

Form factor calculations for mesons and baryons

Hartmut Wittig
Institut für Kernphysik



In collaboration with:

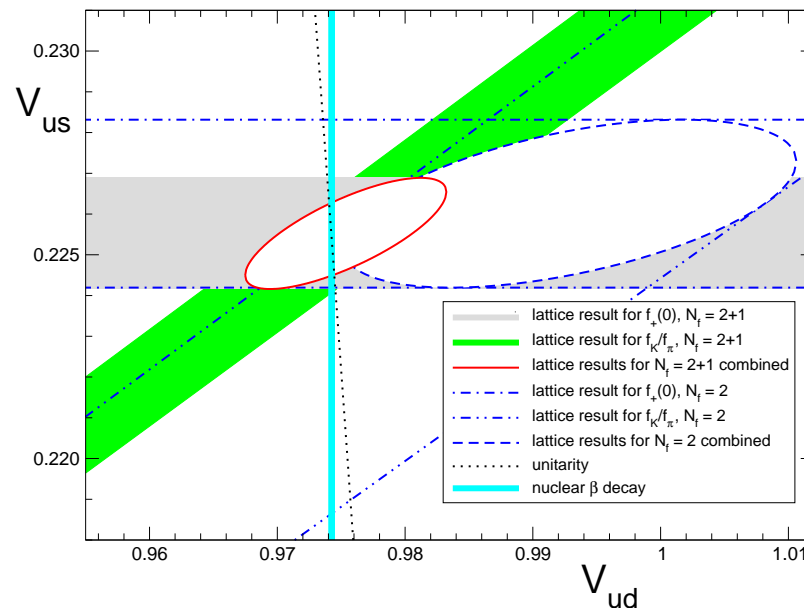
B. Brandt, S. Capitani, M. Della Morte, D. Djukanovic, G. von Hippel,
A. Jüttner, B. Knippschild, H.B. Meyer

Motivation & Outline

Form factors:

- provide information on **hadron structure**:
 - distribution of electric charge and magnetisation; **charge radii**
- accurate experimental data available
- relatively simple to compute on the lattice:
 - precise lattice estimates for $K_{\ell 3}$ -decays

[FLAG Working Group of FLAVIANet]



Motivation & Outline

Form factors:

- provide information on **hadron structure**:
 - distribution of electric charge and magnetisation; **charge radii**
 - accurate experimental data available
 - relatively simple to compute on the lattice:
 - precise lattice estimates for $K_{\ell 3}$ -decays *[FLAG Working Group of FLAVIANet]*
 - Large systematic uncertainties remain for **baryonic** form factors
- ⇒ “Next-generation benchmark” for lattice QCD

Outline:

1. Lattice Set-up
2. Pion electromagnetic form factor
3. Form factors and axial charge of the nucleon
4. Summary & Outlook

1. Lattice Set-up

- Perform a systematic study of mesonic and baryonic form factors with **controlled uncertainties**:
 - lattice artefacts
 - finite-volume effects
 - chiral extrapolations
 - excited state contamination
- Determine form factors with **fine momentum resolution**
- Eventually: include quark disconnected diagrams

1. Lattice Set-up

- Perform a systematic study of mesonic and baryonic form factors with **controlled uncertainties**:
 - lattice artefacts
 - finite-volume effects
 - chiral extrapolations
 - excited state contamination
- Determine form factors with **fine momentum resolution**
- Eventually: include quark disconnected diagrams
- **C**oordinated **L**attice **S**imulations: [\[https://twiki.cern.ch/twiki/bin/view/CLS/WebHome\]](https://twiki.cern.ch/twiki/bin/view/CLS/WebHome)
Berlin – CERN – Madrid – Mainz – Milan – Rome – Valencia – Wuppertal – Zeuthen
- Share configurations and technology

CLS run tables

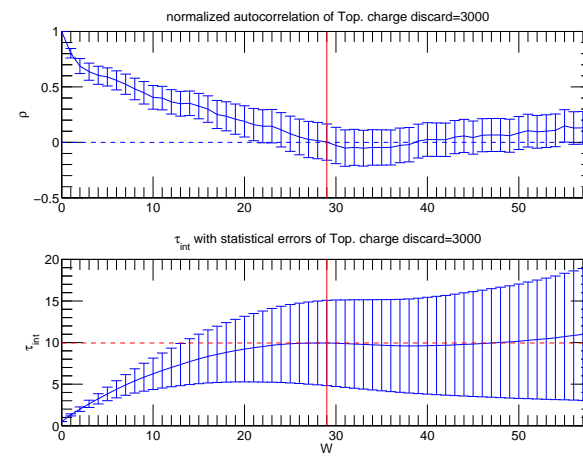
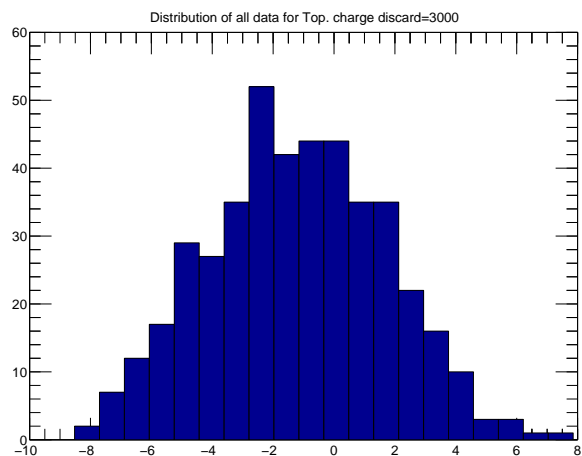
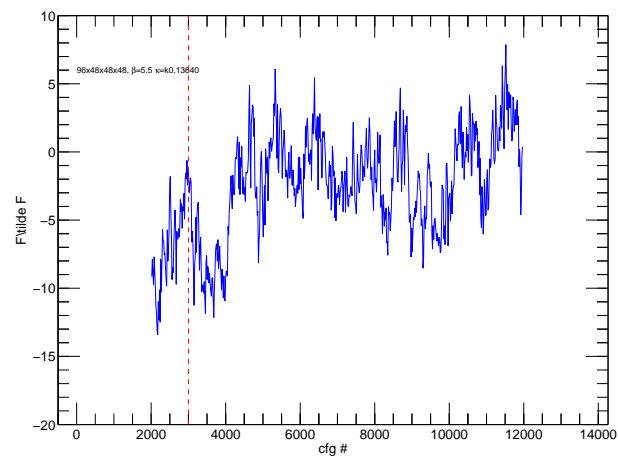
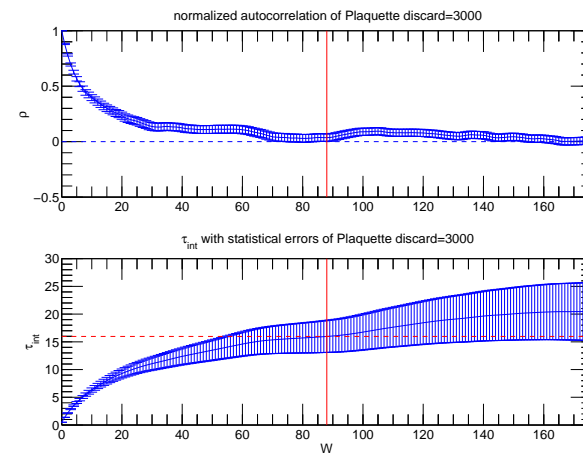
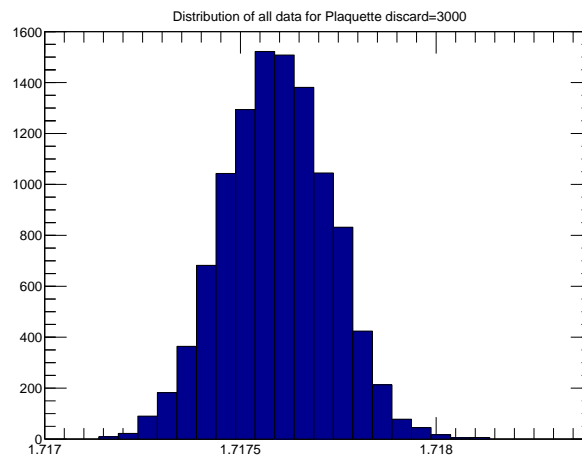
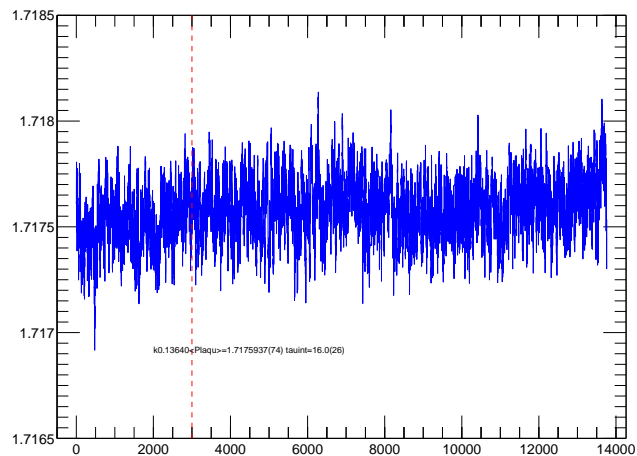
- $N_f = 2$ flavours of non-perturbatively $O(a)$ improved Wilson quarks
- Use deflation accelerated DD-HMC algorithm *[Lüscher 2003–07]*
- Generated ensembles without serious topology problems:

β	a [fm]	lattice	L [fm]	masses	$m_\pi L$	Labels
5.20	0.08	64×32^3	2.6	4 masses	4.8 – 9.0	A1 – A5
5.20	0.08	96×48^3	3.8	1 mass	—	B6
5.30	0.07	48×24^3	1.7	3 masses	4.6 – 7.9	D1 – D3
5.30	0.07	64×32^3	2.2	3 masses	4.7 – 7.9	E3 – E5
5.30	0.07	96×48^3	3.4	2 masses	5.0, 4.2	F6, F7
5.50	0.05	96×48^3	2.5	3 masses	5.3 – 7.7	N3 – N5
5.50	0.05	128×64^3	3.4	1 mass	4.7	O6

[Capitani et al., arXiv:0910.5578, Brandt et al., arXiv:1010.2390]

Autocorrelation times @ $\beta = 5.5$

N3: $\kappa = 0.13640$



The “Wilson” cluster at Mainz

- 280 nodes: 2 AMD “Barcelona” processors @ 2.3 GHz: 2240 cores
- Infiniband network & switch
- Peak speed: ~ 20 TFlop/s
- Sustained speed: 3.7 TFlops/s \Rightarrow 0.30 €/MFlops/s
- Waste heat: 20 kW/TFlops/s (Water-cooled server racks)

The “Wilson” cluster at Mainz

- 280 nodes: 2 AMD “Barcelona” processors @ 2.3 GHz: **2240 cores**
- Infiniband network & switch
- Peak speed: ~ 20 TFlop/s
- Sustained speed: **3.7 TFlops/s** \Rightarrow **0.30 €/MFlops/s**
- Waste heat: **20 kW/TFlops/s** (Water-cooled server racks)

