### 中性子星における Σハイペロン混合の効果について

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### Introduction

Our interest: dense matter EOS including Y

NS matter, SN



[1] Lattimer and Swesty, NPA 535 (1991) 331. [2] Shen et al., PTP 100 1013. [3] Yamamoto, Nishizaki and Takatsuka, NPA 691 (2001) 432c.
[4] Schaffner and Mishustin PRC 53 (1996) 1416. [5] Schaffner-Bielich, NPA 804 (2008) 309..

### Chiral symmetry in hadron phase

- Chiral symmetry in hadron:  $\pi$  as Nambu-Goldstone boson and its chiral partner,  $\sigma$
- Chiral potential in RMF model: suggested from QCD (c.f. NJL model) corresponding to chiral restoration medium
- $\omega$  and  $\rho$  mass modification[12]
  - Evidences of partial chiral symmetry restoration in nuclear medium?
  - Theoretical model overestimates mass modification compared to [13] where all of vector meson mass is gene condensate.





### Motivation

Our interest: dense matter EOS with Y



[1] Lattimer and Swesty, NPA 535 (1991) 331. [2] Shen et al., PTP 100 1013. [3] Yamamoto, Nishizaki and Takatsuka, NPA 691 (2001) 432c.

### Model description

#### RMF Lagrangian



 Long standing problems: Too early chiral restoration, very stiff EOS, instability on chiral potential

What form should the meson potential have?  $\Rightarrow$  Constrained by chiral symmetry, derived from SCL-LQCD

# Chiral potential derived from SCL

- Start point: Action of gluon and fermion on strong coupling limit of Lattice QCD ( $g \rightarrow \infty$ ) at T = 0  $S = \sum_{R} + S_{F} + m_{O} \overline{\chi} \chi$
- After bosonizing fermion pair and integrating fermion  $U = -\frac{a}{2} \log(\sigma) + b\sigma^2 \rightarrow -\frac{a}{2} \log(\det MM^{\dagger}) + b \operatorname{tr}(MM^{\dagger})$

 $SU_{f}(2) \operatorname{case}_{U=-\frac{a}{2}\log(\det MM^{\dagger}) + b\operatorname{tr}(MM^{\dagger}) - c_{\sigma}\sigma} = -\frac{a}{2}\log(\det MM^{\dagger}) + b\operatorname{tr}(MM^{\dagger}) - c_{\sigma}\sigma - c_{\zeta}\zeta} SU_{f}(3) \operatorname{case}_{+o(\det M+\det M^{\dagger})} = -a\left[\log(f_{\pi} + \phi_{\pi}) - \phi_{\tau} + \frac{1}{2}\left(\frac{\phi_{\pi}}{f_{\pi}}\right)^{2}\right] + \frac{1}{2}m_{\sigma'}\phi_{\sigma'} + \frac{1}{2}m_{\pi'}\pi^{2} = -\frac{a}{2}\left[\log(f_{\pi} + \phi_{\pi}) - \left(\frac{\phi_{\pi}}{f_{\pi}}\right) + \frac{1}{2}\left(\frac{\phi_{\pi}}{f_{\pi}}\right)^{2}\right] + \frac{1}{2}m_{\sigma'}\phi_{\sigma'} + \frac{1}{2}m_{\pi'}\pi^{2} = -\frac{a}{2}\left[\log(f_{\zeta} + \phi_{\zeta}) - \phi_{\zeta} + \frac{1}{2}\left(\frac{\phi_{\zeta}}{f_{\zeta}}\right)^{2}\right] + \frac{1}{2}m_{\zeta'}\phi_{\zeta'} + \frac{1}{2}m_{\kappa'}\kappa^{2} + \frac{1}{2}m_{\zeta'}\phi_{\zeta'} + \frac{1}{2}m_{\kappa'}\kappa^{2} + \frac{1}{2}m_{\zeta'}\phi_{\zeta'}\phi_{\zeta'} + \frac{1}{2}m_{\zeta'}\phi_{\zeta'}\phi_{\zeta'} + \frac{1}{2}m_{\zeta'}\phi_{\zeta'}\phi_{\zeta'} + \frac{1}{2}m_{\kappa'}\kappa^{2} + \frac{1}{2}m_{\zeta'}\phi_{\zeta'}\phi_{\zeta'} + \frac{1}{2}$ 

[8] Kawamoto and Smit, NPB 190 (1981) 100. [9] Kawamoto, Miura, Ohnishi and Ohnuma PRD **75** (2007) 014502.

