

The order of the deconfinement phase transition in a heavy quark region - Dependence on the number of flavors -



WHOT-QCD Collaboration:

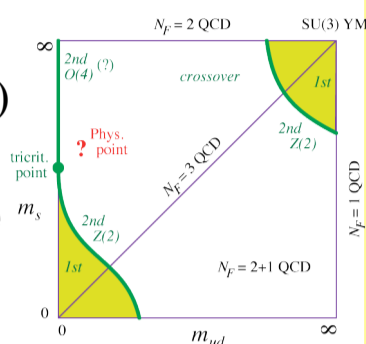
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Phase diagram

The order of the deconfinement phase transition is a fundamental topic. The order of the transition depends on the quark mass.

We investigate the end point of 1st order transition in the (m_{ud}, m_s) plane of $N_f=2+1$ by using lattice QCD simulation. In this study, we discuss it in a heavy quark mass region.



$N_f=2+1$ phase diagram

The N_f -dependence appears in the quark determinant:

$$R(P', \kappa) = \frac{\int \mathcal{D}U \delta(P(U) - P') [\{\det M(U, \kappa_{ud})\}^2 \det M(U, \kappa_s)]}{\int \mathcal{D}U \delta(P(U) - P')}$$

We can consider the many flavors case by replacing into $[\det M(U, \kappa)]^{N_f}$.

Lattice setup

- Action : Standard wilson fermion (hopping parameter $\kappa \approx m_q^{-1}$)

and Plaquette action

- HeatBath / Simulation Setup
- The $\det M(\kappa)$ is evaluated by using a hopping parameter expansion.

The LO of the quark determinant factor has a coefficient $N_f \kappa_{ep}$. The shape of the potential is unchanged when the coefficient is constant.

β	Confs
5.68	100,000
5.685	430,000
5.69	500,000
5.6925	670,000
5.7	100,000

Probability distribution function

Probability distribution function is useful to identify the order of a phase transition.

❖ Probability distribution function:

$$w(P') = \int \mathcal{D}U \mathcal{D}\psi \mathcal{D}\bar{\psi} \delta(P(U) - P') e^{-S(\beta, \kappa)}$$

β : coupling

❖ An effective potential:

$$V(P) = -\ln w(P)$$

❖ Nature of the potential

- The shape of the potential \Leftrightarrow order of the transition

double-well : 1st order
single well : 2nd/crossover

- The derivative of the potential

- Turning point in the two shapes

- \Leftrightarrow End point of 1st order

❖ Reweighting method

- (i) **β -Direction** : To get dV/dP in a wide range

$$\frac{dV(P, \beta)}{dP} = \frac{dV(P, \beta')}{dP} - 6N_{\text{site}}(\beta - \beta')$$

- (ii) **κ -Direction** : To evaluate the end point

$$V(P, \kappa, \beta) = -\ln R(P, \kappa) + V_0(P, \beta)$$

$$\text{Reweighting term : } R(P', \kappa) = \frac{\int \mathcal{D}U \delta(P(U) - P') [\det M(U, \kappa)]}{\int \mathcal{D}U \delta(P(U) - P')}$$

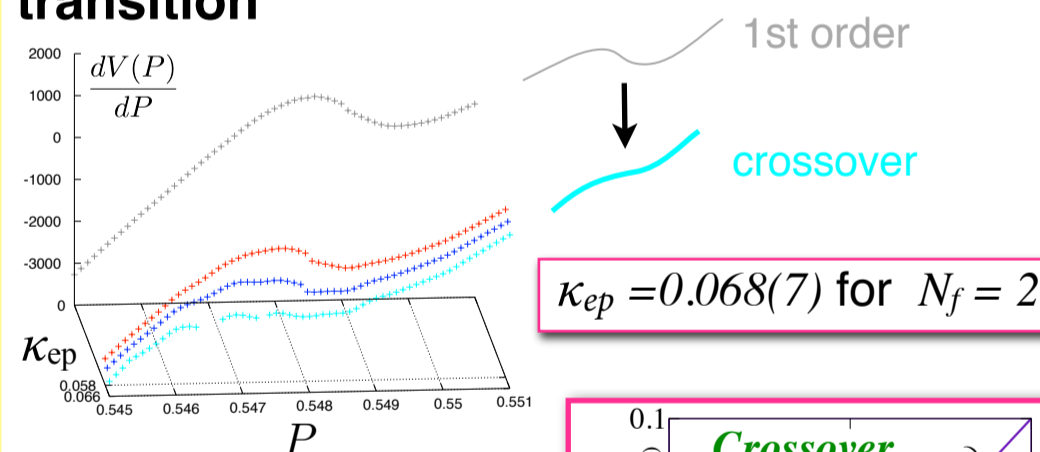
The change of the order is caused by dynamical quark effect which is included by the quark determinant factor.

$$V_0(P', \beta) = -\ln \left[\int \mathcal{D}U \delta(P(U) - P') e^{-S_g(\beta)} \right]$$

HS talk in Lattice2010

Result

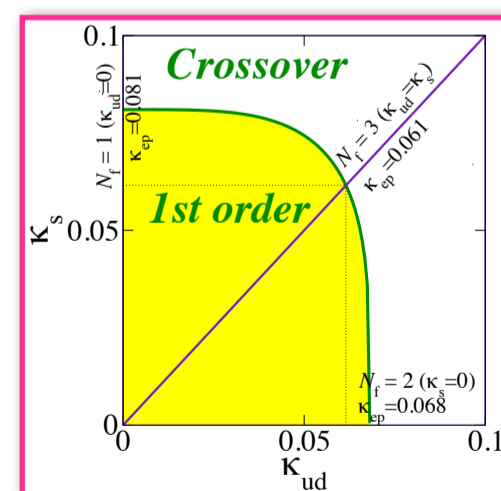
❖ The order of the deconfinement phase transition



$$\kappa_{ep} = 0.068(7) \text{ for } N_f = 2$$

❖ The $N_f=2+1$ phase diagram

The result is consistent with the result of Z(3) model reported by Alexandrou *et al*(1990).



Summary & Future prospect

- The κ_{ep} value decreases as N_f increases.
- The method can be applied to the case of non-zero μ . S. Ejiri's talk
- The method of probability distribution function to identify the order of the transition works well.