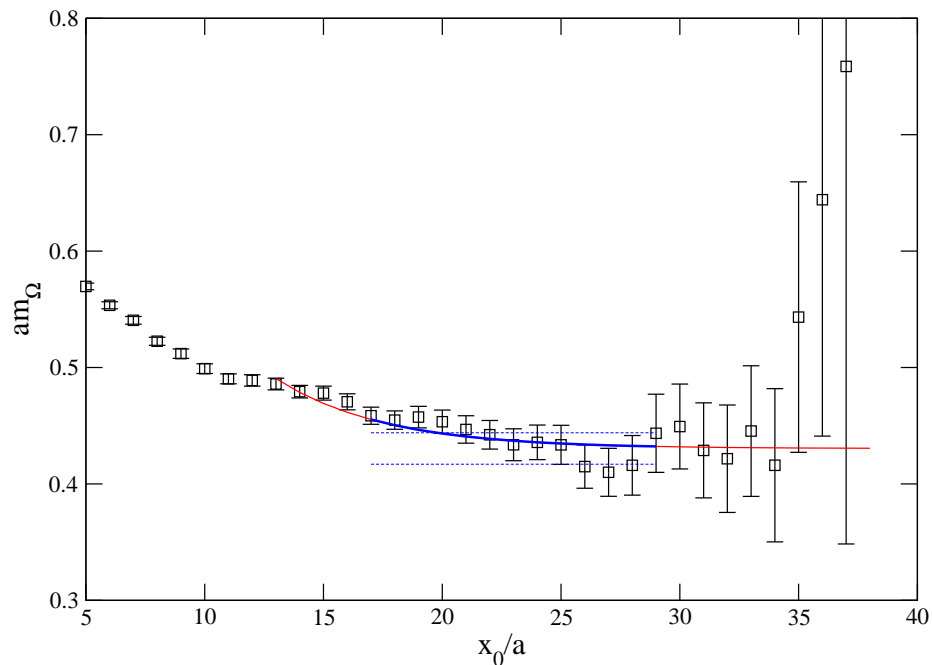


S

Scale setting

- Use mass of Ω -baryon to set the scale
- Effective masses from smeared-local correlator (Jacobi smearing)

$$\text{N5 : } \kappa_l = 0.1366, \quad \kappa_{\text{val}} = 0.1363, \quad \kappa_s = 0.13629(2)$$



- Fit to **ground** plus **1st excited** state: $m_1 \equiv m_\Omega$, $m_2 \equiv m_\Omega + 2m_\pi$

Scale setting

- Observe weak sea quark mass dependence of m_Ω

$$\rightarrow \beta = 5.5 : \quad a_\Omega = 0.053(1) \text{ fm} \quad (\text{preliminary})$$

- Determination of am_Ω still on-going at $\beta = 5.3$

$$\beta = 5.3 : \quad "a_\Omega" = \left(\frac{a_{\text{ref}}(\beta = 5.3)}{a_{\text{ref}}(\beta = 5.5)} \right) 0.053(1) \text{ fm} = 0.069(2) \text{ fm} \quad (\text{preliminary})$$

- To come: comparison with r_0, f_K, t_0

4. The pion form factor

- Provides information on pion structure:

$$\langle \pi^+(\vec{p}_f) | \frac{2}{3} \bar{u} \gamma_\mu u - \frac{1}{3} \bar{d} \gamma_\mu d | \pi^+(\vec{p}_i) \rangle = (p_f + p_i)_\mu f_\pi(q^2)$$

$$q^2 = (p_f - p_i)^2 : \quad \text{momentum transfer}$$

- Pion charge radius derived from form factor at zero q^2 :

$$f_\pi(q^2) = 1 - \frac{1}{6} \langle r^2 \rangle q^2 + O(q^4) \quad \Rightarrow \quad \langle r^2 \rangle = 6 \left. \frac{df_\pi(q^2)}{dq^2} \right|_{q^2=0}$$

- Restriction on minimum accessible momentum transfer:

$$\vec{p}_{i,f} = \vec{n} \frac{2\pi}{L} \quad \Rightarrow \quad |q^2| \geq 2m_\pi \left(m_\pi - \sqrt{m_\pi^2 + (2\pi/L)^2} \right)$$

$$L = 2.5 \text{ fm}, \quad m_\pi = 300 \text{ MeV} \quad \Rightarrow \quad |q^2| \geq 0.17 \text{ GeV}^2 = (0.41 \text{ GeV})^2$$

→ Lack of accurate data points near $q^2 = 0$

Twisted boundary conditions

[Bedaque 2004; de Divitiis, Petronzio & Tantalò 2004; Flynn, Jüttner & Sachrajda 2005]

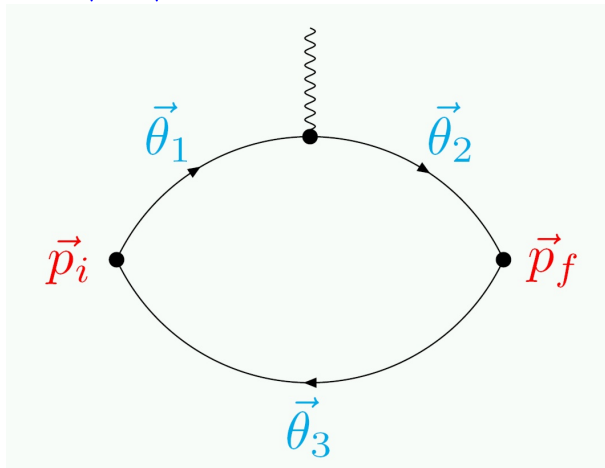
- Apply “twisted” spatial boundary conditions;

Impose periodicity up to a phase $\vec{\theta}$:

$$\psi(x + L\hat{e}_k) = e^{i\theta_k}\psi(x) \quad \Rightarrow \quad p_k = n_k \frac{2\pi}{L} + \frac{\theta_k}{L}, \quad k = 1, 2, 3$$

- Can tune $|q^2|$ to any desired value

[Boyle, Flynn, Jüttner, Sachrajda, Zanotti, hep-lat/0703005]



$$\begin{aligned} \vec{\theta}_i &= \vec{\theta}_1 - \vec{\theta}_3, \\ \vec{\theta}_f &= \vec{\theta}_2 - \vec{\theta}_3 \end{aligned}$$

$$\Rightarrow q^2 = (p_i - p_f)^2 = \left(E_\pi(\vec{p}_i) - E_\pi(\vec{p}_f) \right)^2 - \left[\left(\vec{p}_i + \frac{\vec{\theta}_i}{L} \right) - \left(\vec{p}_f + \frac{\vec{\theta}_f}{L} \right) \right]^2$$

Preliminary results

[Brandt, Della Morte, Djukanovic, Endreß, Gegelia, Jüttner, H.W.]

β	$L^3 \cdot T$	a [fm]	L [fm]	m_π [MeV]	Lm_π
5.50	$48^3 \cdot 96$	0.053	2.5	600	7.7
5.50	$48^3 \cdot 96$	0.053	2.5	510	6.5
5.50	$48^3 \cdot 96$	0.053	2.5	410	5.3
5.30	$48^3 \cdot 96$	0.069	3.3	290	5.0

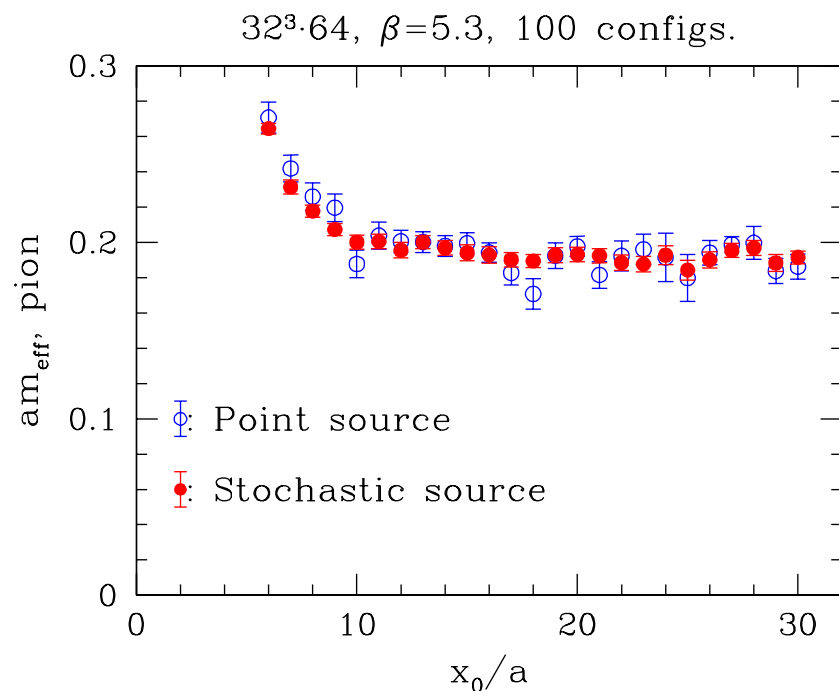
Preliminary results

[Brandt, Della Morte, Djukanovic, Endreß, Gegelia, Jüttner, H.W.]

β	$L^3 \cdot T$	a [fm]	L [fm]	m_π [MeV]	Lm_π
5.50	$48^3 \cdot 96$	0.053	2.5	600	7.7
5.50	$48^3 \cdot 96$	0.053	2.5	510	6.5
5.50	$48^3 \cdot 96$	0.053	2.5	410	5.3
5.30	$48^3 \cdot 96$	0.069	3.3	290	5.0

- Use **stochastic noise source** (“one-end trick”)

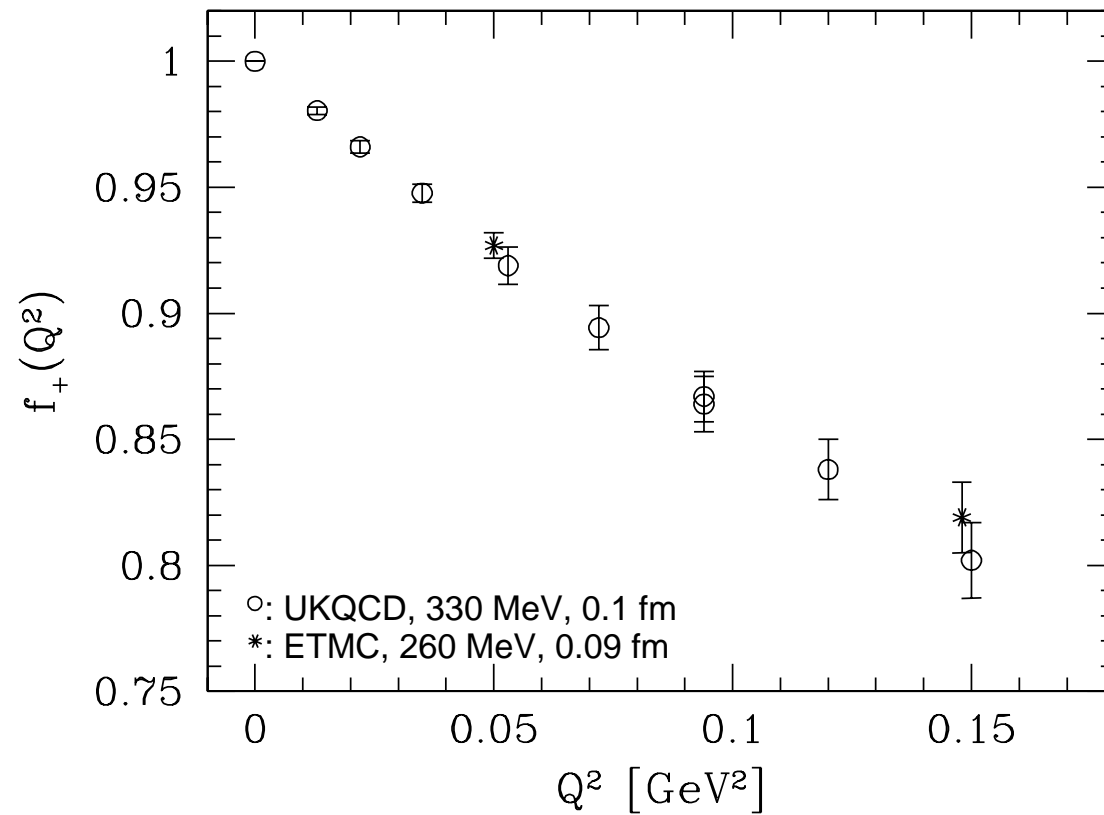
[E. Endreß, Diploma thesis, 2009]



Pion form factor

[Brandt, Della Morte, Djukanovic, Endreß, Gegelia, Jüttner, H.W.]

