# Quarkonium mass splittings in three-flavor lattice QCD 

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## Outline

1. Objectives
2. Challenges
3. Results
4. Conclusions

## Charmonium spectroscopy before the B-factories



## Charmonium spectroscopy after the B-factories



## The message

- We want to guide the discovery and classification of excited quarkonium states.
- There has been significant recent progress in methods and results for "gold-plated" levels (i.e. ground states).
- These successes point the way to improvements in all quantities.


## Lattice challenges for heavy quarks

- Reducing lattice cutoff $1 / a$ effects, especially for heavy quarks. $\mathcal{O}(M a)$ errors are bad for charm and bottom when, typically $1 / a \approx 1.8-3 \mathrm{GeV}$.
- Nonrelativistic QCD: expansion in $p / M$ slow for charm, good for bottom.
- Fermilab quarks good for both. Further improvements under study [Oktay, Kronfeld]
- Highly Improved Staggered Quarks (HISQ) errors first at $\mathcal{O}\left(\alpha_{s}^{2}\left(a M_{c}\right)^{2}\right)$ : good for charm, not so good for bottom with today's lattices.
- Excited states. Ground state properties are easiest. Excited states more difficult.
- Multihadronic states, e.g., open charm are complicated. We are just beginning to learn how to treat them.


## Quarkonium spectroscopy

- Hadron masses are determined from propagators

$$
C_{i j}(t)=\langle 0| \mathcal{O}_{i}(t) \mathcal{O}_{j}(0)|0\rangle \rightarrow z_{i}^{*} z_{j} \exp (-m t) \text { for large } t,
$$

where $\mathcal{O}_{i}$ 's are suitable interpolating operators. For example, for the $J / \psi$ we could use $\mathcal{O}=\bar{q} \gamma_{\mu} q$.

- Operators are classified according to lattice symmetries.
- The cubic group replaces the rotation group. $A_{1}$ replaces $J=0, T_{1}$ replaces $J=1$, etc.
- Ambiguities: $A_{1}$ sees $J=0,4,6 ; T_{1}$ could be $1,3,4$, etc. With a little effort we can often resolve them.
- If we use a large basis set $\mathcal{O}_{j}$ for the same quantum numbers, we get a correlation matrix. The eigenvalues contain information about the ground and excited states. Several groups now use this method successfully.
- In the next slides I show a sample of results obtained by Dudek et al. (2007) using this method.


## JLab quarkonium: $J^{--}$states


[Dudek, Edwards, Mathur, and Richards, arXiv:0707.4162]

## JLab quarkonium: $J^{--}$states



- Comparison with experimental masses and quark potential model masses. Tabulated masses in MeV .
[Dudek, Edwards, Mathur, and Richards, op cit.]


## JLab quarkonium $J^{--}$states

- How well does this work? Look at the $T_{1}^{--}$channel, where the ground state $J / \psi$ has a very clean signal.

| state | predict | PDG | difference |
| :--- | :--- | :--- | :--- |
| $J / \psi$ | $3109(2)$ | 3097 | 12 |
| $\psi(2 S)$ | $3722(24)$ | 3686 | 36 |
| $\psi(3 S)$ | $3855(12)$ | 3773 | 82 |
| $\psi(4040)$ | $3843(18) ?$ | 4039 | -196 |
| $\psi(4160)$ | $4472(79) ?$ | 4153 | 319 |
| $\psi(4415)$ | $4442(48) ?$ | 4421 | 21 |

- Here ? means the authors did not make any assignment.
- Higher excitations: more difficult to assign J.
- Errors grow with excitation, as expected.
- When the ground state signal is not so clean, less can be extracted.


## Hadron Spectrum Collaboration quarkonium future

These results are pioneering and impressive, but ...

- No sea quarks (quenched approximation).
- Only one lattice spacing, so no continuum limit.
- No open charm states (none, anyway, in quenched approximation.)

The collaboration is currently remedying these shortcomings.
[Ryan, Lattice 2010]

## Open charm mixing?

Causes level shifts. Is it significant?

