensate vanishes in deconfinement phase
- Must emerge as a real (thermal) component of deconfinement plasma

[V.I.Zakharov, M.N.Chernodub,'07] [Liao and Shuryak, '06-'07]

- Similar to electrically neutral electron-positron plasma:
   individual particles exist at high temperatures in a heat bath
   annihilate at low temperatures, but still present in the vacuum
- <u>Check</u>: if the suggestion is true, then the monopoles must contribute to the equation of state of the gluon plasma

Gauge fixing (MA gauge): maximize

$$R = \sum_{s,\mu} Tr[\sigma_3 U_{\mu}(s)\sigma_3 U_{\mu}^{\dagger}(s)]$$

[A.S.Kronfeld, M.L.Laursen, G.Schierholz, U.J.Wiese '87]

 Define particular singular gluon objects (monopoles) [T.A.DeGrand, D.Toussaint '80]

 $k_{\mu}(s) = \epsilon_{\mu\nu\rho\sigma}\partial_{\nu}n_{\rho\sigma}(s+\hat{\mu})/2$ 

- Extract the plaquettes around the monopole
- Decompose the trace anomaly into two parts
   The contribution around the monopole and the rest [for center vortex ; M.N.Chernodub, A.Nakamura, V.I.Zakharov '08]

Action density:

$$\begin{split} \langle S_P \rangle &= \langle S_P \rangle^{\text{mon}} + \langle S_P \rangle^{\text{rest}} \\ &= \frac{1}{6N_s^3 N_t} \left[ \langle \sum_P \rho_P S_P \rangle + \langle \sum_P (1 - \rho_P) S_P \rangle \right] \\ \rho_P &= 1 \ (P \in \Sigma) \text{ or } 0 \ (P \notin \Sigma) \\ \Sigma \text{ : plaquettes around monopoles} \\ \end{split}$$
Trace anomaly (naive regularization):
$$\begin{aligned} \theta_{\text{naive}}^{\text{mon}} \\ \frac{\theta_{\text{naive}}^{\text{mon}}}{T^4} &= 6N_t^4 \left( \frac{\partial \beta}{\partial \log a} \right) \left[ \langle S_P \rangle_T^{\text{mon}} - \rho(T) \langle S_P \rangle_0 \right] \end{aligned}$$

 $\frac{\theta_{\text{naive}}^{\text{rest}}}{T^4} = 6N_t^4 \left(\frac{\partial\beta}{\partial\log a}\right) \left[\langle S_P \rangle_T^{\text{rest}} - (1 - \rho(T)) \langle S_P \rangle_0\right]$ 

Trace anomaly (naive regularization):



T = 0 :<br/>16<sup>4</sup> lattice,<br/>1000 conf.

T > 0: 16<sup>3</sup> × 4 lattice, 5000 conf.

Calculated by RICC at RIKEN

Not sensitive to the phase transition ⇒ regularization not appropriate

#### Specific action density:



(action density per an elementary plaquette)

Regularized trace anomaly :

$$\frac{\theta_{\text{reg}}^{\text{mon}}}{T^4} = 6N_t^4 \left(\frac{\partial\beta}{\partial\log a}\right) \rho(T) \left[\langle s_P \rangle_T^{\text{mon}} - \langle s_P \rangle_0^{\text{mon}}\right]$$

#### Specific action density :



The difference between  $\langle s_P \rangle^{
m mon}$  and  $\langle s_P \rangle^{
m rest}$  is seen clearly.



•Sensitive to the phase transition.

•The behavior is similar to the case of center vortex. [M.N.Chernodub et.al. '08]

### Conclusion and future works

#### Conclusion

Found: strong contributions from the plaquettes around Abelian monopoles to the trace anomaly, and, consequently, to the pressure and to the energy density of the gluon plasma.

□ Gluonic configurations around the Abelian monopoles are similar to the worldsheets of the center vortex.

#### Future works

- Check of scaling for trace anomaly (wrapped monopole (T>0) and the largest monopole cluster (T=0))
- □ What is the correct regularization scheme?