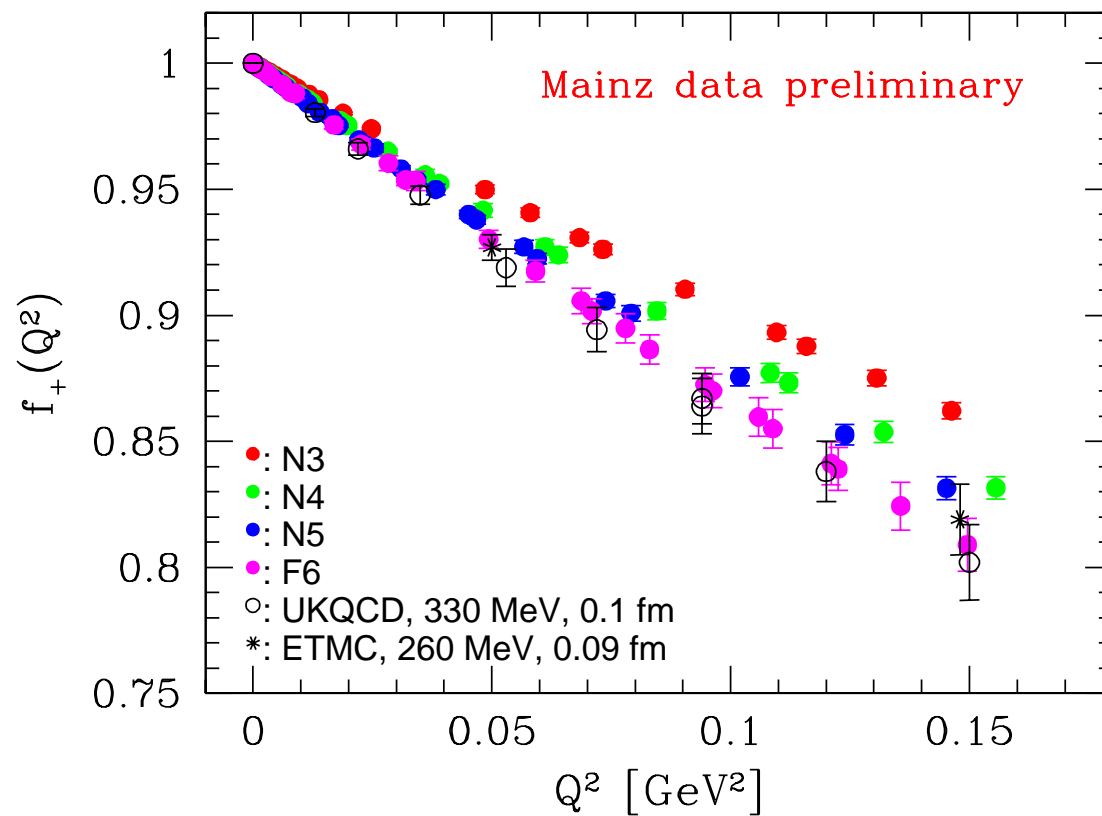
- Recently published results



Pion form factor

[Brandt, Della Morte, Djukanovic, Endreß, Gegelia, Jüttner, H.W.]

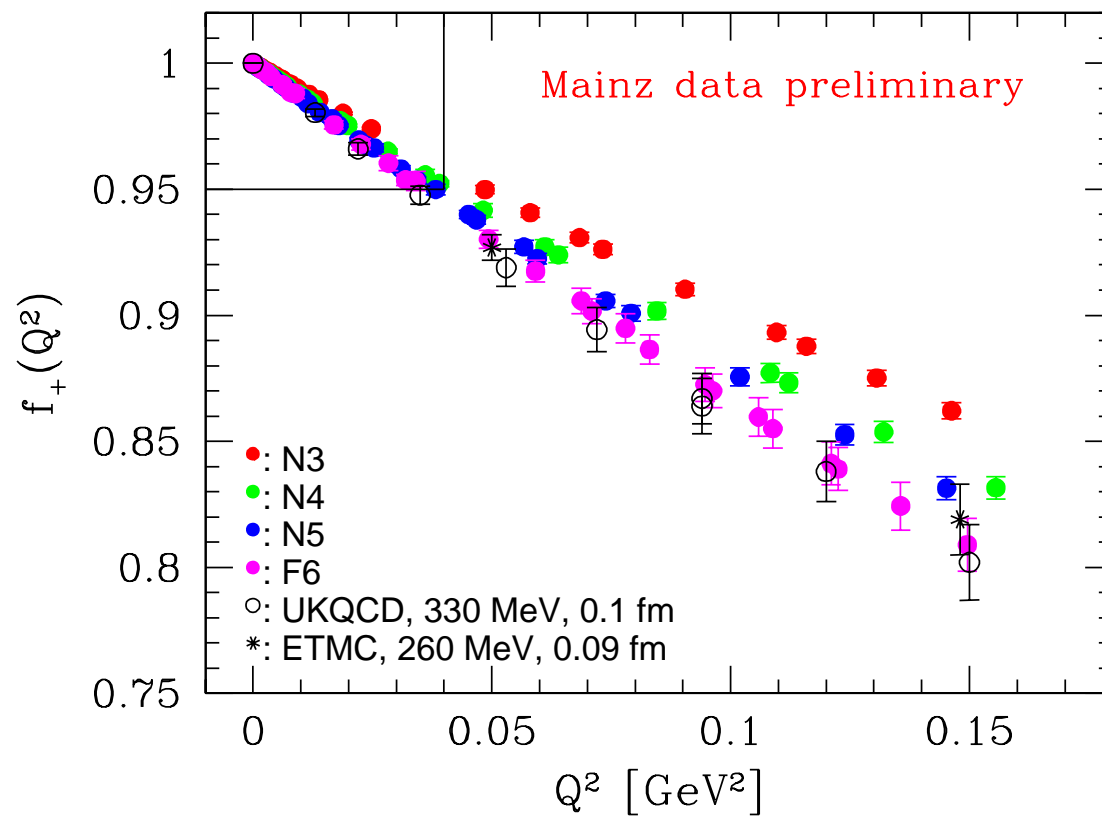
- Comparison with Mainz data



Pion form factor

[Brandt, Della Morte, Djukanovic, Endreß, Gegelia, Jüttner, H.W.]

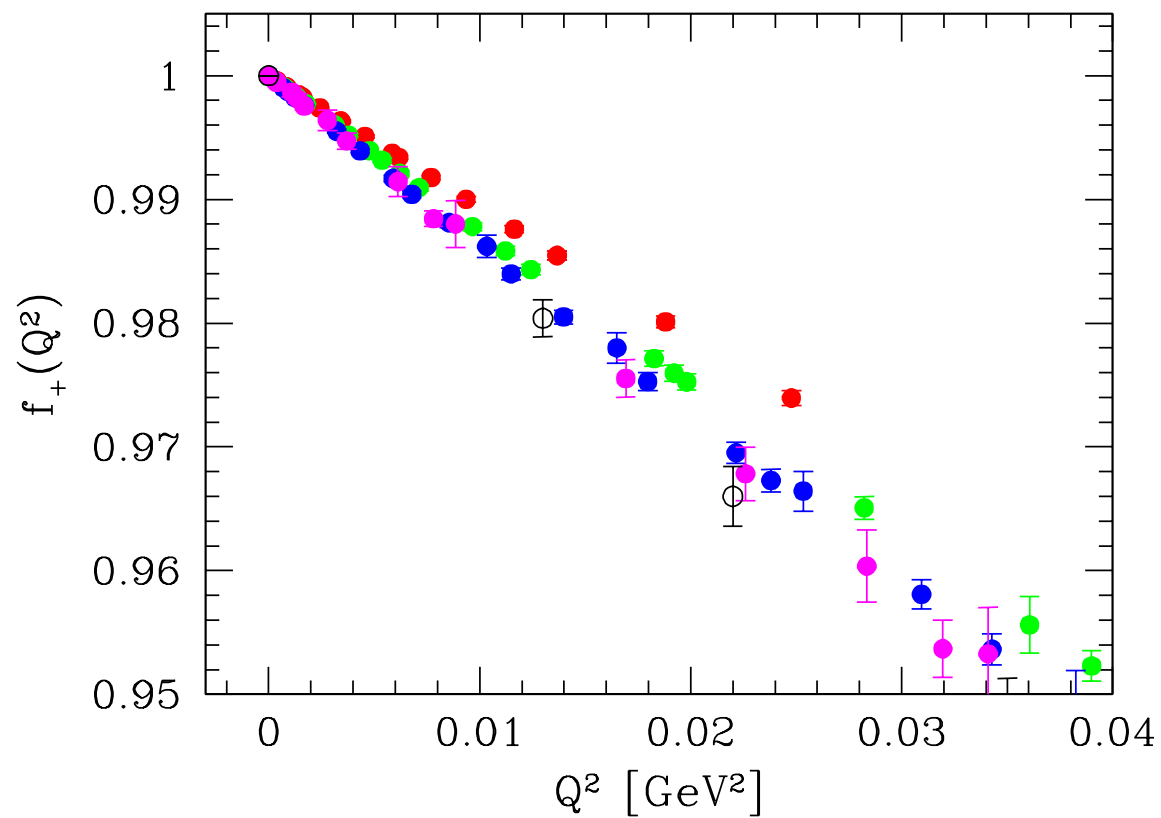
- Comparison with Mainz data



Pion form factor

[Brandt, Della Morte, Djukanovic, Endreß, Gegelia, Jüttner, H.W.]

