

General consideration

$$G_3 = \bar{F}_3 - \bar{\Phi} H_3$$

$$\widehat{F}_3 = (1 + *_{10}) d\alpha(y) \wedge dx^0 \wedge dx^1 \wedge dx^2 \wedge dx^3$$

$$ds_{10}^2 = e^{2A(y)} \sum \eta_{\mu\nu} dx^\mu dx^\nu + e^{-2A(y)} \sum \widehat{g}_{mn} dy^m dy^n$$

$$\textcircled{1} R_{MN} = \dots$$

$$M = \mu, N = \nu.$$

$$\widehat{\nabla}^2 A = e^{-2A} \left(\frac{G_{mnp} \bar{G}^{mnp}}{48 \operatorname{Im} \bar{\Phi}} + \dots \right)$$

... 2-1

② EOM for \widehat{F}_5 :

$$d\widehat{F}_5 = H_3 \wedge \widehat{F}_3 + \text{localised source} \quad \text{L2-3}$$

$$\widehat{\nabla}^2 \alpha = \frac{i e^{2A} G_{mnp} \star_6 \overline{G}^{mnp}}{12 \text{Im} \widehat{\Phi}} + e^{-6A} d_m \alpha d^m e^{4A}$$

$$+ 2 k_{10}^2 e^{2A} T_3 P_3^{\text{loc}}$$

... L2-2

$$(\star_6 G)_{mnp} = \epsilon_{mnpqrs} G^{qrs}$$

L2-1, L2-2 \Rightarrow

$$\hat{\nabla}^2 (e^{4A} - \alpha) = \frac{e^{2A}}{6 \text{Im} \tilde{\phi}} |iG_3 - \star_6 G_3|^2$$

$$+ 2e^{-6A} |d(e^{4A} - \alpha)|^2$$

$$+ 2K_{10}^2 e^{2A} \left[\frac{1}{4} (T_m^m - T^m_m) - T_3 f_3^{\text{loc}} \right]$$

vanishes for D7|07 & D3|03

$$\left(\begin{array}{l} \text{for D3|03} \\ -T_\mu^{\mu \text{loc}} + T_m^m \text{loc} = T_3 - 4 \delta(\Sigma) \end{array} \right) \text{ planes}$$

Conclude

$$iG_3 = \star_6 G_3$$

→ Imaginary self dual

(L2-5)

$$e^{4A} = \alpha$$

$$(\star_6)^2 = -1$$

Other EOM :

$$dF_3 = d * F_3 = 0$$

① $F_{(3)}$ & $H_{(3)}$ Harmonic, closed ... L2-6

② $\hat{g}_{mn}(y)$, $\bar{\Phi}(y)$: solve EOM without flux

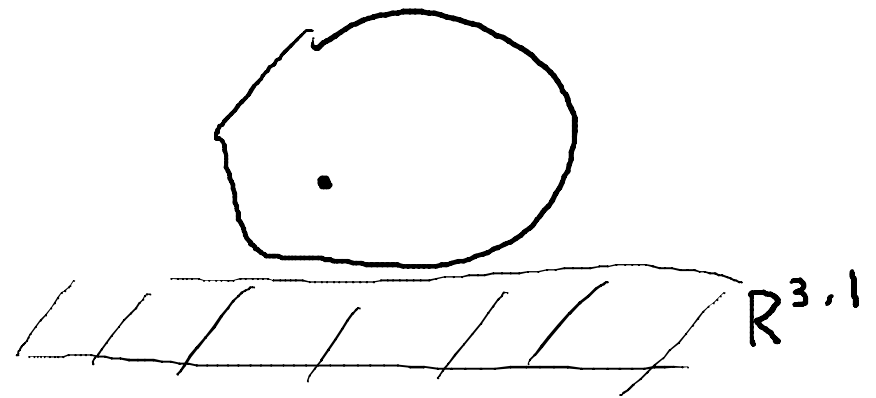
(This means CY_3 ($-O_3^+$) are solutions
with $\phi(y) = \text{const.}$)

(L2-3)

$$\int_{\widehat{M}} \widehat{F}_5 = \int H_3 \wedge F_3 + 2K_{10}^2 T_3 (N_{D3} - N_{\overline{D3}})$$

$$\int_{\widehat{M}} H_3 \wedge F_3 + 2K_{10}^2 T_3 N_{D3} = 2K_{10}^2 T_3 N_{\overline{D3}}$$

... (L2-4)



$H^3(M)$: 3rd cohomology group of \widehat{M}

(L2-6) F_3 & $H_3 \in H^3(\widehat{M})$

\int
D string F string

F_3 & H_3 : Dirac quantisation

$$\alpha^i \in H^3(M, \mathbb{Z}) \quad , \quad \gamma_j \in H_3(M)$$

$$\int_{\gamma_j} \alpha^i = p^i_j \quad , \quad p^i_j \in \mathbb{Z}$$

$$\frac{1}{(2\pi)^2} \int \alpha^i F_3 = \sum m_i \alpha^i$$

$$\frac{1}{(2\pi)^2} \int \alpha^i H_3 = \sum n_i \alpha^i$$

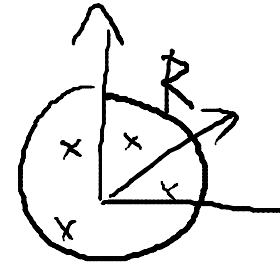
$$m_i, n_i \in \mathbb{Z}$$

$$\dim(H^3(M)) = b_3$$

Total number of choices for \bar{F}_3, H_3

$$\sim (\sqrt{|N_{03}|})^{2b_3}$$

Douglas



$$\sim \frac{R^{2b_3}}{a^{2b_3}}$$

Number of diff. fluxes grows like $(|N_{03}|)^{b_3}$

$b_3 \sim 10^2 - 10^3$ (Typical) can be large

F. theory

$$N_{03} \rightarrow \frac{\chi(\chi_4)}{24}$$

Moduli stabilisation : $G_3 = F_3 - \Phi H_3$

$$\star_6 G(3) = i G(3)$$

$$\sqrt{\frac{2}{3}} \lambda^3 [m_1 m_2 m_3 m_4 m_5 m_6] \underset{1/\lambda}{\widehat{g}^{m_4 n_4}} \underset{1/\lambda}{\widehat{g}^{m_5 n_5}} \underset{1/\lambda}{\widehat{g}^{m_6 n_6}} (\widehat{G}_3)_{n_4 n_5 n_6} = i (G_3)_{m_1 m_2 m_3}$$

$$\widehat{g}_{mn}, \Phi$$

$$\widehat{g}_{mn} \rightarrow \lambda \widehat{g}_{mn}$$

Moduli scalar field arise from \hat{g}_{mn}
Some of these fixed.