- For static quarks, string breaking studies on the lattice suggest mixing is weak.
- For dynamical quarks, very little is known from lattice studies.
- Bali and Ehmann (Lattice 2009) studied mixing between S-wave charmonium and a $D \bar{D}$ "molecule" using a variational method.

| state | $(c \bar{c})_{l}$ | $(c \bar{c})_{n}$ | $(c \bar{c} \bar{c} u)_{l}$ | $(c \bar{u} \bar{c} u)_{n}$ |
| :---: | :---: | :---: | :---: | :---: |
| $\eta_{c}$ | $0.54(3)$ | $-0.02(1)$ | $-0.1(1)$ | $-0.31(5)$ |
| $D_{1} D^{*}$ | $0.07(1)$ | $0.01(1)$ | $-0.46(8)$ | $0.14(2)$ |
| $J / \psi$ | $0.51(4)$ | $-0.03(1)$ | $0.09(1)$ | $0.21(6)$ |
| $D_{1} \bar{D}$ | $0.08(6)$ | $0.04(1)$ | $-0.18(1)$ | $0.53(4)$ |
| $\chi_{c 1}$ | $0.39(5)$ | $0.69(3)$ | $-0.22(3)$ | $-0.49(4)$ |
| $D D^{*}$ | $0.63(4)$ | $-0.23(3)$ | $-0.73(4)$ | $0.12(3)$ |

- (" 1 " and " $n$ " refer to different basis wave functions.)
- In some cases such as the $\chi_{c 1}$ the mixing appears to be large, but a more thorough study is now needed.
[Bali, Ehmann, arXiv:0911.1238]


## MILC Collaboration gauge field ensembles

- Ensembles to help us reach the physical point and continuum.
- Parameters of a publicly available archive of gauge configurations based on $u, d$, and $s$ sea quarks.

| ensemble | $a$ (approx) $(\mathrm{fm})$ | sea quark ratio $m_{u d} / m_{s}$ |
| :--- | :--- | :--- |
| Extra coarse | 0.18 | $0.6,0.4,0.2,0.1$ |
| Medium coarse | 0.15 | $0.6,0.4,0.2,0.1$ |
| Coarse | 0.12 | $0.6,0.4,0.2,0.15,0.1$ |
| Fine | 0.09 | $0.4,0.2,0.1,0.05$ |
| Superfine | 0.06 | $0.4,0.2,0.1$ |
| Ultrafine | 0.045 | 0.2 |

- With these one can carry out an extrapolation to physical $m_{u d}$ ( $\approx 0.037 m_{s}$ ) and $a=0$ (continuum).
[MILC, Rev Mod Phys 82, 1349 (2010)]


## Progress in lattice actions and analysis campaigns

Both campaigns use MILC lattices to extrapolate to the physical point.

| Study | Charm | Bottom | Ops | States | Comment |
| :--- | :--- | :--- | :--- | :--- | :--- |
| OId FNAL/MILC | FNAL | FNAL | 2 | low S, P | $a \geq 0.09 \mathrm{fm}$ |
| New FNAL/MILC | FNAL | - | JLab | many | $a \geq 0.06 \mathrm{fm}$ |
| HPQCD | HISQ | NRQCD | 2 | low S, P | $a \geq 0.06 \mathrm{fm}$ |
| Future FNAL/MILC | HISQ | FNAL | JLab | many | $a \geq 0.06 \mathrm{fm}$ |

## FNAL/MILC Overview

- Charmonium levels constructed from splittings from the spin-averaged $\overline{1 S}$ level.


- Lattice spacings are $0.18 \mathrm{fm}, 0.15 \mathrm{fm}, 0.12 \mathrm{fm}, 0.09 \mathrm{fm}$.
[MILC/FNAL, PRD 81, 034508 (2010)]


## HPQCD Overview

- Gold-plated meson spectrum based on a subset of the MILC ensembles.