- Comparison with Mainz data



Pion charge radius

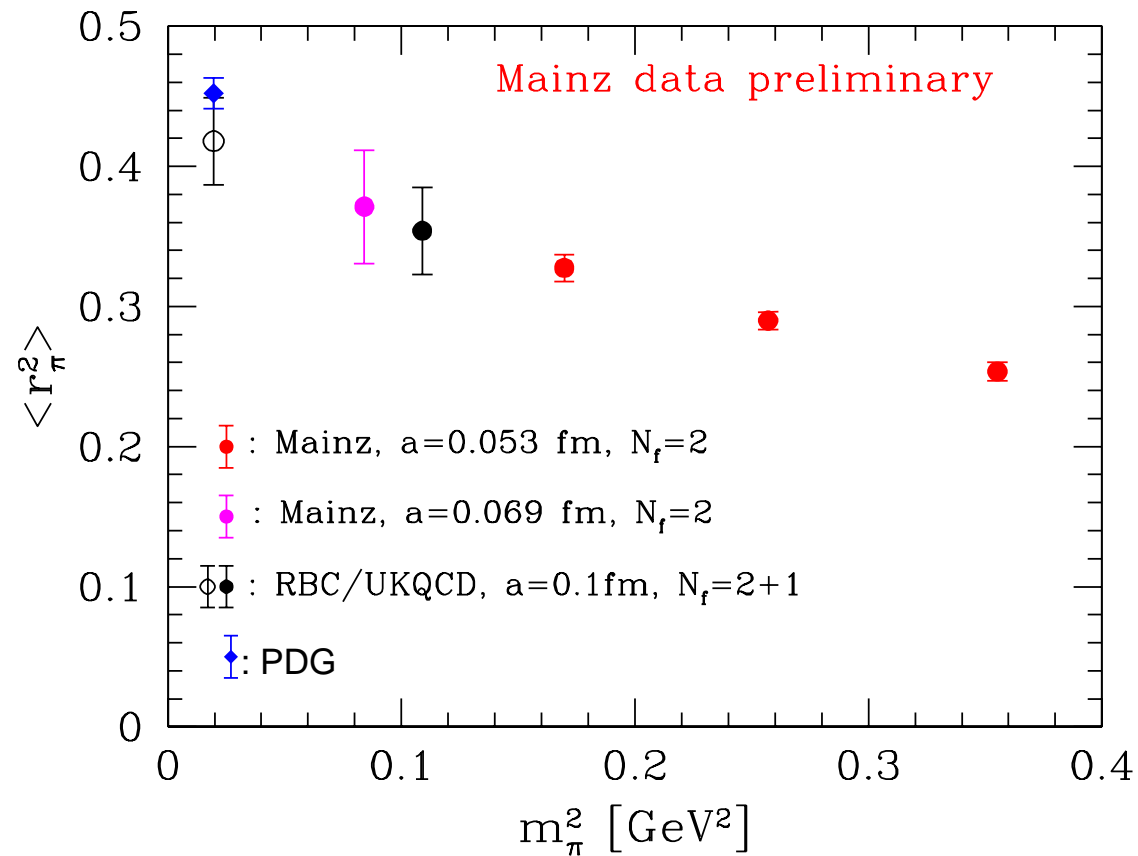
[Brandt, Della Morte, Djukanovic, Endreß, Gegelia, Jüttner, H.W.]

- Twisted boundary conditions: accurate data near $Q^2 = 0$
 - extract charge radius from linear slope

Pion charge radius

[Brandt, Della Morte, Djukanovic, Endreß, Gegelia, Jüttner, H.W.]

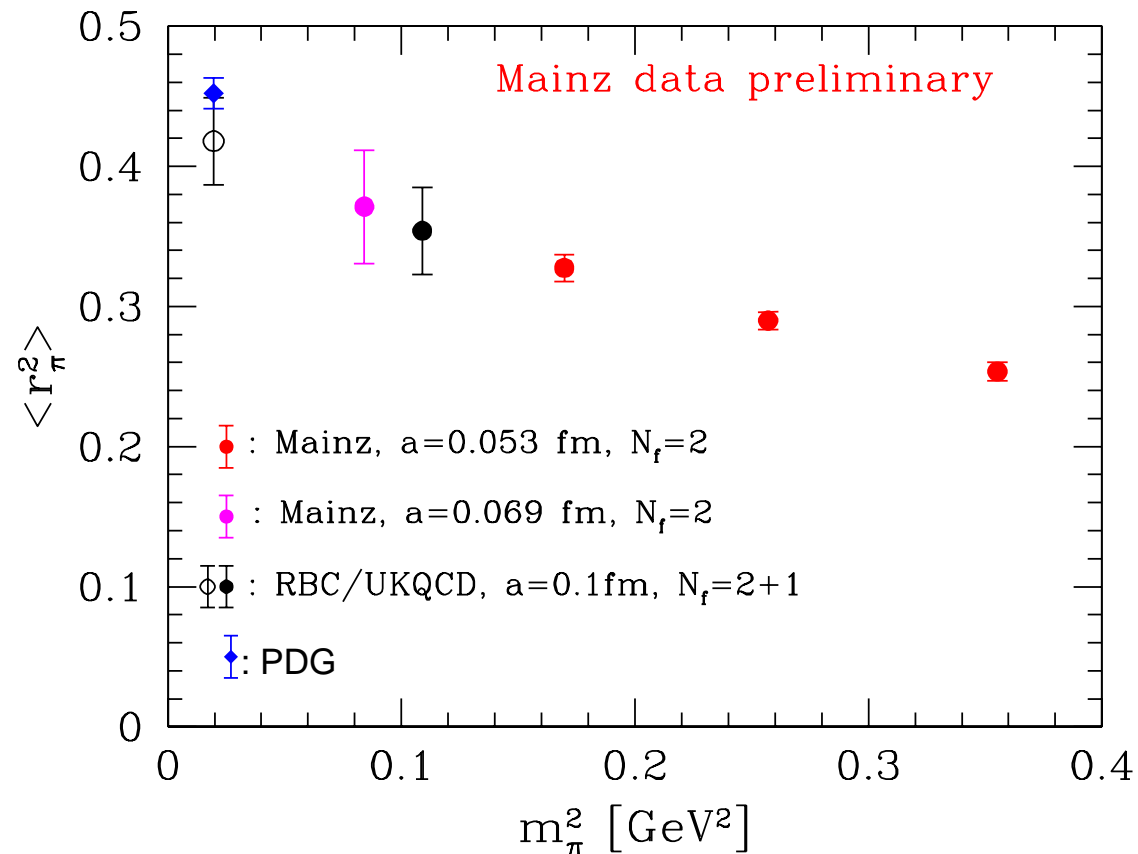
- Twisted boundary conditions: accurate data near $Q^2 = 0$
 - extract charge radius from linear slope



Pion charge radius

[Brandt, Della Morte, Djukanovic, Endreß, Gegelia, Jüttner, H.W.]

- Twisted boundary conditions: accurate data near $Q^2 = 0$
 - extract charge radius from linear slope



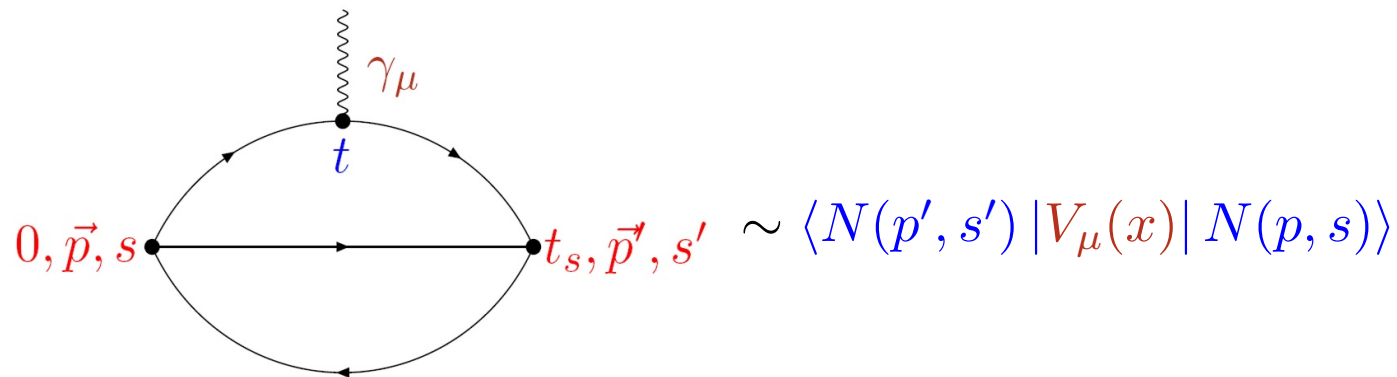
- Still to come: Fits to ChPT including vector degrees of freedom

3. Form factors and axial charge of the nucleon

- Dirac and Pauli form factors

$$\langle N(p', s') | V_\mu(x) | N(p, s) \rangle = \bar{u}(p', s') \left[\gamma_\mu F_1(q^2) + \sigma_{\mu\nu} \frac{q_\nu}{2m_N} F_2(q^2) \right] u(p, s)$$

$$G_E(q^2) = F_1(q^2) - \frac{q^2}{(2m_N)^2} F_2(q^2), \quad G_M(q^2) = F_1(q^2) + F_2(q^2)$$

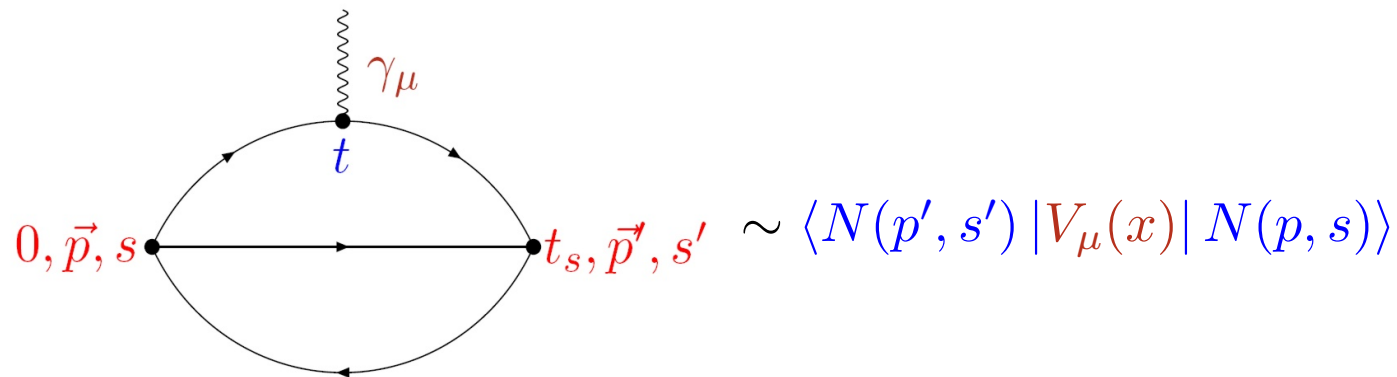


3. Form factors and axial charge of the nucleon

- Dirac and Pauli form factors

$$\langle N(p', s') | V_\mu(x) | N(p, s) \rangle = \bar{u}(p', s') \left[\gamma_\mu F_1(q^2) + \sigma_{\mu\nu} \frac{q_\nu}{2m_N} F_2(q^2) \right] u(p, s)$$

$$G_E(q^2) = F_1(q^2) - \frac{q^2}{(2m_N)^2} F_2(q^2), \quad G_M(q^2) = F_1(q^2) + F_2(q^2)$$