### SCL3 RMF model

 Logarithmic SCL Chiral Potential OProperty of chiral potential: Neither instability nor abnormal vacuum OConsistent with reproducing *BE* of normal nuclei,  $S_{\Lambda}$  of single  $\Lambda$  nuclei,  $\Delta B_{\Lambda\Lambda}$  of  ${}^{6}_{\Lambda\Lambda}$  He and FP EOS  $(K \sim 210 MeV \text{ softened by hidden})$ strange condensate  $\zeta = \overline{ss}$ Calculated NS mass underlies the most

□ Calculated NS mass underlies the most reliable observation, 1.44 M<sub>☉</sub> when we include hyperon degrees of freedom

What are key ingredients which are efficient to refine ?

[10] KT, Maekawa, Matsumiya and Ohnishi, arXiv:0909.5058



## Scalar and vector potential in matter

- Repulsive potential from  $\omega$ : Good agreement with DBHF result by  $3\rho_0$  but seems to be insufficient from higher  $\rho_B$
- C<sub>ω</sub>: Strength of ω meson potential and known as the suppresser of vector meson field



To reduce this strength, we would like to suggest "Density dependent type coupling ( $\sigma\omega N$ )"

## Density dependent type coupling

- May be derived from NLO calculation on SCL-LQCD
- Scalar potential(∝ σ)⇒saturate
   Vector potential (∝ ω)
   ⇒increase linearly
- $\sigma\omega N$  coupling: effective around  $\rho_{0}$  and not important in high  $\rho_{B}$



Here, we introduce  $\sigma\omega N$  coupling  $g_{\omega} = g_{\omega 0} + g_{\sigma\omega N} \phi_{\sigma}$ and we examine its effect to the property of sym. nuclear matter and NS matter

### Interesting feature

- Is the ratio of hyperon fraction in NS changed by  $\boldsymbol{\Sigma}$  hyperon potential?
  - Determined by  $SU_{\rm f}(3)$  symmety and the atomic shifts(AS) of  $\Sigma^-$
  - Repulsive but weak vector-isovector meson coupling: Σ hyperon may appear on NS



### Aim of this study

- Is SCL3 RMF model be able to support maximum mass of neutron star?
  - Can density dependent type coupling resolve this problem?
- On another RMF model, is the ratio of hyperon fraction in NS also changed by Σ hyperon potential determined by the atomic shifts(AS) of Σ<sup>-</sup>?

### Result(1)

- Σ hyperon potential: determined by the atomic shifts(AS) of Σ<sup>-</sup>
- $\Sigma^{-}$  appear at lower  $\rho_{B}$  than  $\frac{3}{4}$ 
  - $\Lambda \rightarrow \text{contradict to other calc.}^{10}$

where AS are not reproduced.





### Result(2)

 To confirm this situation, we also use the phenomenological RMF model with non-linear potential(TM1)



- We can find the same situation occurred in this calculation and  $\Sigma$  can appear in NS matter even if it has a repulsive potential

### Results(3)

#### Maximum mass of Neutron Star

Neutron Star EOS with Hyperon effect: →Inclusion of Hyperon make NS EOS softened at high density phase 2.5 **Recent observation** NS mass 2 may confirm heavier M/M<sub>sun</sub> maximum mass of 1.5 neutron star,  $1.7 M_{\odot}$ 1 New SCL3 DDC RMF 0.5 model seems not to 0 be enough to explain 0.5 0 1.5 this observation.

2

### Summary

- RMF model with chiral SU(3) potential (SCL3)
  - Saturation property, incompressibility, BE, S<sub>Λ</sub>, and ΔB<sub>ΛΛ</sub> are well reproduced in appropriate parameter range.
     Calculated NS matter EOS underestimates observed NS mass.
- Key for the property in high  $\rho_{\scriptscriptstyle B}$  phase:  $C_{_{\!\omega}}$ 
  - □ Density dependent type coupling ( $\sigma\omega$ N type coupling) ⇒This needs not too large C<sub>w</sub> to reproduce nuclear property. Calculated results seems not to support 1.7 M<sub>☉</sub> if hyperon effects are taken into account properly.
- Ordering of hyperon appearance
  - $\hfill\square$  Atomic shifts of  $\Sigma^-$ : one of key ingredients of this topic
  - Ordering may be changed but that does not affect the final result so much so far.
- In future ······
  - Finite temperature EOS for supernovae simulation

### Thank you for listening!!