The gold-plated meson spectrum from lattice QCD - HPQCD 2008


- Five states are used to get the four quark masses and lattice spacing.
- Three states were predictions.
[HPQCD-PoS LATTICE2008, 118]


## HPQCD $D_{s}$

Example of achievable precision: $D_{s}: 3 \mathrm{MeV}$ !


- Based on splitting $M(c \bar{s})-\frac{1}{2} M\left(\eta_{c}\right)$
- PACS-CS Lattice 2009 result using their relativistic heavy quark action: $1.972(2) \mathrm{GeV}$ vs expt 1.968. (Error is statistical only.)
[HPQCD, arXiv:1008.4018; PACS-CS arXiv:0911.5362]


## HPQCD $B_{c}$

Example of achievable precision: $B_{c}$ : 10 MeV !


- Based on splitting $M\left(B_{c}\right)-M(b \bar{b}) / 2-M\left(\eta_{c}\right) / 2$.
- Light shaded band includes all errors.
[HPQCD, arXiv:1010.3848]


## MILC/FNAL 1S hyperfine splitting



- Charm result 117(11) MeV.
- Bottom result 53(9) MeV.
- Annihilation effects are ignored here.
[FNAL/MILC PRD 81, 034508 (2010)]


## HPQCD 1S hyperfine splitting

HPQCD PRELIMINARY


- Annihilation effects are estimated from perturbation theory. They and other corrections shift the result up to the pink band, which indicates all errors: 2 MeV
[HPQCD, private communication, 2010]


## Annihilation contribution to charm HFS

Calculated from the lattice rather than perturbation theory.


Connected and disconnected diagrams

- Decreases the HFS splitting - by about 2 MeV .
- Sign contrary to perturbation theory (partly due to axial anomaly).
[Levkova and DeTar, Lattice 2010 and forthcoming paper.]


## MILC/FNAL $\overline{\mathbf{2 S}}-\overline{\mathbf{1 S}}$ splitting

Where we don't do so well.


Charmonium


Bottomonium

- The calculation does not treat the open charm threshold.
- Note the open bottom threshold is safely off scale here.
- Is the disgreement in charmonium caused by open charm?
[FNAL/MILC PRD 81, 034508 (2010)]


## MILC/FNAL $D_{s}^{*}-D_{s}$



- Based on our "old" data.
- Splittings are in $r_{1}=0.311 \mathrm{fm}$ units here. $\left(1 / r_{1}=635 \mathrm{MeV}\right)$.
- The green point includes discretization errors. Accuracy: 4 MeV .
- PACS-CS Lattice 2009 result with their relativistic heavy quark action: $135(3) \mathrm{MeV}$ vs expt 144. (Statistical only.)
[FNAL/MILC arXiv:1003.1937, PACS-CS arXiv:0911.5362]


## Fermilab-Lattice/MILC quarkonium preview



## Toward precision lattice charmonium results

Through tackling "gold-plated" quantities, we have learned what is required to do good lattice charm physics for all quantities:

- A heavy quark action with small discretization errors and an accurate tuning of the heavy quark masses.
- An accurate determination of the lattice scale.
- A full treatment of sea quarks. Simulate at the physical light quark masses or do a controlled extrapolation.
- A careful extrapolation to zero lattice spacing.
- Good interpolating operators and an adequate data sample.


## Conclusions

- Improved lattice charm quark formulations yield high precision for gold-plated quantities.
- This experience is teaching us how to do good charm physics on the lattice.
- Gauge field ensembles at smaller lattice spacing enable significant reductions in errors for all quantities.
- Expect improvements in excited state and exotic predictions from the Hadron Spectrum and MILC/Fermilab Lattice collaborations.
- Treating two-hadron mixing (e.g. open charm) remains a challenge.


Purple Mangosteen＝die Mangostane＝マンゴスチン

The Purple Mangosteen（Garcinia mangostana），colloquially known simply as＂the mangosteen＂，is a tropical evergreen tree believed to have originated in the Sunda Islands and the Moluccas of Indonesia．Botanically an aril，the fragrant edible flesh can be described as sweet and tangy，citrusy with peach flavor and texture．（Wikipedia）