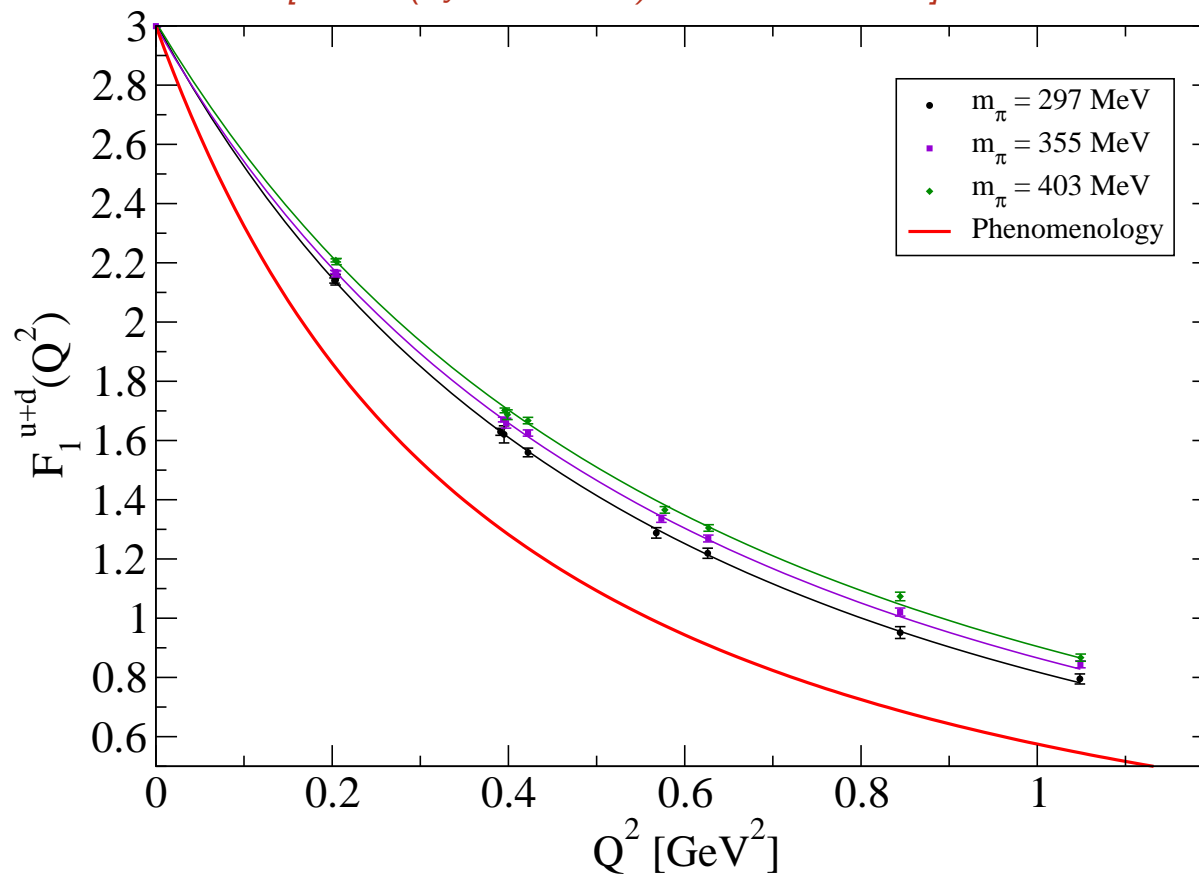
- Disconnected diagrams contribute to iso-scalar form factors
- Twisted boundary conditions: may incur large finite-volume effects?

Current Status

[Dru Renner @ Lattice 2009, Dina Alexandrou @ Lattice 2010]

- **Form factors:** experimental Q^2 -dependence & charge radii **not** reproduced

[LHPC (Syritsin et al.), arXiv:0907.4194]

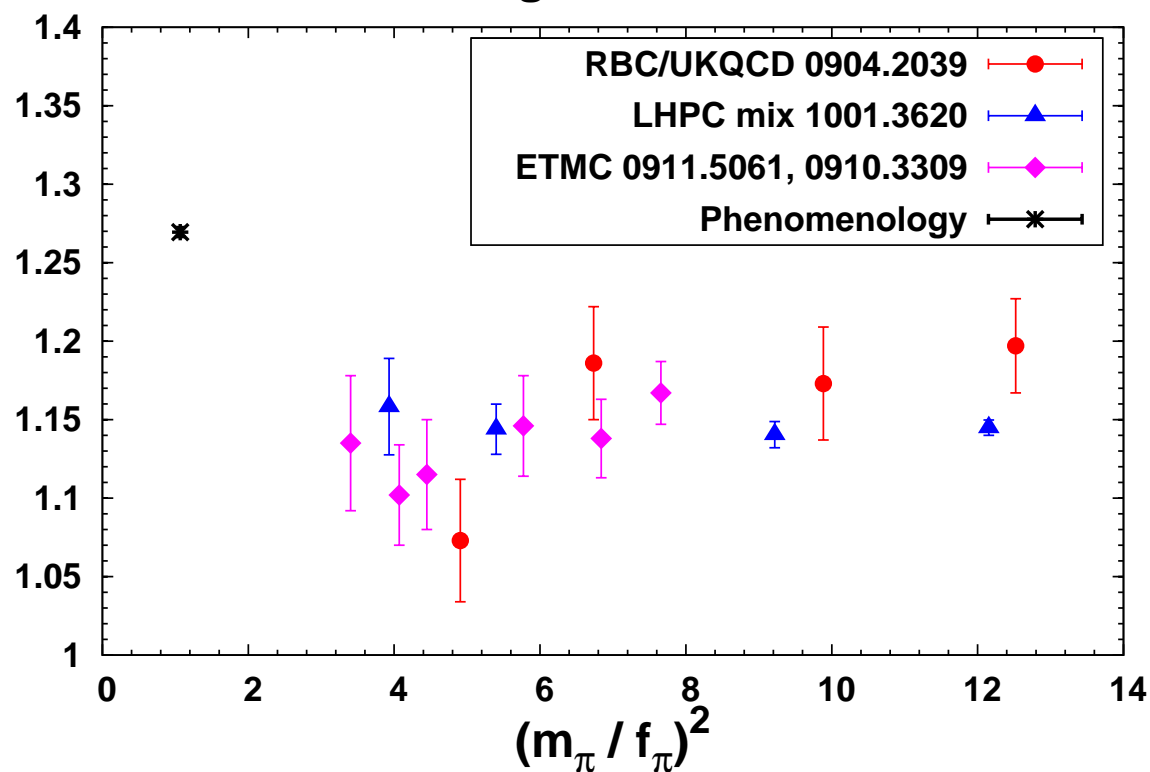


Current Status

[Dru Renner @ Lattice 2009, Dina Alexandrou @ Lattice 2010]

- **Form factors:** experimental Q^2 -dependence & charge radii **not** reproduced
- Lattice simulations produce low values for **axial charge** g_A

Axial Charge of the Nucleon



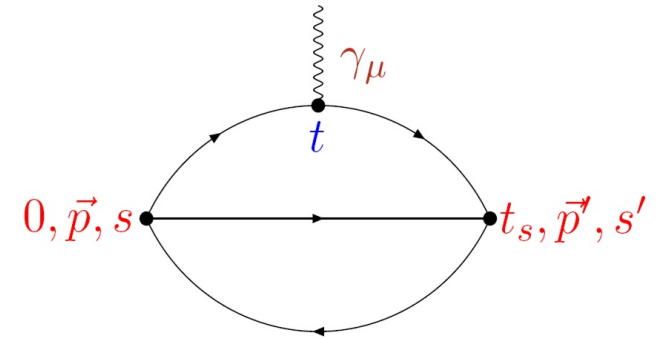
Current Status

[Dru Renner @ Lattice 2009, Dina Alexandrou @ Lattice 2010]

- **Form factors:** experimental Q^2 -dependence & charge radii **not** reproduced
- Lattice simulations produce low values for **axial charge** g_A
- Possible origin:
 - Lattice artefacts
 - Chiral extrapolations (pion masses too large)
 - Finite-volume effects
 - Contamination from excited states

Standard method

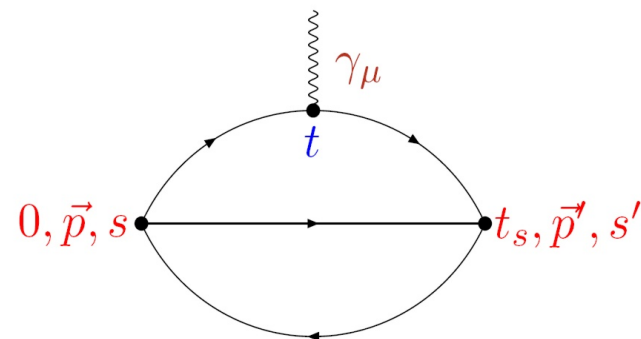
- Extract nucleon form factors from ratios of three- and two-point functions:



$$R(\vec{q}; t, t_s) = \frac{C_3(\vec{q}, t, t_s)}{C_2(\vec{0}, t_s)} \cdot \left\{ \frac{C_2(\vec{q}, t_s - t) C_2(\vec{0}, t) C_2(\vec{0}, t_s)}{C_2(\vec{0}, t_s - t) C_2(\vec{q}, t) C_2(\vec{q}, t_s)} \right\}^{1/2} \propto G_E(q^2), G_M(q^2)$$

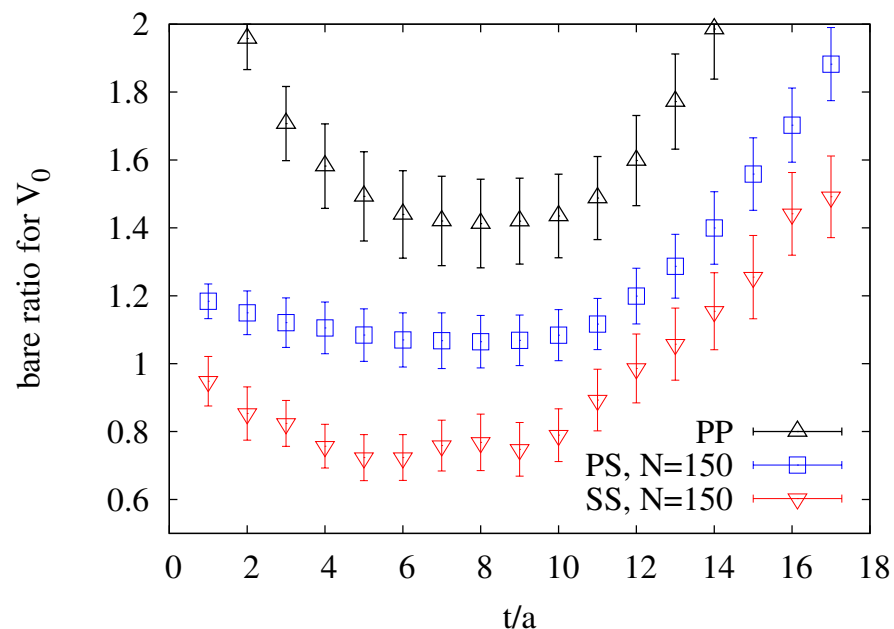
Standard method

- Extract nucleon form factors from ratios of three- and two-point functions:



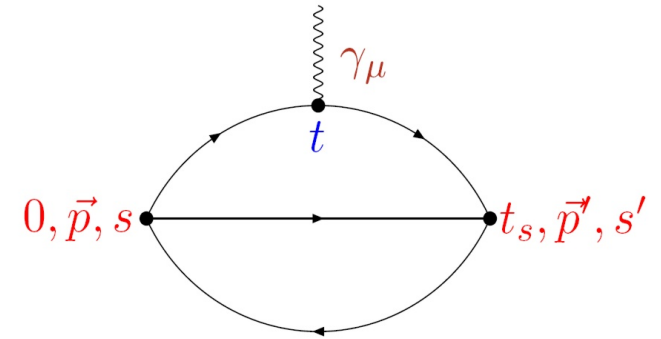
$$R(\vec{q}; t, t_s) = \frac{C_3(\vec{q}, t, t_s)}{C_2(\vec{0}, t_s)} \cdot \left\{ \frac{C_2(\vec{q}, t_s - t) C_2(\vec{0}, t) C_2(\vec{0}, t_s)}{C_2(\vec{0}, t_s - t) C_2(\vec{q}, t) C_2(\vec{q}, t_s)} \right\}^{1/2} \propto G_E(q^2), G_M(q^2)$$

- $\beta = 5.3$, $32^3 \cdot 64$, $a = 0.069$ fm
- $R(\vec{q}; t, t_s)$, iso-scalar V_0 (connected)
at $Q^2 = 0.87$ GeV, $t_s = 18$
- Several source/sink combinations



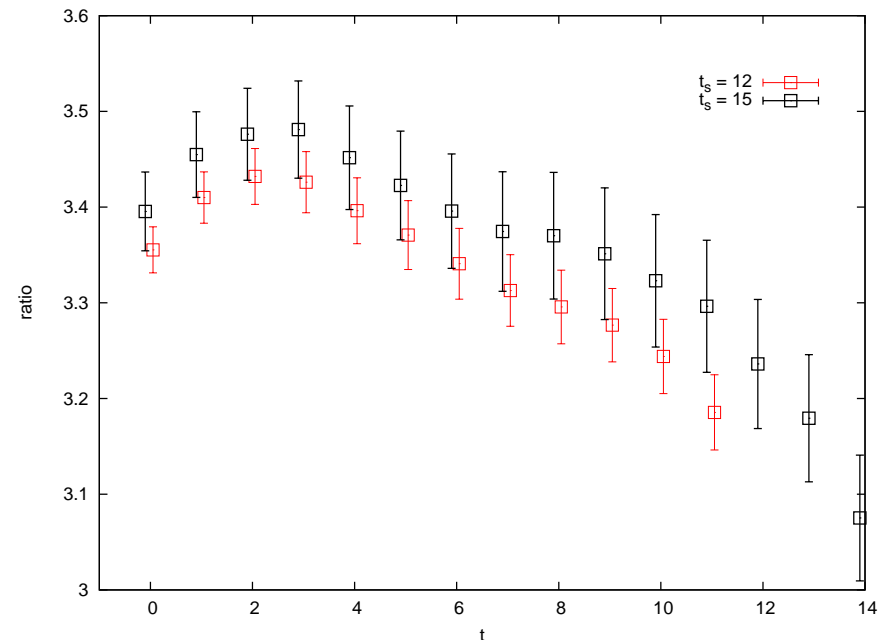
Standard method

- Extract nucleon form factors from ratios of three- and two-point functions:



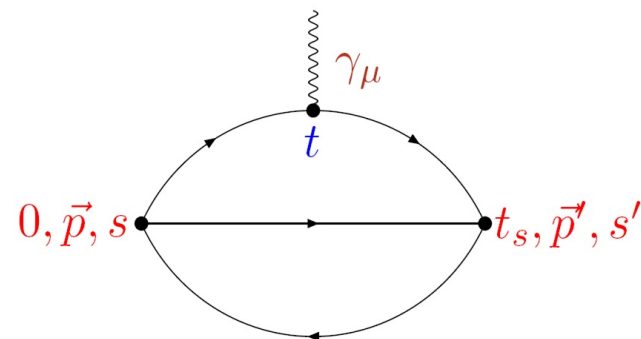
$$R(\vec{q}; t, t_s) = \frac{C_3(\vec{q}, t, t_s)}{C_2(\vec{0}, t_s)} \cdot \left\{ \frac{C_2(\vec{q}, t_s - t) C_2(\vec{0}, t) C_2(\vec{0}, t_s)}{C_2(\vec{0}, t_s - t) C_2(\vec{q}, t) C_2(\vec{q}, t_s)} \right\}^{1/2} \propto G_E(q^2), G_M(q^2)$$

- $\beta = 5.3, \quad 32^3 \cdot 64, \quad a = 0.069 \text{ fm}$
- Iso-vector magnetic form factor at $Q^2 = 0.30 \text{ GeV}, t_s = 12, 15$
- Smeared-local correlator



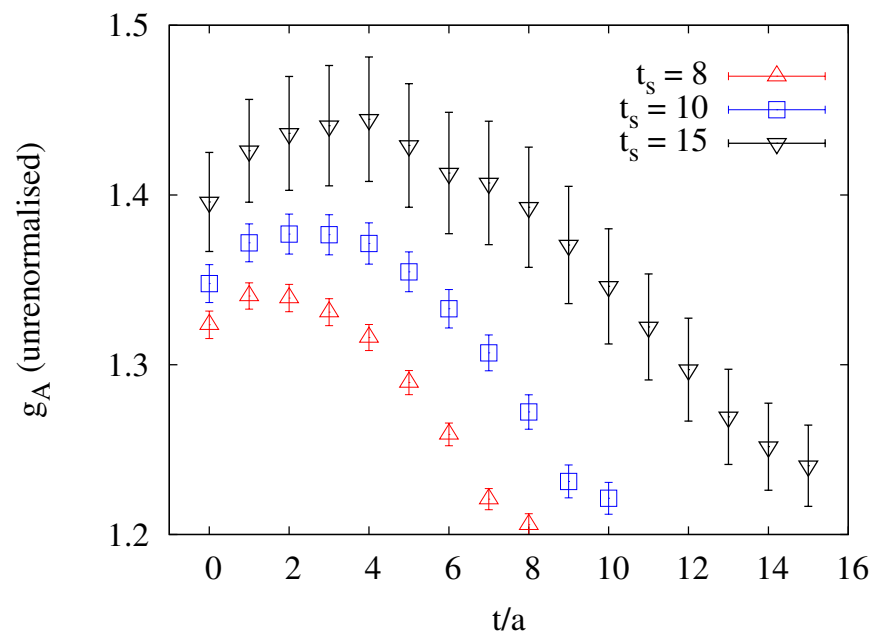
Standard method

- Extract nucleon form factors from ratios of three- and two-point functions:



$$R(\vec{q}; t, t_s) = \frac{C_3(\vec{q}, t, t_s)}{C_2(\vec{0}, t_s)} \cdot \left\{ \frac{C_2(\vec{q}, t_s - t) C_2(\vec{0}, t) C_2(\vec{0}, t_s)}{C_2(\vec{0}, t_s - t) C_2(\vec{q}, t) C_2(\vec{q}, t_s)} \right\}^{1/2} \propto G_E(q^2), G_M(q^2)$$

- $\beta = 5.3$, $32^3 \cdot 64$, $a = 0.069$ fm
- Iso-vector axial charge at $m_\pi = 550$ MeV, $t_s = 12, 15$
- Smeared-local correlator



Summed insertions

[Maiani et al. 1987; B. Knippschild @ Lattice2010]

- Standard method:

$$R(\vec{q}, t, t_s) = R_G(\vec{q}) + O(e^{-\Delta t}) + O(e^{-\Delta'(t_s-t)})$$

- Summed insertion:

$$\sum_{t=0}^{t_s} R(\vec{q}, t, t_s) = R_G(\vec{q}) \cdot t_s + K(\Delta, \Delta') + O(e^{-\Delta t_s}) + O(e^{-\Delta' t_s})$$

- Excited state contributions more strongly suppressed
- Determine $R_G(\vec{q})$ from linear slope of summed ratio

Summed insertions

[Maiani et al. 1987; B. Knippschild @ Lattice2010]

- Standard method:

$$R(\vec{q}, t, t_s) = R_G(\vec{q}) + O(e^{-\Delta t}) + O(e^{-\Delta'(t_s-t)})$$

- Summed insertion:

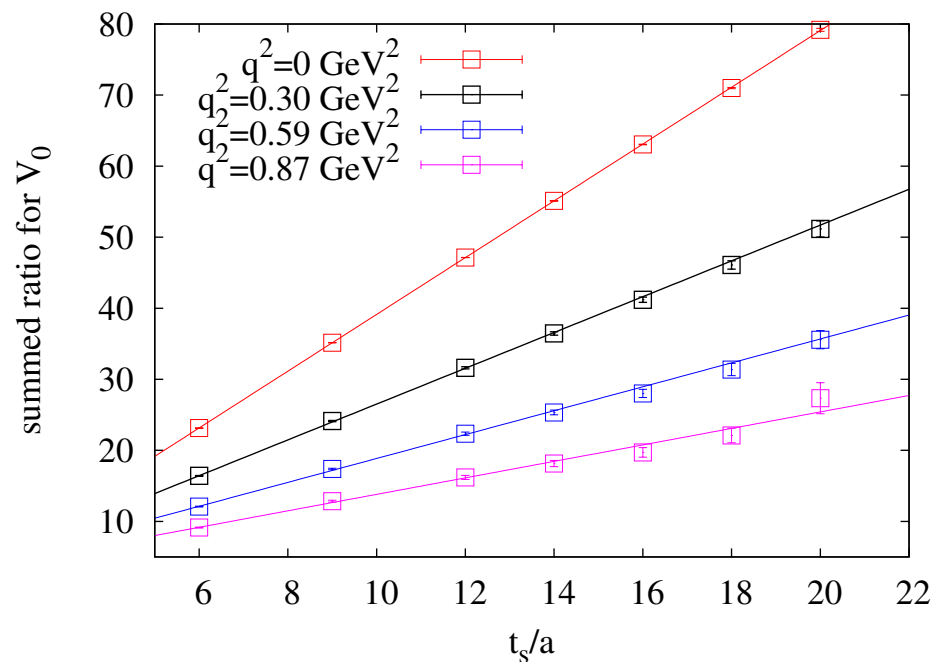
$$\sum_{t=0}^{t_s} R(\vec{q}, t, t_s) = R_G(\vec{q}) \cdot t_s + K(\Delta, \Delta') + O(e^{-\Delta t_s}) + O(e^{-\Delta' t_s})$$

- $\beta = 5.3$, $32^3 \cdot 64$,

$$a = 0.069 \text{ fm}$$

- Connected Iso-scalar form factor

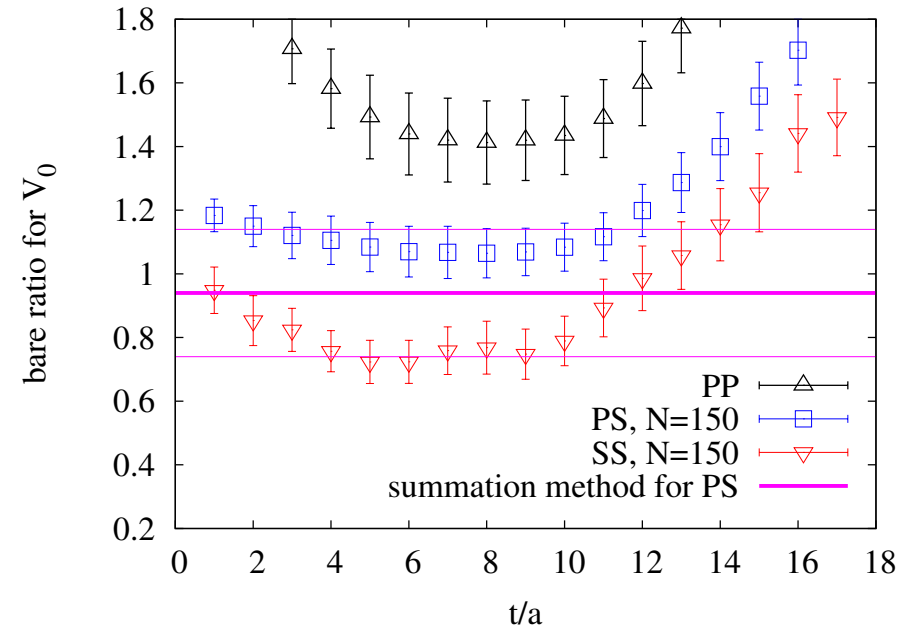
- Smeared-local correlators



Summed insertions

- $\beta = 5.3$, $32^3 \cdot 64$, $a = 0.069$ fm
- $R(\vec{q}; t, t_s)$, iso-scalar V_0 (connected)
at $Q^2 = 0.87$ GeV, $t_s = 18$
- Several source/sink combinations

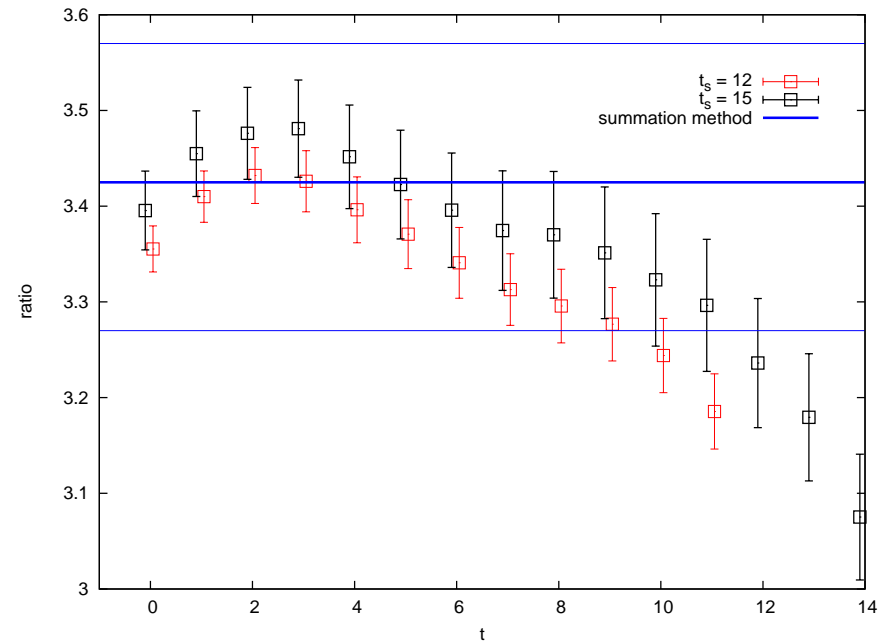
[Maiani et al. 1987; B. Knippschild @ Lattice2010]



Summed insertions

- $\beta = 5.3$, $32^3 \cdot 64$, $a = 0.069$ fm
- Iso-vector magnetic form factor at $Q^2 = 0.30$ GeV, $t_s = 12, 15$
- Smeared-local correlator

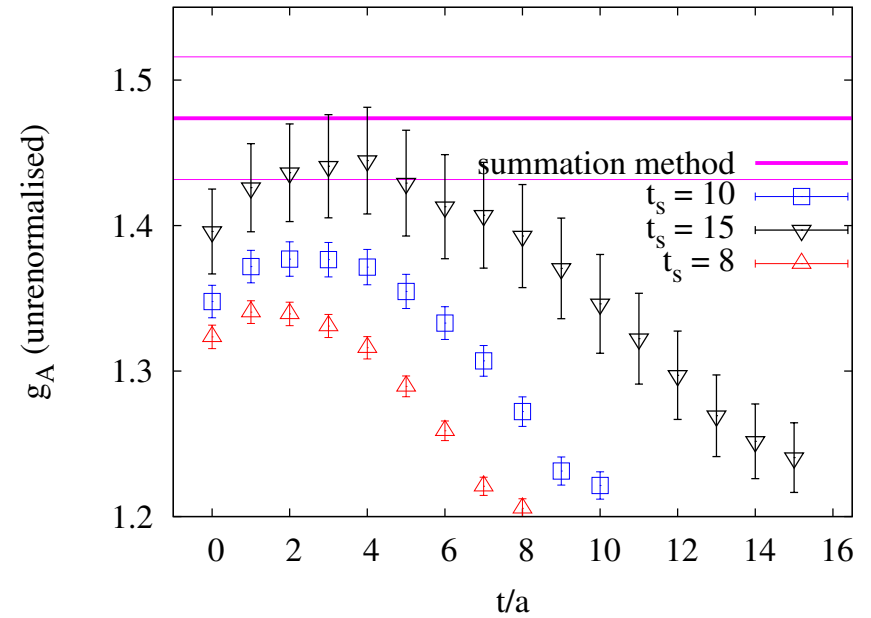
[Maiani et al. 1987; B. Knippschild @ Lattice2010]



Summed insertions

- $\beta = 5.3$, $32^3 \cdot 64$, $a = 0.069$ fm
- Iso-vector axial charge at $m_\pi = 550$ MeV, $t_s = 12, 15$
- Smearred-local correlator

[Maiani et al. 1987; B. Knippschild @ Lattice2010]



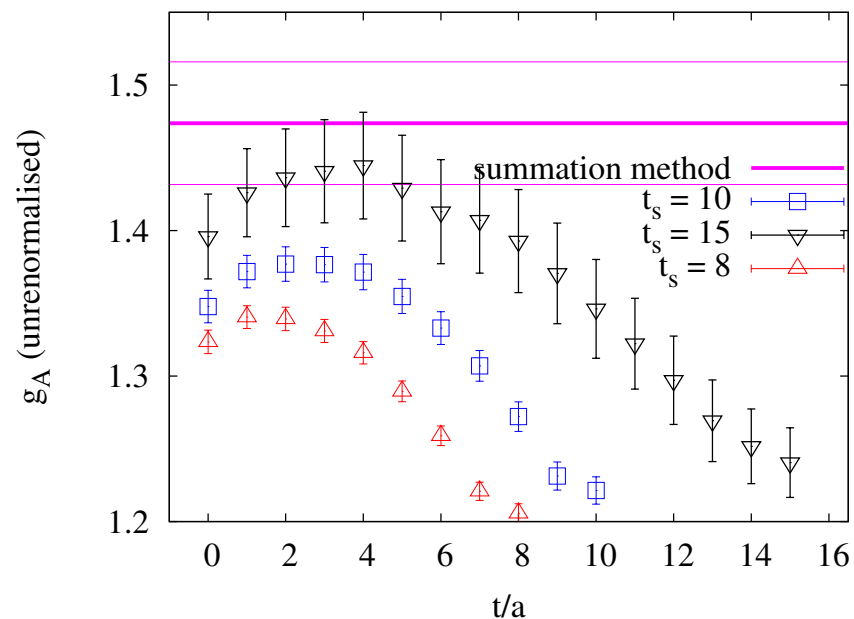
Summed insertions

[Maiani et al. 1987; B. Knippschild @ Lattice2010]

- $\beta = 5.3$, $32^3 \cdot 64$, $a = 0.069$ fm

- Iso-vector axial charge at $m_\pi = 550$ MeV, $t_s = 12, 15$

- Smearred-local correlator



- Better control over excited state contamination

- Larger statistical errors

- Requires more values of t_s

Results

[Capitani, Della Morte, Jüttner, Knippschild, Meyer, H.W.]

- Ensembles at $\beta = 5.3$, $a = 0.069$ fm, $32^3 \cdot 64$ and $48^3 \cdot 96$

$$m_\pi = 290 - 590 \text{ MeV}, \quad L = 2.2 \text{ and } 3.3 \text{ fm}, \quad Lm_\pi \gtrsim 5$$

Employ summed insertions for ≈ 6 values of t_s

Results

[Capitani, Della Morte, Jüttner, Knippschild, Meyer, H.W.]

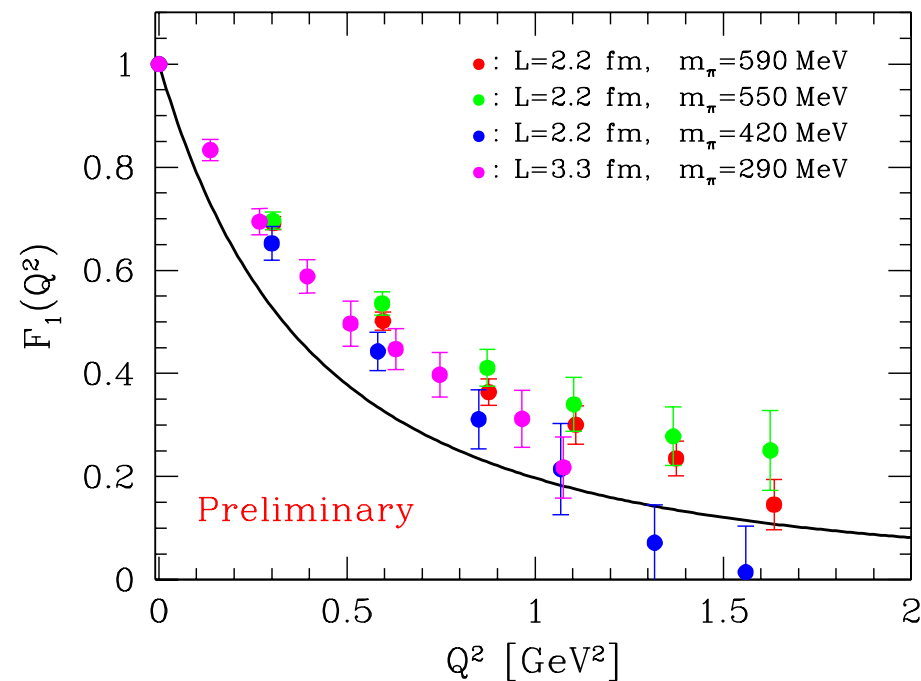
- Ensembles at $\beta = 5.3$, $a = 0.069$ fm, $32^3 \cdot 64$ and $48^3 \cdot 96$

$$m_\pi = 290 - 590 \text{ MeV}, \quad L = 2.2 \text{ and } 3.3 \text{ fm}, \quad Lm_\pi \gtrsim 5$$

Employ summed insertions for ≈ 6 values of t_s

- Dirac form factor:

[Knippschild @ Lattice 2010]



Results

[Capitani, Della Morte, Jüttner, Knippschild, Meyer, H.W.]

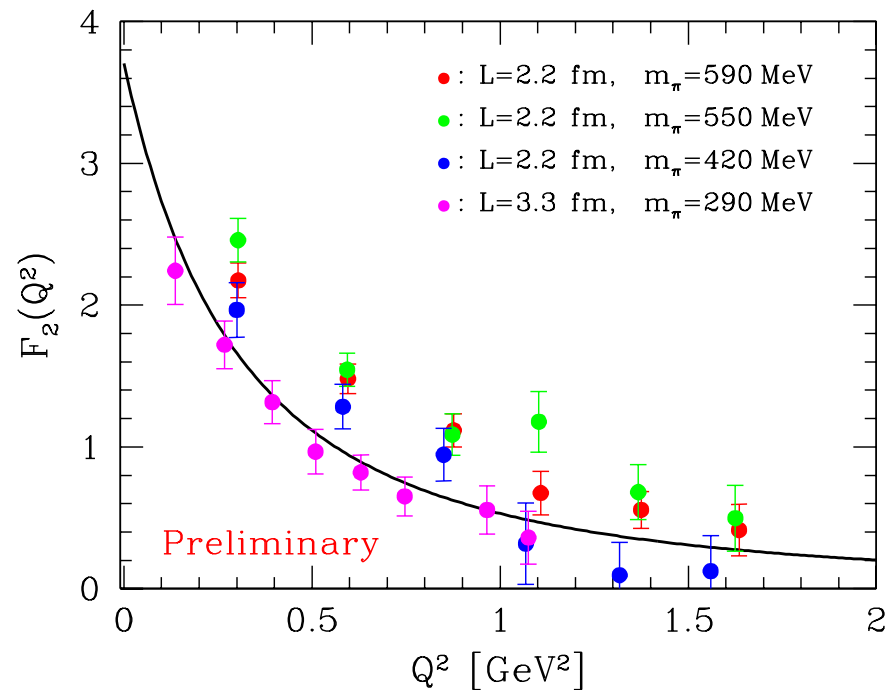
- Ensembles at $\beta = 5.3$, $a = 0.069$ fm, $32^3 \cdot 64$ and $48^3 \cdot 96$

$$m_\pi = 290 - 590 \text{ MeV}, \quad L = 2.2 \text{ and } 3.3 \text{ fm}, \quad Lm_\pi \gtrsim 5$$

Employ summed insertions for ≈ 6 values of t_s

- Pauli form factor:

[Knippschild @ Lattice 2010]



Results

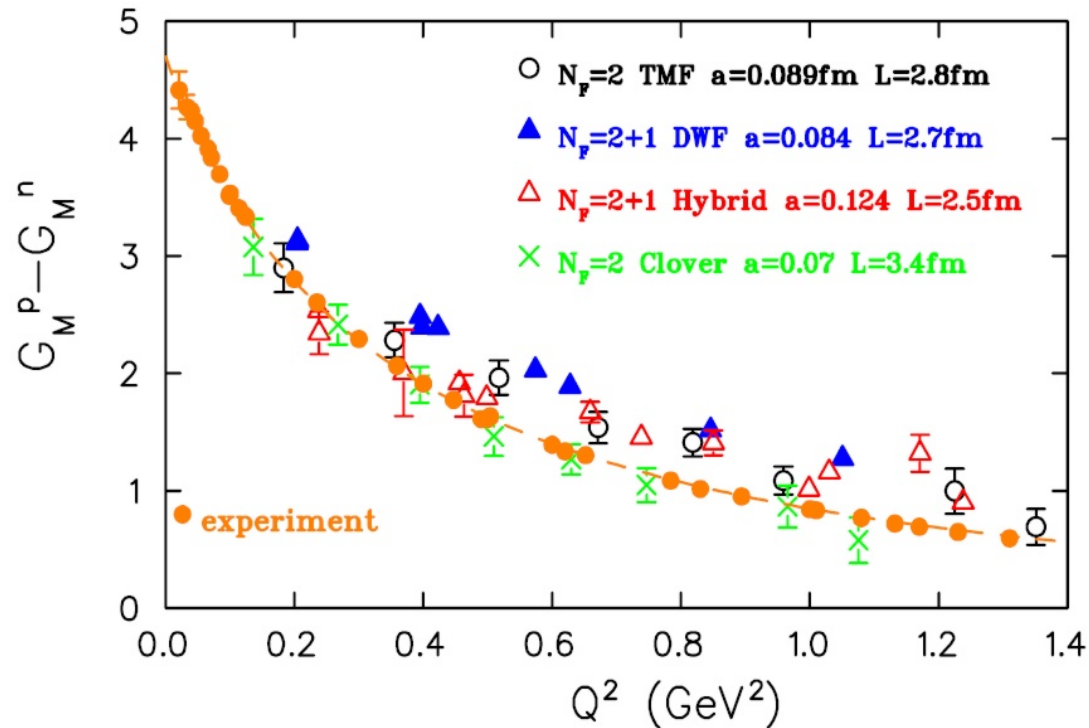
[Capitani, Della Morte, Jüttner, Knippschild, Meyer, H.W.]

- Ensembles at $\beta = 5.3$, $a = 0.069$ fm, $32^3 \cdot 64$ and $48^3 \cdot 96$

$$m_\pi = 290 - 590 \text{ MeV}, \quad L = 2.2 \text{ and } 3.3 \text{ fm}, \quad Lm_\pi \gtrsim 5$$

Employ summed insertions for ≈ 6 values of t_s

- Comparison: magnetic form factor @ $m_\pi \simeq 300$ MeV [D. Alexandrou @ Lattice 2010]



Results

[Capitani, Della Morte, Jüttner, Knippschild, Meyer, H.W.]

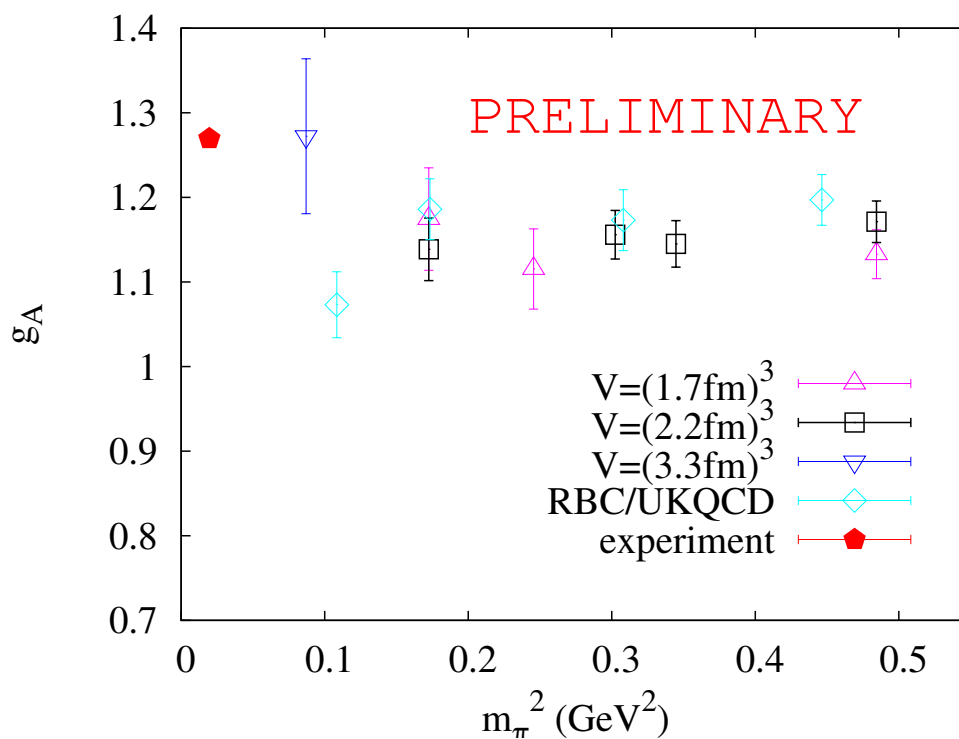
- Ensembles at $\beta = 5.3$, $a = 0.069$ fm, $32^3 \cdot 64$ and $48^3 \cdot 96$

$$m_\pi = 290 - 590 \text{ MeV}, \quad L = 2.2 \text{ and } 3.3 \text{ fm}, \quad Lm_\pi \gtrsim 5$$

Employ summed insertions for ≈ 6 values of t_s

- Axial charge:

[Knippschild @ Lattice 2010]



Results

[Capitani, Della Morte, Jüttner, Knippschild, Meyer, H.W.]

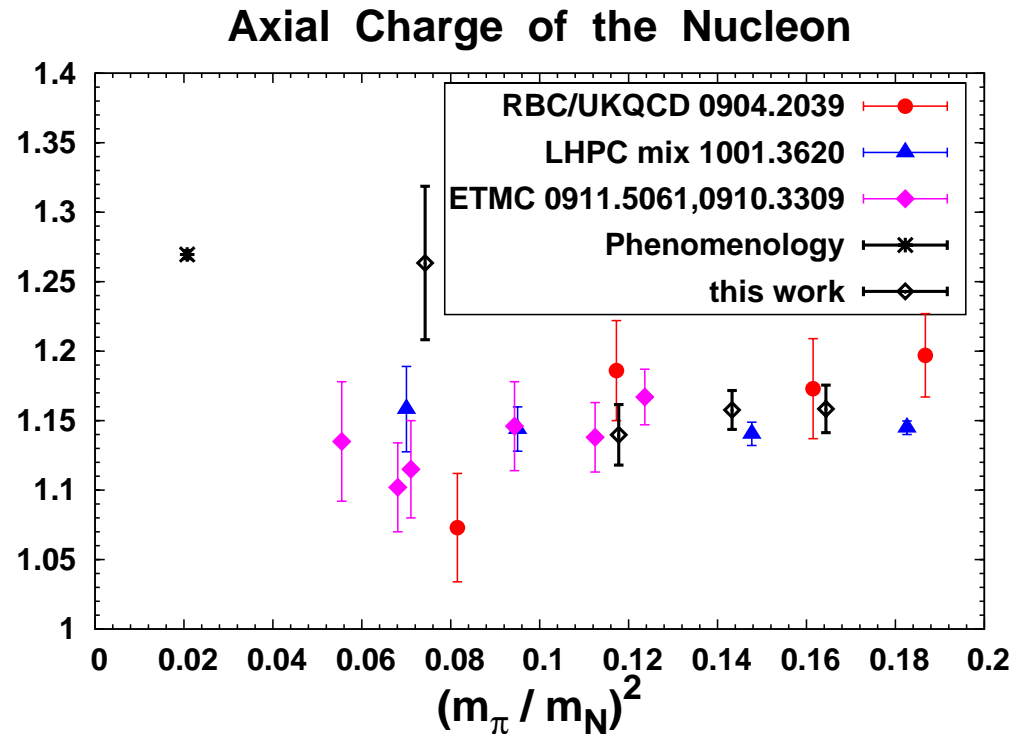
- Ensembles at $\beta = 5.3$, $a = 0.069$ fm, $32^3 \cdot 64$ and $48^3 \cdot 96$

$$m_\pi = 290 - 590 \text{ MeV}, \quad L = 2.2 \text{ and } 3.3 \text{ fm}, \quad Lm_\pi \gtrsim 5$$

Employ summed insertions for ≈ 6 values of t_s

- Axial charge:

[Knippschild @ Lattice 2010]



Still to come. . .

- Process CLS ensembles at $\beta = 5.2, a = 0.08 \text{ fm}$ and $\beta = 5.5, a = 0.05 \text{ fm}$
→ Lattice artefacts
- Comparison of $L \simeq 2.6 \text{ fm}$ and $L \simeq 3.8 \text{ fm}$ at $\beta = 5.2$

Summary

- Nucleon form factors and axial charge:
“2nd Generation Benchmark” for lattice QCD
- Progress in controlling systematic uncertainties in form factor calculations using CLS ensembles
- Pion form factor:
 - precise, model-independent estimates of $\langle r_\pi^2 \rangle$ via **twisted boundary conditions**
- Nucleon form factors and g_A :
 - **summed insertions** help control excited state contamination
 - situation still not settled (pion masses, volumes, discretisation effects)